

# **A COUPLED 1D AND 2D SPH-SWES MODEL FOR OPEN CHANNEL FLOWS IN COMPLEX CHANNEL GEOMETRIES**

**SPEAKER: DR. KAO-HUA CHANG (張高華)**

**2016.12.01**

**ACKNOWLEDGEMENT:**

**PROF. TONY WEN-HANN SHEU (許文翰 教授);**

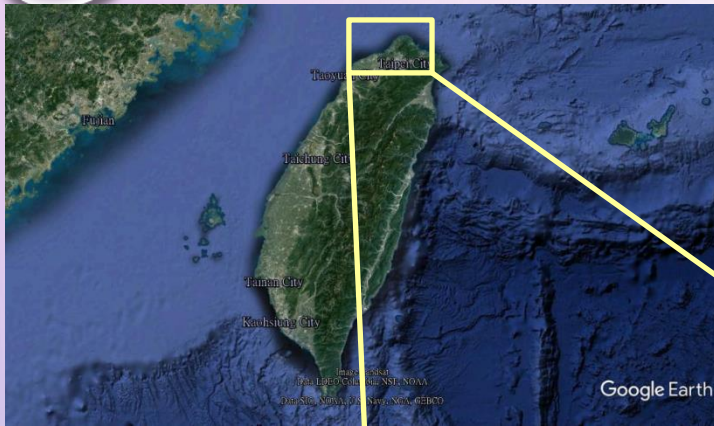
**PROF. TSANG-JUNG CHANG (張倉榮 教授)**

# OUTLINE

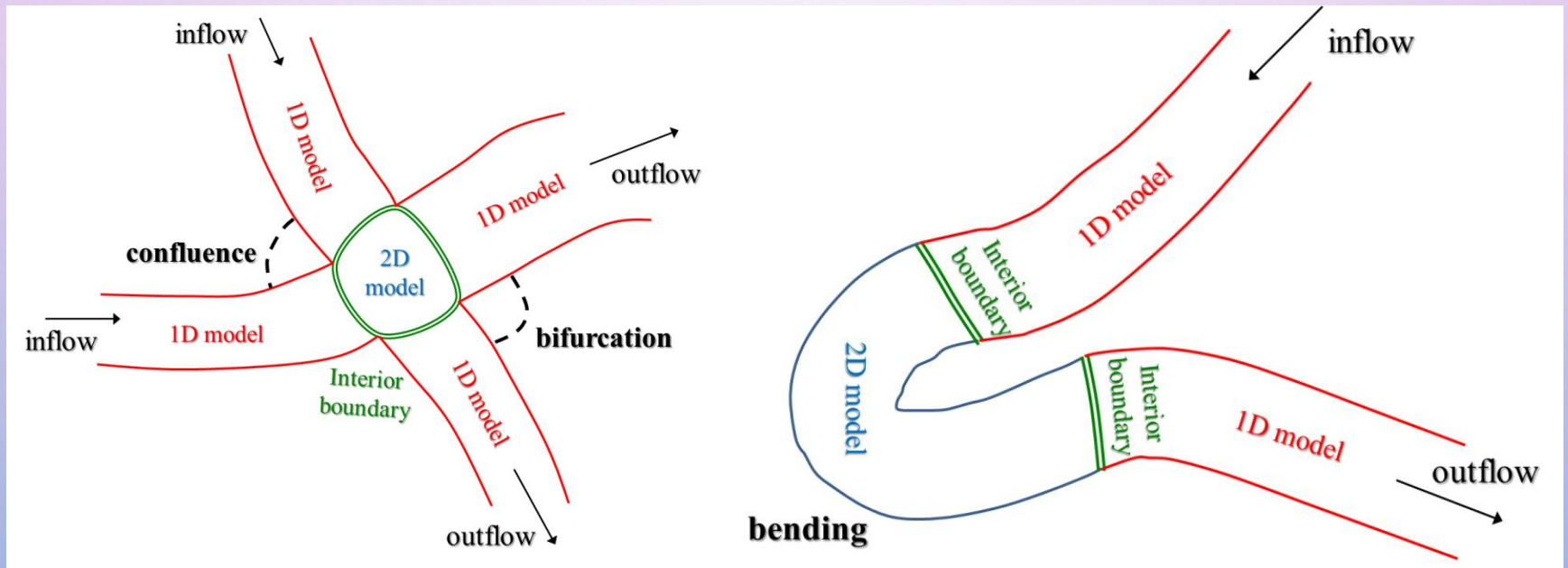
1. Motivation
2. SPH-SWEs Model
3. Boundary Treatments
4. Results and Discussion
5. Summary

# MOTIVATION

- ❑ In hydraulic engineering, **river flooding** (overbank flow and mud flow) and **river pollution** are both concerned problems.
- ❑ How to predict **discharge**, **depth** and **concentration** is a key issue.



□ For large river systems. **1D hydrodynamic models** are extensively used due to their **efficiency**. Besides, **2D hydrodynamic models** are also adopted to enhance the **accuracy**.

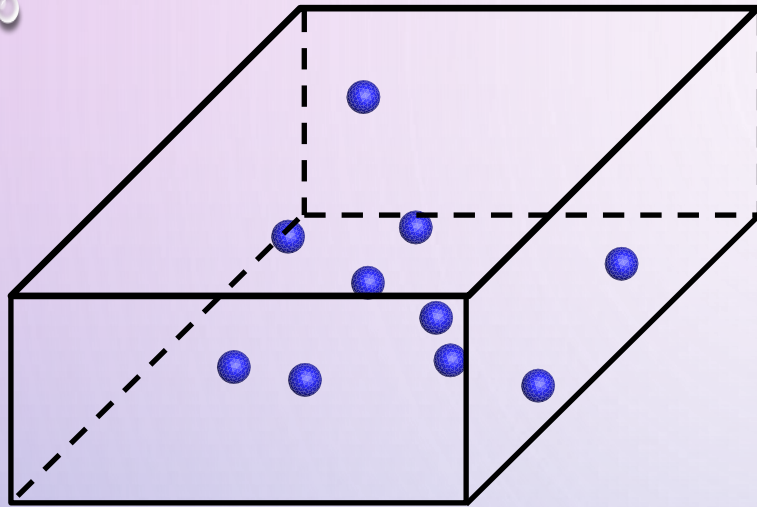


## Goal

➤ Combining the efficiency of 1D model and the accuracy of 2D model.

# **SPH-SWE<sub>s</sub> MODEL**

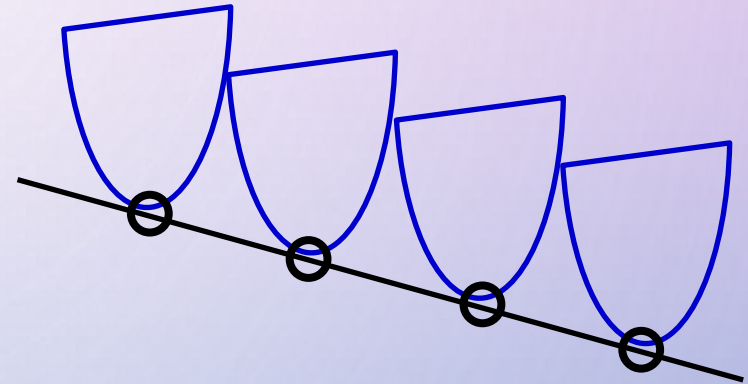
# SHALLOW WATER EQUATIONS



## *3D Navier-Stokes equations*

Density  $\rho = \rho_w$

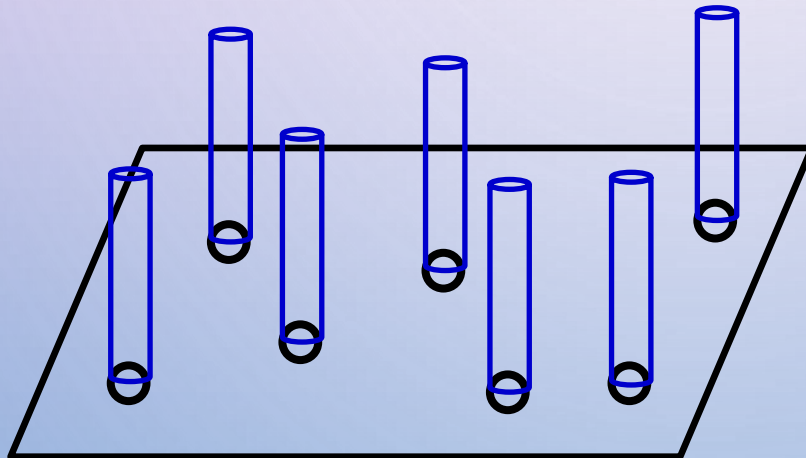
Mass  $m = \rho V = \rho_w (\Delta x_0)^3$



## *1D shallow water equations*

Density  $\rho = \rho_w A$  *Area-averaged*

Mass  $m = \rho V = \rho_w A \Delta x_0$



## *2D shallow water equations*

Density  $\rho = \rho_w d_w$  *Depth-averaged*

Mass  $m = \rho V = \rho_w d_w (\Delta x_0)^2$

Pressure is assumed to be hydrostatic

$$P = \gamma d_w = \rho g d_w = \rho c^2,$$

where  $c$  is the celerity.  $c = \sqrt{g d_w}$

## 1D shallow water equations

$$\frac{DA}{Dt} = -A \frac{\partial u}{\partial x}$$

$$\frac{DQ}{Dt} = -Q \frac{\partial u}{\partial x} - gA \frac{\partial (d_w + z_b)}{\partial x} - gAS_f$$

## 2D shallow water equations

$$\frac{Dd_w}{Dt} = -d_w \nabla \cdot \vec{U}$$

$$\frac{D\vec{U}}{Dt} = -g \nabla (d_w + z_b) - g\vec{S}_f$$

$A$  is the wetted cross-section area,  
 $d_w$  is the water depth,  
 $Q$  is the discharge,  
 $U$  is the velocity,  
 $z_b$  is the ground elevation,  
 $S_f$  is the friction slope and  
 $g$  is the gravity acceleration.

In the above,

$$\vec{U} = (u, v), \quad \nabla = \frac{\partial}{\partial x} \vec{x} + \frac{\partial}{\partial y} \vec{y}, \quad \frac{D(\quad)}{Dt} = \frac{\partial(\quad)}{\partial t} + \vec{U} \cdot \nabla(\quad)$$



# SMOOTHED PARTICLE HYDRODYNAMICS

## Kernel Function ( $\omega$ )

- (1) The integral of a kernel function within its compact domain should be equal to one.

$$\int_{\Omega} \omega(r_{ij}, h_i) dV = 1$$

- (2) As the smoothing length  $h$  approaching zero, the kernel function will become a Dirac delta function.

$$\lim_{h_i \rightarrow 0} \omega(r_{ij}, h_i) = \delta(r_{ij})$$

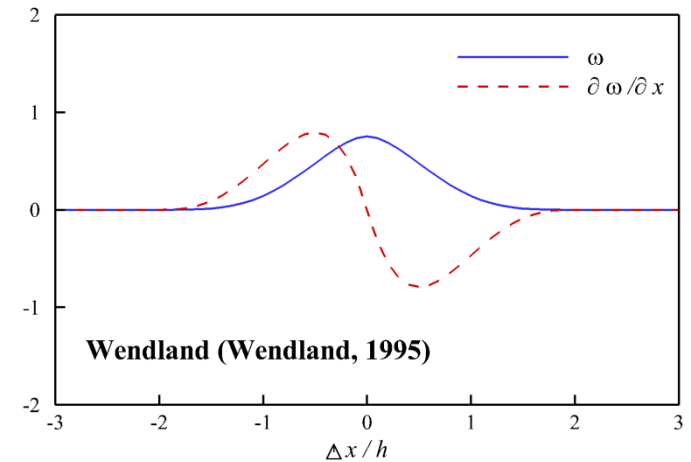
- (3) The kernel function should be compactly supported.

$$\omega(r_{ij}, h_i) = 0 \quad \text{when } r_{ij} > \kappa h_i$$

- (4) The kernel function is assumed to be symmetric.

$$\omega(r_{ij}, h_i) = \omega(r_{ji}, h_i)$$

$$\nabla_i \omega(r_{ij}, h_i) = -\nabla_j \omega(r_{ij}, h_i)$$



$$\omega(s, h) = \frac{3}{4h} \begin{cases} (1+2s)(1-\frac{s}{2})^4 & 0 \leq s \leq 2 \\ 0 & s > 2 \end{cases}$$

where  $s = \frac{\Delta x}{h}$

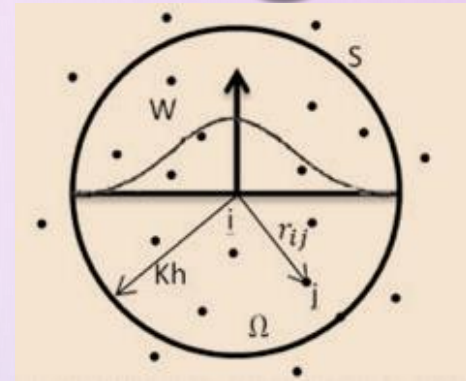
## ◆ Summation operator

$$\phi_i \doteq \int_{\Omega} \phi(\vec{r}) \delta(\vec{r}_i - \vec{r}) dV$$

$$\langle \phi \rangle_i = \int_{\Omega} \phi(\vec{r}) \omega(|\vec{r}_i - \vec{r}|, h) dV$$

$$= \sum_j \frac{m_j}{\rho_j} \phi_j \omega(r_{ij}, h_i) \quad \longleftarrow \text{Discretization into particles}$$

where  $m$  is the mass of a particle,  
 $\rho$  is the density of a particle,  
 $\omega$  is the kernel function,  
 $h$  is the smoothing length and  
 $\langle \phi \rangle_i$  denotes approximated  $\phi_i$ .



## ◆ Differential operator

$$\langle \nabla \phi \rangle_i = \int_{\Omega} \nabla \phi(\vec{r}) \omega(|\vec{r}_i - \vec{r}|, h) dV$$

$$= \int_{\Omega} \nabla [\phi(\vec{r}) \omega(|\vec{r}_i - \vec{r}|, h)] dV - \int_{\Omega} \phi(\vec{r}) \nabla \omega(|\vec{r}_i - \vec{r}|, h) dV$$

$$= \int_{\partial\Omega} \phi(\vec{r}) \omega(|\vec{r}_i - \vec{r}|, h) \vec{n} dS - \int_{\Omega} \phi(\vec{r}) \nabla \omega(|\vec{r}_i - \vec{r}|, h) dV$$

$$= - \int_{\Omega} \phi(\vec{r}) \nabla \omega(|\vec{r}_i - \vec{r}|, h) dV \quad \nabla(\ ) = -\nabla_i(\ )$$

$$= \sum_j \frac{m_j}{\rho_j} \phi_j \nabla_i \omega(r_{ij}, h_i) \quad \longleftarrow \text{Discretization into particles}$$

$$\nabla(\phi\omega) = \phi\nabla\omega + \omega\nabla\phi$$

## ◆ Divergent operator in the asymmetric form

$$\nabla \cdot \bar{\phi} = \frac{1}{\rho} \left[ \nabla \cdot (\rho \bar{\phi}) - \bar{\phi} \cdot \nabla \rho \right]$$

Using the differential operator  $\langle \nabla \phi \rangle_i = \sum_j \frac{m_j}{\rho_j} \phi_j \nabla_i \omega_{ij}$

$$\langle \nabla \cdot \bar{\phi} \rangle_i = \frac{1}{\rho_i} \sum_j m_j (\bar{\phi}_i - \bar{\phi}_j) \cdot \nabla_i \omega_{ij}$$

## ◆ Gradient operator in the symmetric form

$$\nabla \phi = \rho \left[ \nabla \left( \frac{\phi}{\rho} \right) + \frac{\phi}{\rho^2} \nabla \rho \right]$$

Using the differential operator  $\langle \nabla \phi \rangle_i = \sum_j \frac{m_j}{\rho_j} \phi_j \nabla_i \omega_{ij}$

$$\langle \nabla \phi \rangle_i = \rho_i \sum_j m_j \left( \frac{\phi_i}{\rho_i^2} + \frac{\phi_j}{\rho_j^2} \right) \nabla_i \omega_{ij}$$

# DIFFUSIVE TERMS

## I. Artificial viscous term

$$\left( \frac{D\vec{U}}{Dt} \right)_i = -g \sum_j m_j \frac{d_{w,i} + d_{w,j}}{d_{w,j}} \nabla_i \omega_{ij} - g \vec{S}_f + \sum_j \nu_{ij} V_j (\vec{U}_i - \vec{U}_j) \frac{\vec{r}_{ij} \cdot \nabla_i \omega_{ij}}{|\vec{r}_{ij}|^2}$$

**Kinetic viscosity**  $\nu_{ij} = \alpha h_{ij} c_{ij}$ , where  $h_{ij} = 0.5(h_i + h_j)$ ,  $c_{ij} = 0.5(c_i + c_j)$ ,  
 $\alpha = 0.2$ .

## 2. Density diffusive term

$$\left( \frac{Dd_w}{Dt} \right)_i = \sum_j m_j (\vec{U}_i - \vec{U}_j) \cdot \nabla_i \omega_{ij} + \sum_j D_{ij} V_j (d_{w,i} - d_{w,j}) \frac{\vec{r}_{ij} \cdot \nabla_i \omega_{ij}}{|\vec{r}_{ij}|^2}$$

$D_{ij} = \beta h_{ij} c_{ij}$ , where  $h_{ij} = 0.5(h_i + h_j)$ ,  $c_{ij} = 0.5(c_i + c_j)$ ,  $\beta = 0.2$ .

Ref. 1) JJ Monaghan, Rep Prog Phys, 68: 1703-1759, 2005.

2) M Antuono et al., Comput Phys Commun, 183: 2570-2580, 2012.

# TIME MARCHING METHOD

## Modified Verlet method

$$\begin{aligned} \vec{U}_i^{n+1/2} &= \vec{U}_i^n + 0.5 \cdot \Delta t \left( \frac{D\vec{U}}{Dt} \right)_i^n & \Longrightarrow & \vec{U}_i^{n+1} = \vec{U}_i^n + \Delta t \left( \frac{D\vec{U}}{Dt} \right)_i^{n+1/2} \\ \rho_i^{n+1/2} &= \rho_i^n + 0.5 \cdot \Delta t \left( \frac{D\rho}{Dt} \right)_i^n & & \Downarrow \\ & & & \vec{r}_i^{n+1} = \vec{r}_i^{n+1/2} + 0.5 \cdot \Delta t \cdot \vec{U}_i^{n+1} \\ \vec{r}_i^{n+1/2} &= \vec{r}_i^n + 0.5 \cdot \Delta t \cdot \vec{U}_i^n & & \Downarrow \\ & & & \rho_i^{n+1} = \rho_i^{n+1/2} + 0.5 \cdot \Delta t \left( \frac{D\rho}{Dt} \right)_i^{n+1} \end{aligned}$$

$$\Delta t \leq CFL \cdot \min_i \left( \frac{h_i}{|\vec{U}_i| + c_i} \right), \text{ where } CFL = 0.8.$$

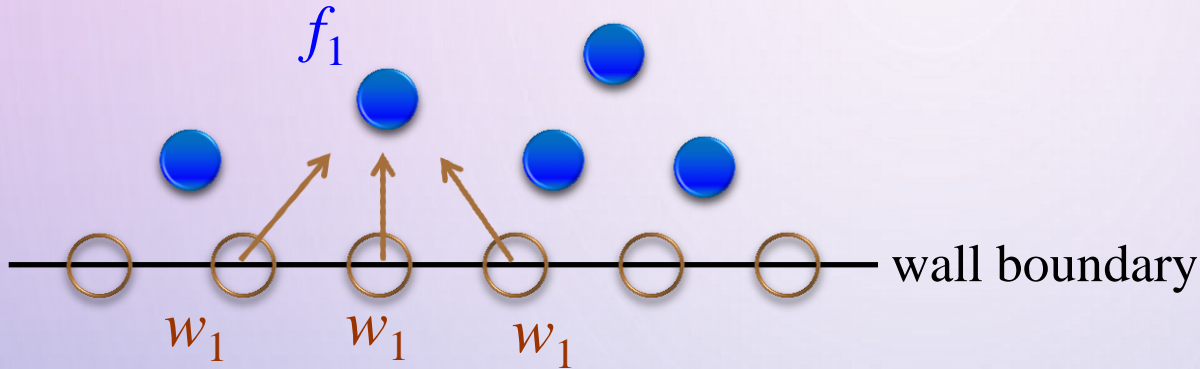
Ref. 1) D Molteni, A Colagrossi, *Comput Phys Commun*, 180: 861-872, 2009.

# **BOUNDARY TREATMENTS**

# WALL BOUNDARY

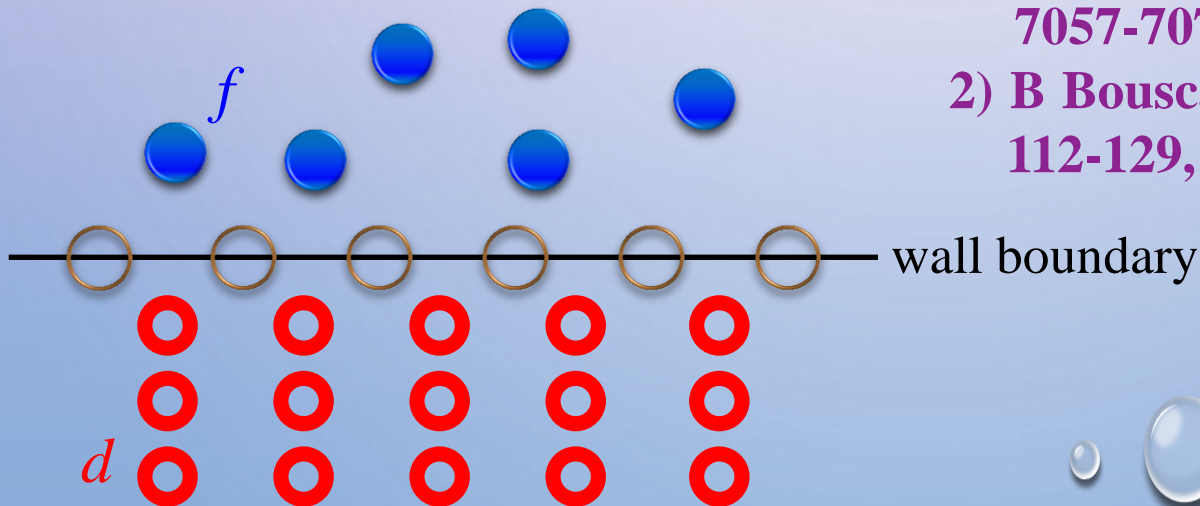
## □ Repellent particle

Ref. 1) JJ Monaghan, A Kos, J Waterw Port C-ASCE, 125: 145-154, 1999.



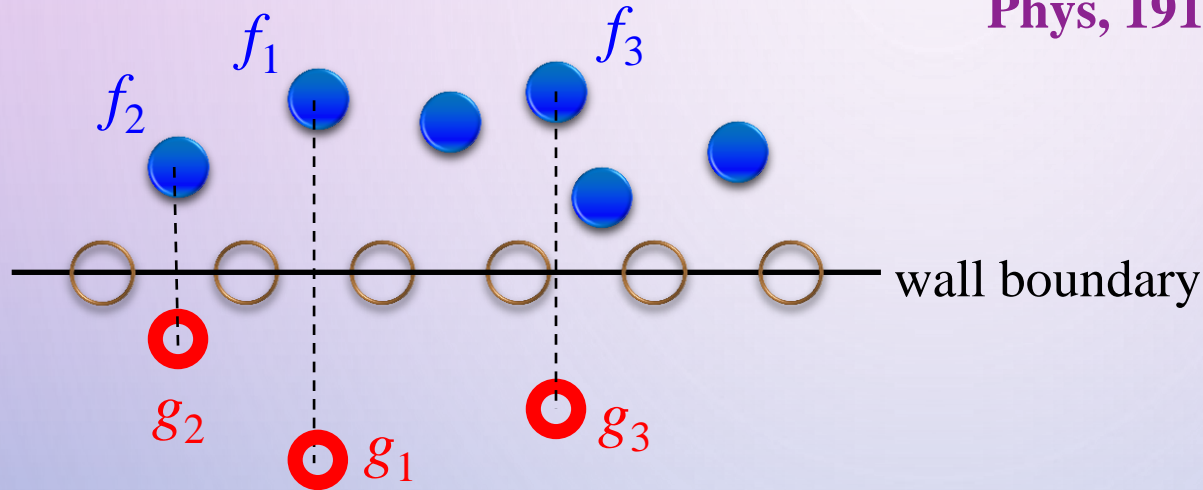
## □ Dummy particle

Ref. 1) S Adami et al., J Comput Phys, 231: 7057-7075, 2012.  
2) B Bouscasse et al., J Fluid Struct, 42: 112-129, 2013.



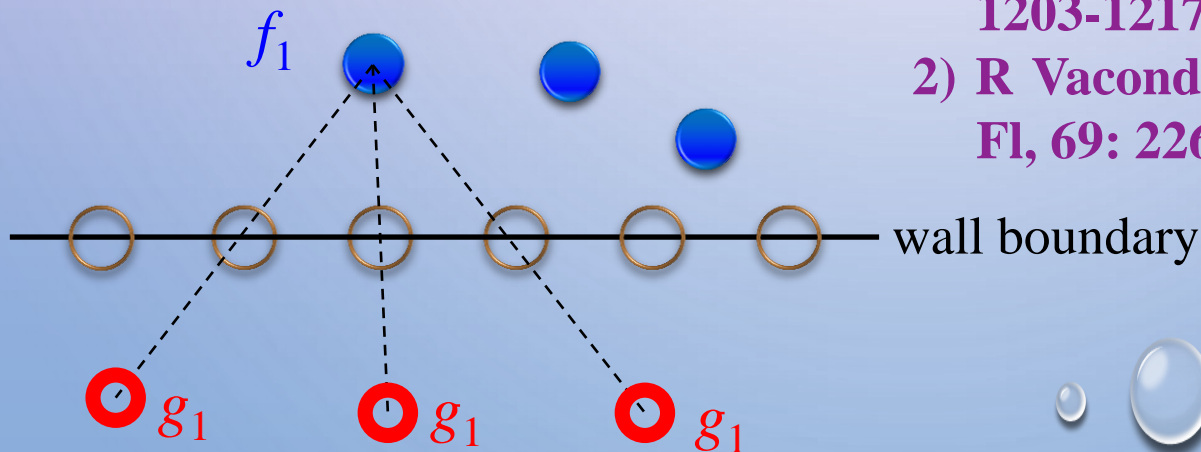
## □ Ghost particle

### ✓ Plane-symmetry



Ref. 1) A Colagrossi, M Landrini, J Comput Phys, 191: 448-475, 2003.

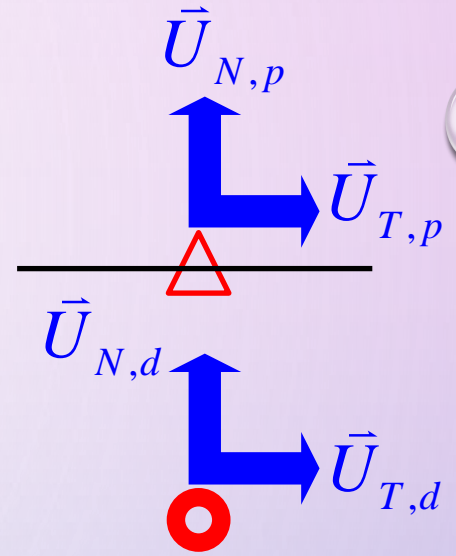
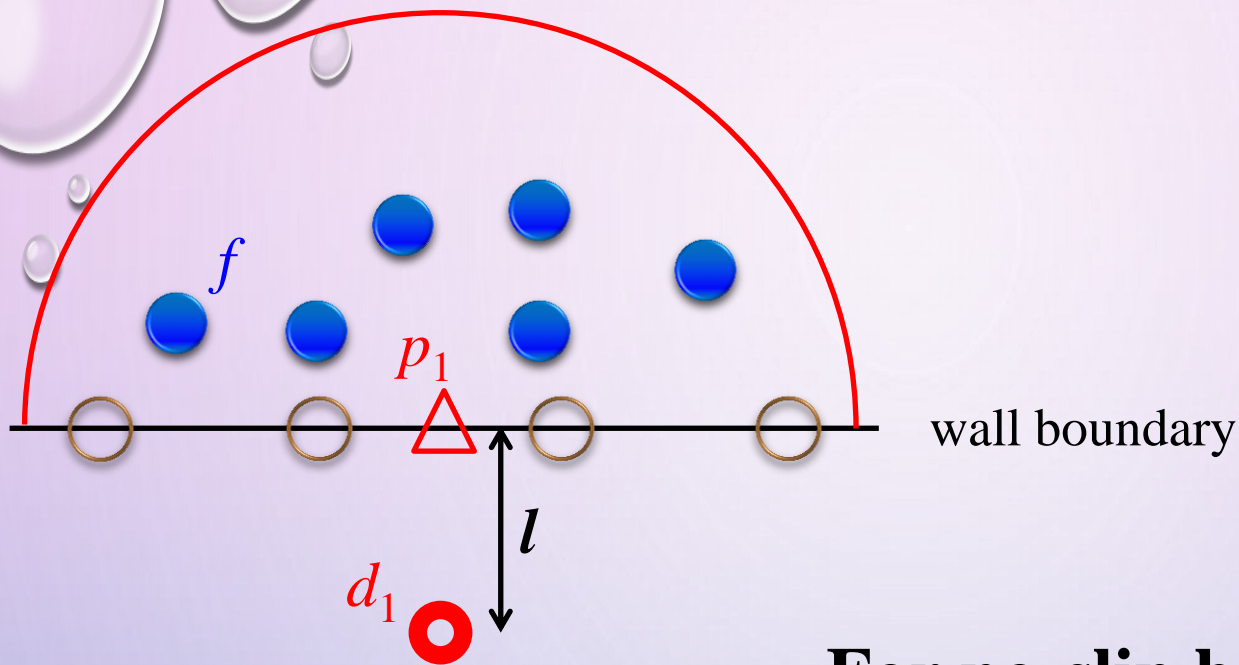
### ✓ Point-symmetry



Ref. 1) A Ferrari et al., Comput Fluids, 38: 1203-1217, 2009.

2) R Vacondio et al., Int J Numer Meth Fl, 69: 226-253, 2012.





**For no-slip boundary**

$$\vec{U}_{i \in p} = \frac{\sum_{j \in f} V_j \vec{U}_j \omega_{ij}}{\sum_{j \in f} V_j \omega_{ij}}$$

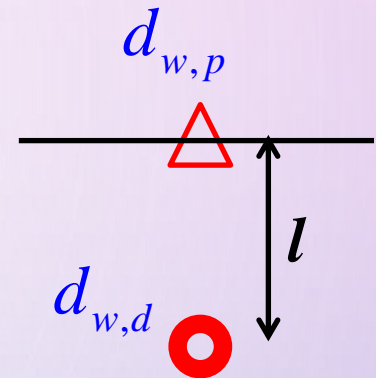
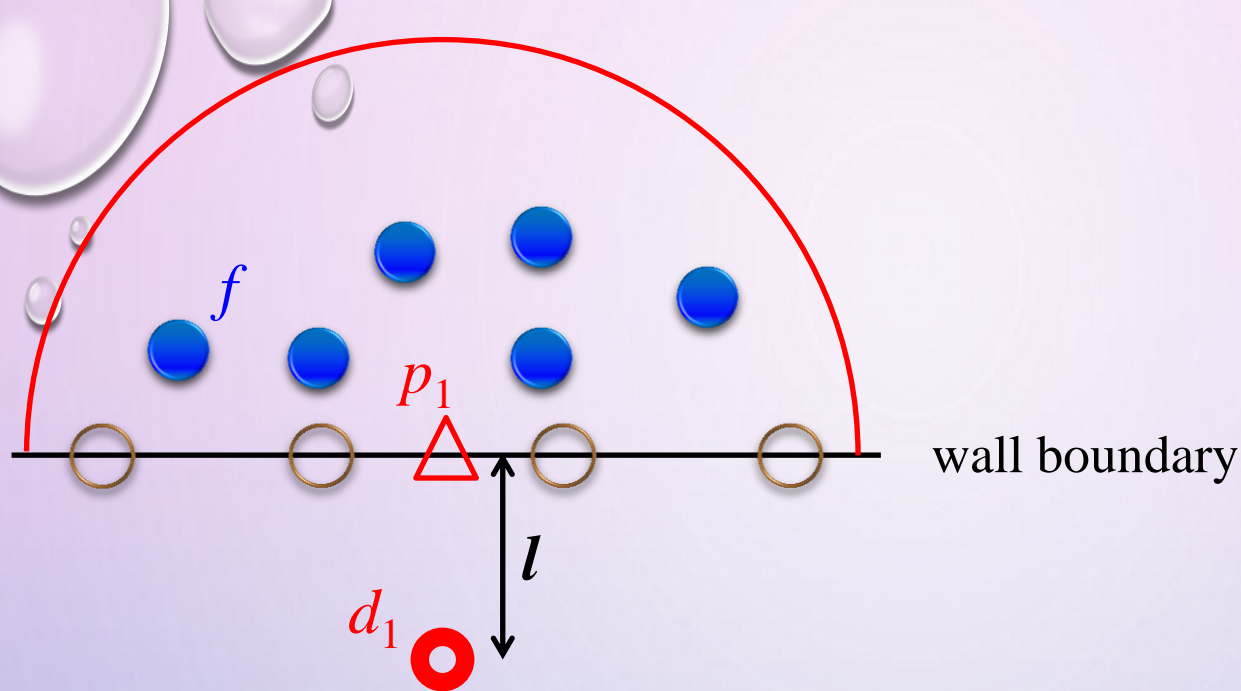


$$\nabla \cdot \vec{U} \begin{cases} \vec{U}_{N,p} = 0 \\ \vec{U}_{T,p} = 0 \end{cases}$$

$$= \sum_{j \in f} V_j \vec{U}_j \tilde{\omega}_{ij}$$

$$\nabla^2 \vec{U} \begin{cases} \vec{U}_{N,d} = \vec{U}_{N,p} \\ \vec{U}_{T,d} = -\vec{U}_{T,p} \end{cases}$$

- Ref. 1) S Adami et al., J Comput Phys, 231: 7057-7075, 2012.  
 2) B Bouscasse et al., J Fluid Struct, 42: 112-129, 2013.



## Impermeable boundary condition

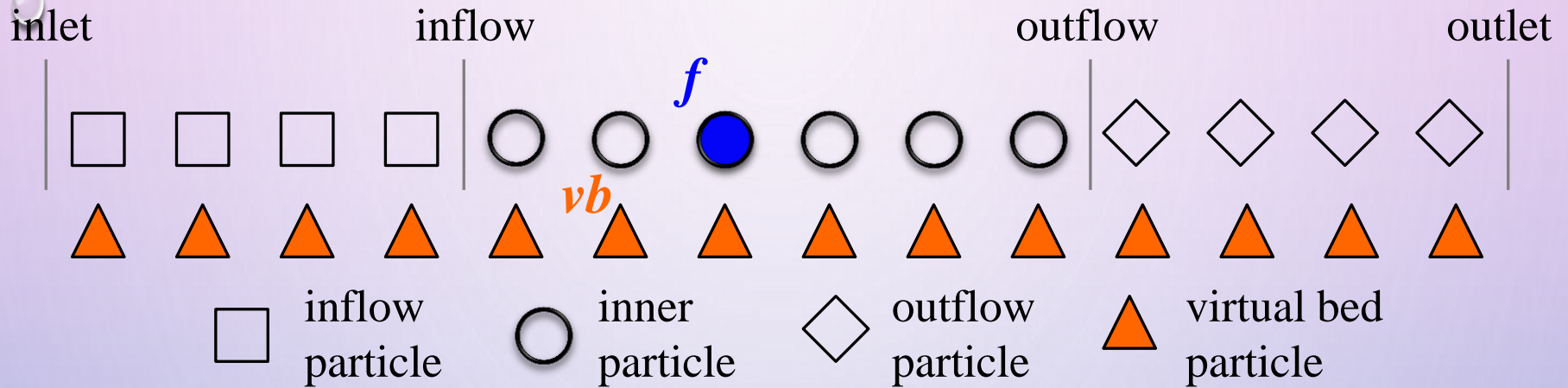
$$\frac{d\bar{u}_d}{dt} \cdot \bar{N} = \left[ -g\nabla(d_w + z_b) - g\bar{S}_f \right] \cdot \bar{N} = 0$$

$$\Rightarrow \nabla d_w \cdot \bar{N} = \left( -\nabla z_b - \bar{S}_f \right) \cdot \bar{N} \Rightarrow \frac{d_{w,j \in p} - d_{w,i \in d}}{l} = \left[ \sum_{j \in f} V_j \left( -\nabla z_{b,j} - \bar{S}_{f,j} \right) \tilde{\omega}_{ij} \right] \cdot \bar{N}_{ij}$$

$$\Rightarrow d_{w,i \in d} = \sum_{j \in f} V_j \left[ d_{w,j} + \left( -\nabla z_{b,j} - \bar{S}_{f,j} \right) \cdot \bar{N}_{ij} \right] \tilde{\omega}_{ij}$$

- Ref. 1) S Adami et al., J Comput Phys, 231: 7057-7075, 2012.  
 2) B Bouscasse et al., J Fluid Struct, 42: 112-129, 2013.

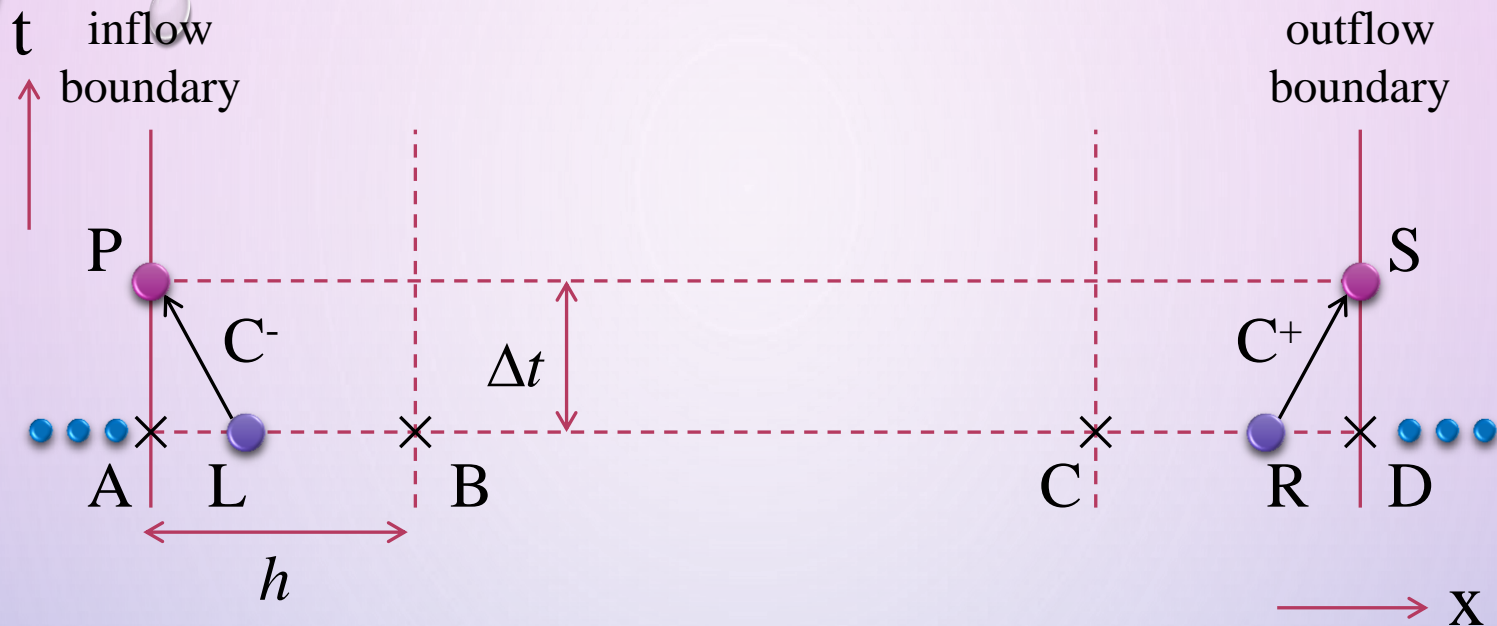
# IN/OUT-FLOW BOUNDARY



$$\langle z_b \rangle_{i \in f} = \sum_{j \in vb} V_j z_{b,j} \tilde{\omega}_{ji}, \quad \langle n_a \rangle_{i \in f} = \sum_{j \in vb} V_j n_{a,j} \tilde{\omega}_{ji},$$

$$\text{where } \tilde{\omega}_{ji} = \frac{\omega_{ji}}{\sum_{j \in vb} V_j z_{b,j} \omega_{ji}}$$

- Ref. 1) R Vacondio et al., Int J Numer Meth Fl, 69: 226-253, 2012.  
 2) TJ Chang, KH Chang, J Hydraul Eng-ASCE, 139:1142-1149, 2013.



## Characteristics equations

$$C^+ : \begin{aligned} u_S &= u_R - \frac{g}{c_R} (d_{w,S} - d_{w,R}) + g (S_{0,R} - S_{f,R}) \Delta t \\ \frac{x_S - x_R}{\Delta t} &= u_R + c_R \end{aligned}$$

$$C^- : \begin{aligned} u_P &= u_L + \frac{g}{c_L} (d_{w,P} - d_{w,L}) + g (S_{0,L} - S_{f,L}) \Delta t \\ \frac{x_P - x_L}{\Delta t} &= u_L + c_L \end{aligned}$$

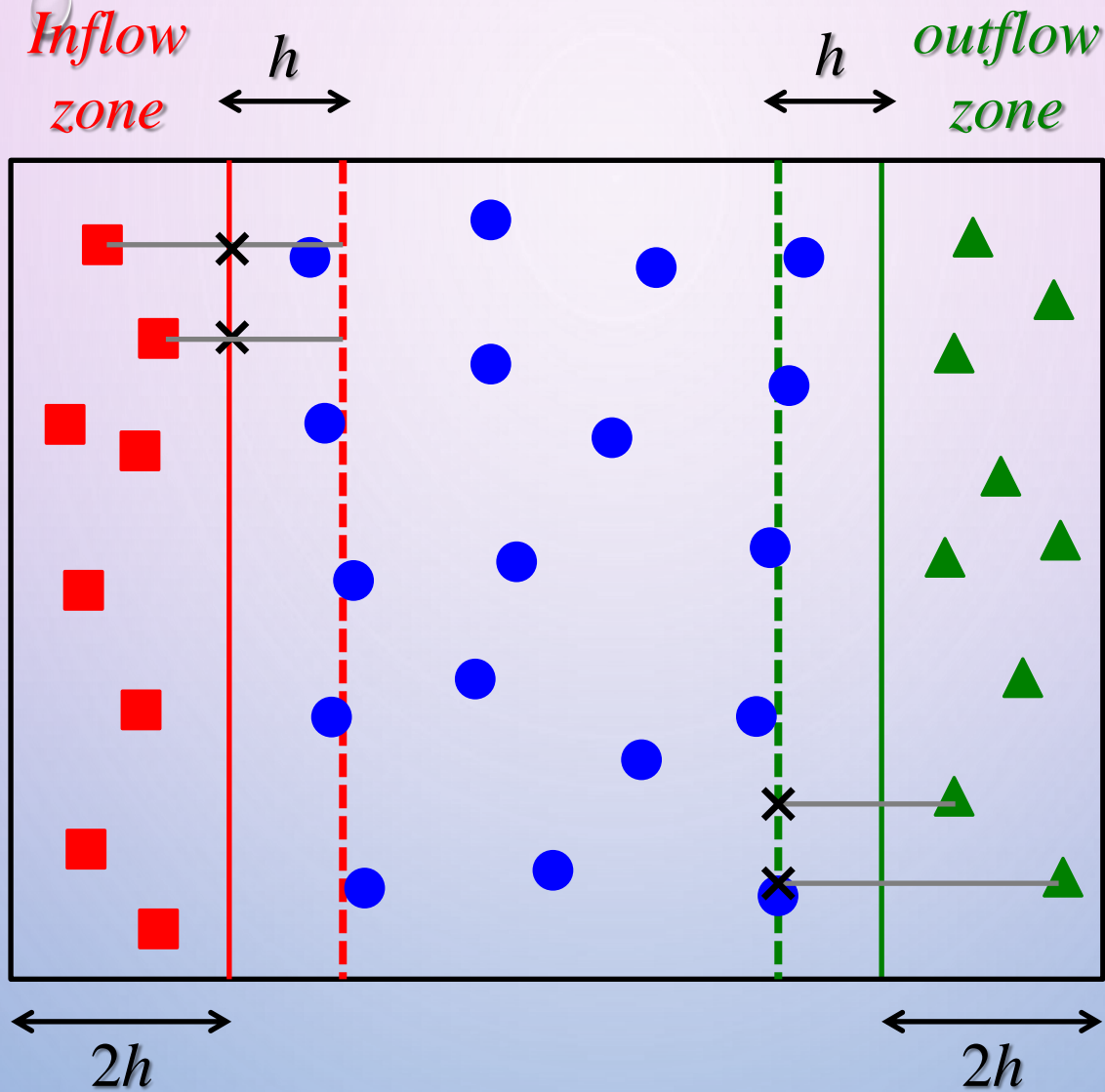
### □ Subcritical flows at outflow boundaries

$$Q_S = A_S u_S = A_S \left[ u_R - \frac{g}{c_R} (d_{w,S} - d_{w,R}) + g (S_{0,R} - S_{f,R}) \Delta t \right]$$

### □ Subcritical flows at inflow boundaries

$$Q_P = A_P u_P = A_P \left[ u_L + \frac{g}{c_L} (d_{w,P} - d_{w,L}) + g (S_{0,L} - S_{f,L}) \Delta t \right]$$

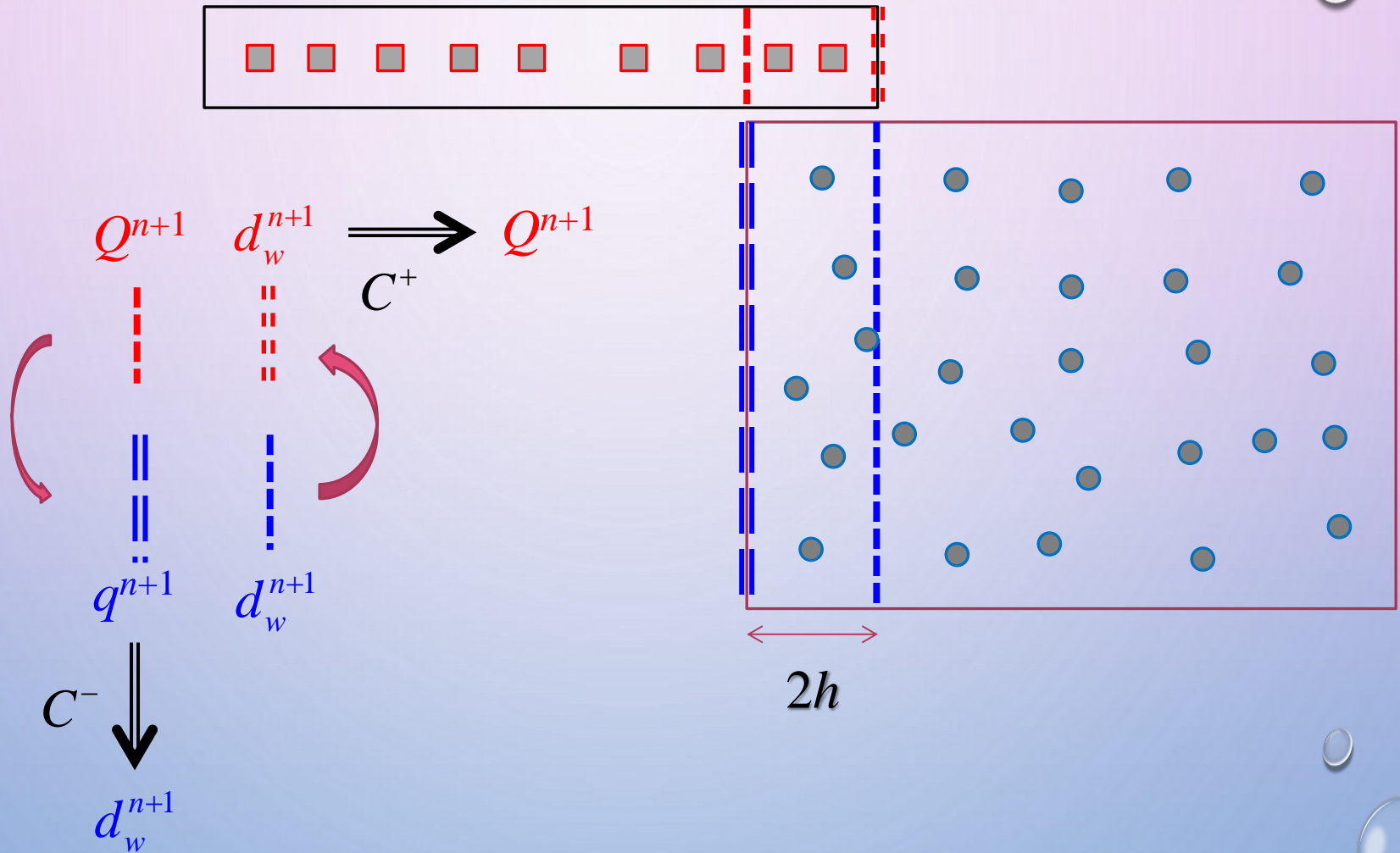
$$f(d_{w,P}) = A_P \left[ u_L + \frac{g}{c_L} (d_{w,P} - d_{w,L}) + g (S_{0,L} - S_{f,L}) \Delta t \right] - Q_P = 0$$



Ref. 1) R Vacondio et al., J Hydraul Eng-ASCE, 138: 530-541, 2012.

2) TJ Chang, KH Chang, J Hydraul Eng-ASCE, 139:1142-1149, 2013.

# COUPLED BOUNDARY

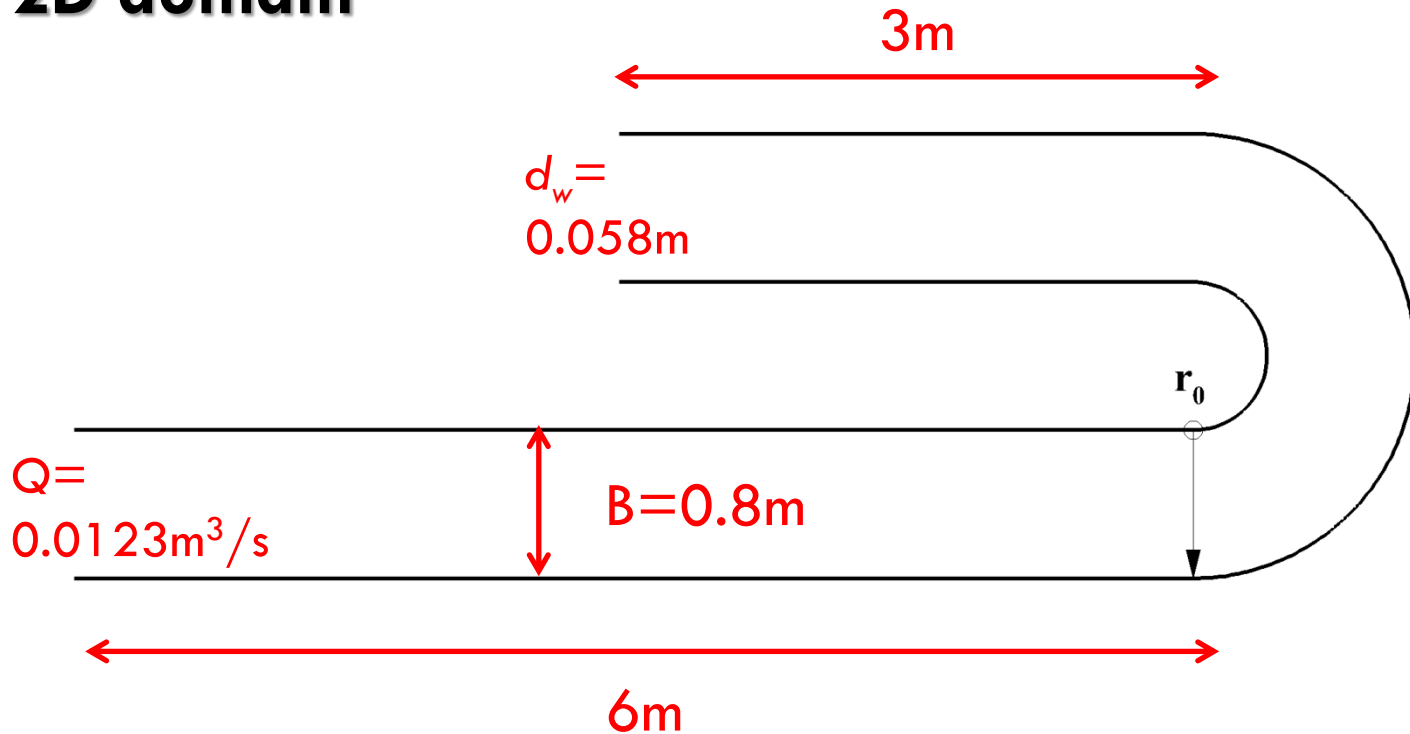


Ref. 1) M Narayanaswamy et al., J Hydraul Res, 48:85-93, 2010.

# RESULTS & DISCUSSION

# CASE 1 – 180° BENDING CHANNEL

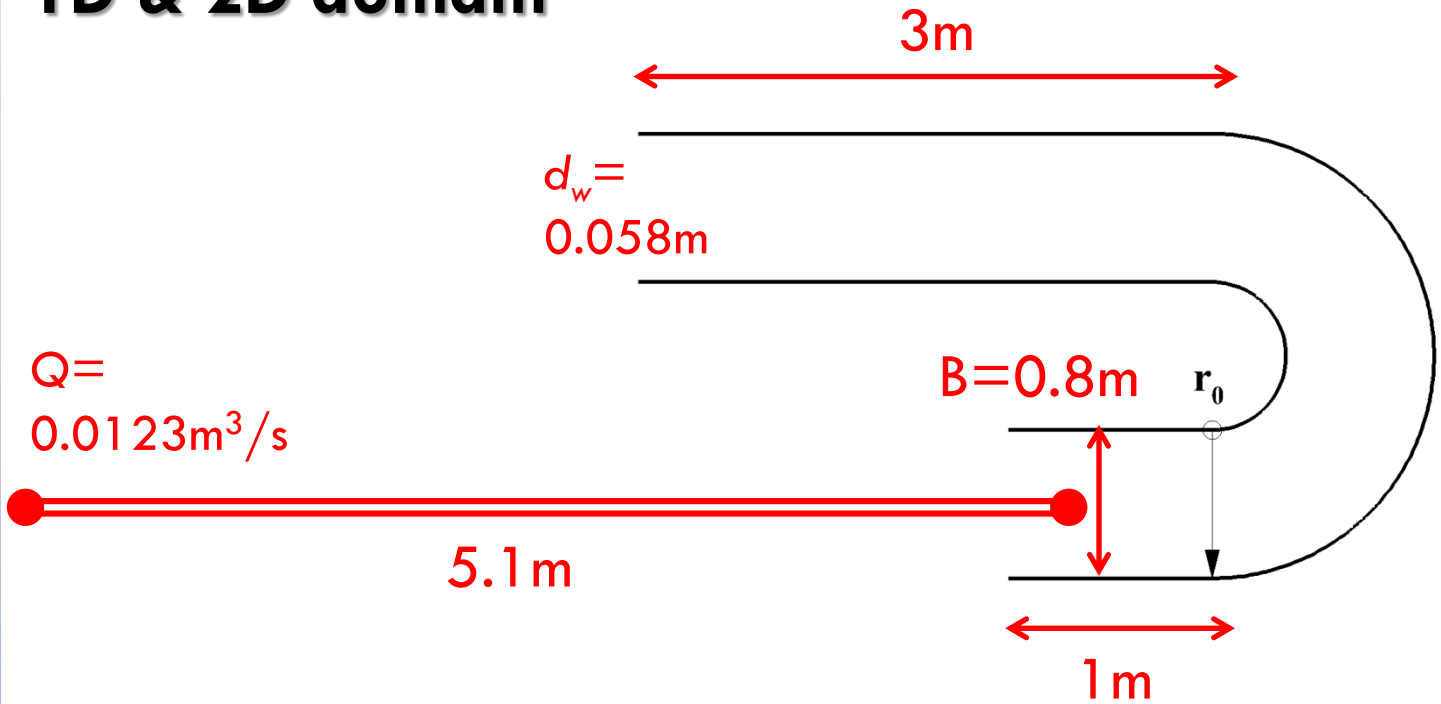
## 2D domain

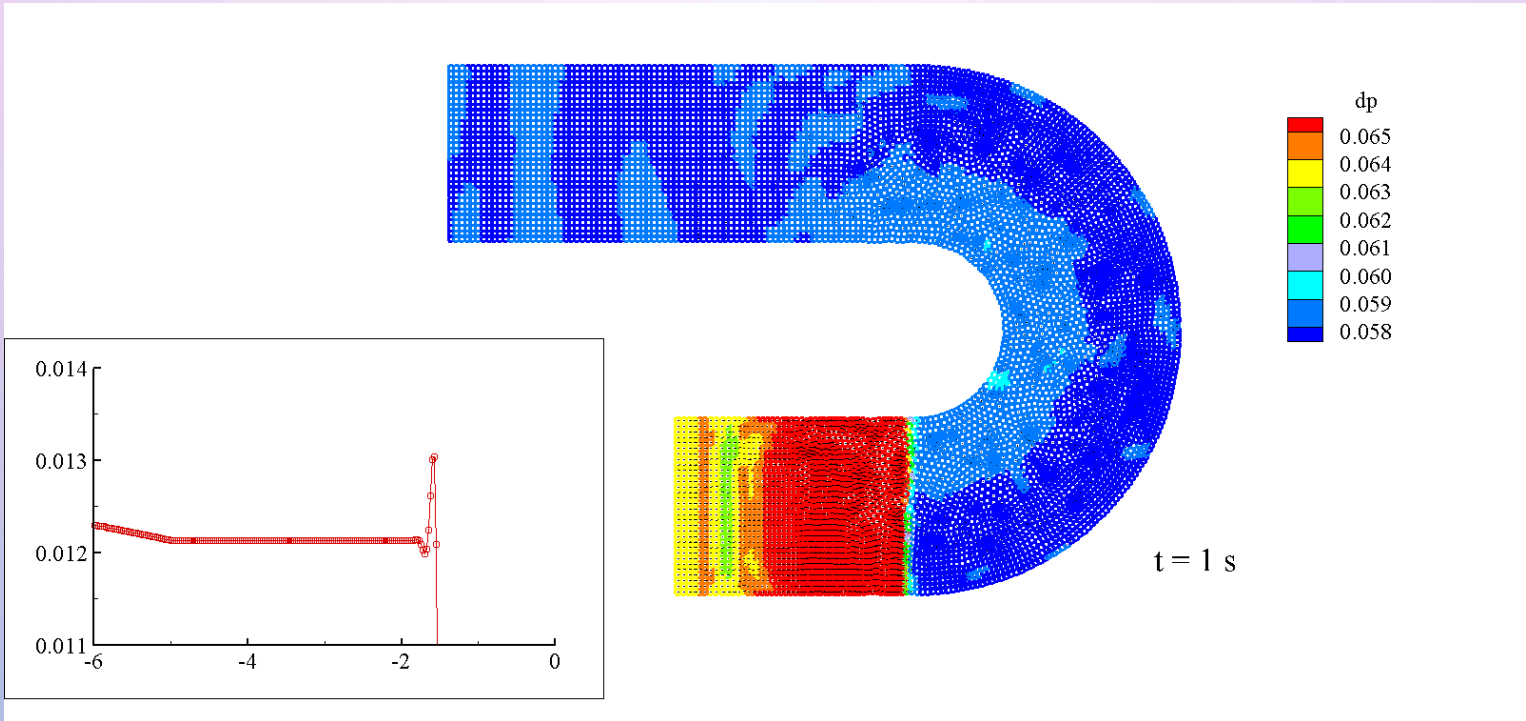


1.  $d_w/B \ll 1$  ( $=0.075$  in this case)
2. No slip boundary condition will produce flow separation from the inner sidewall which was not observed in the laboratory.



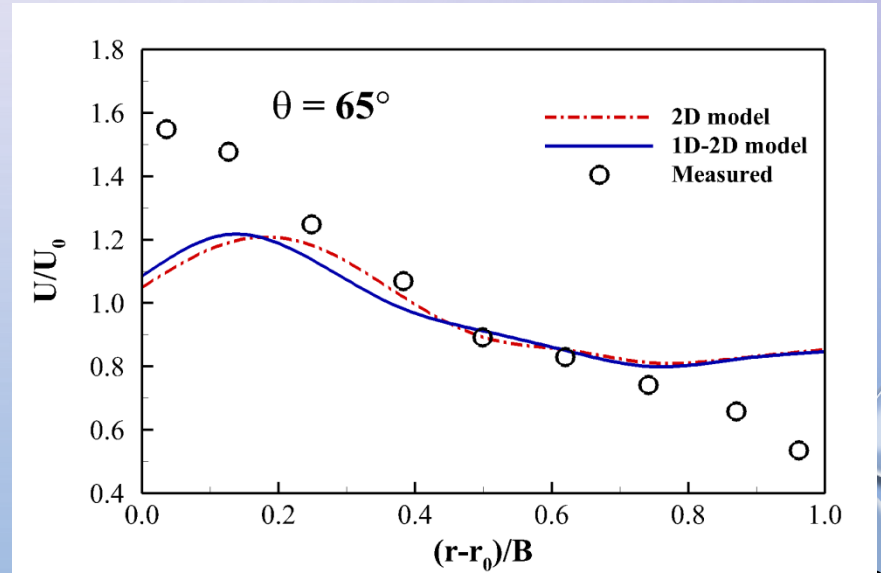
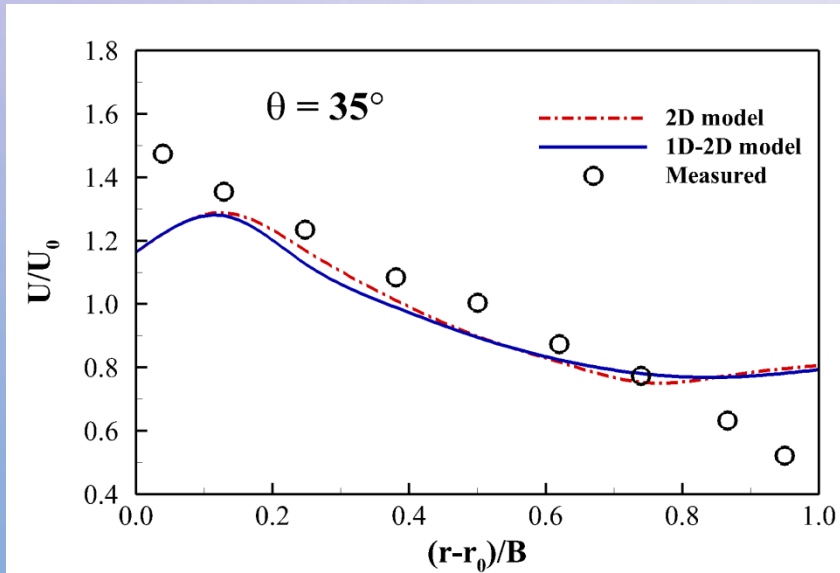
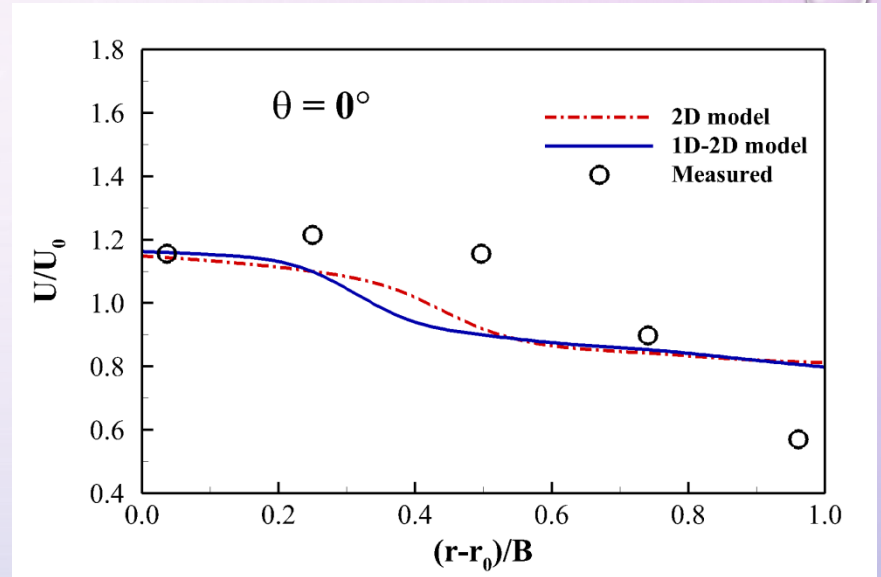
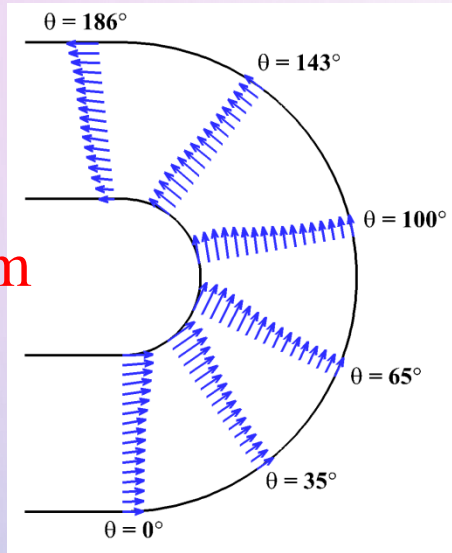
# 1D & 2D domain

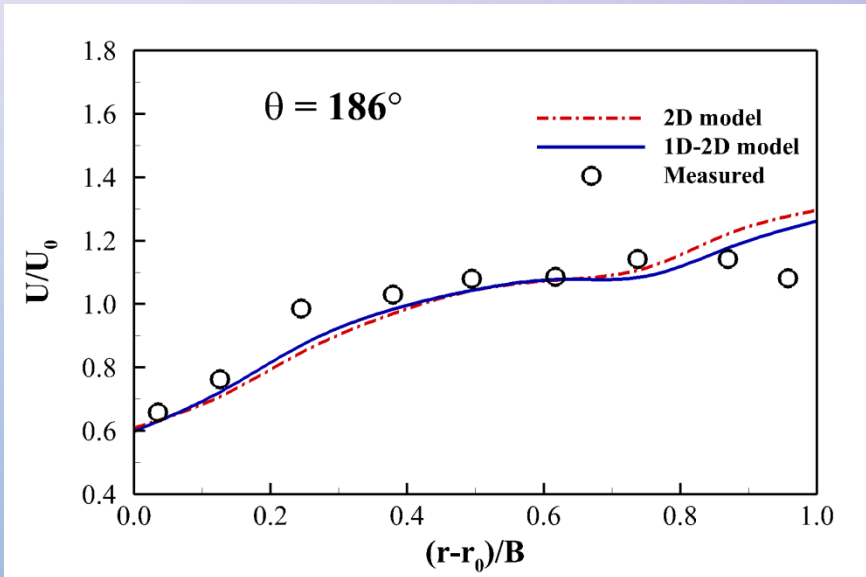
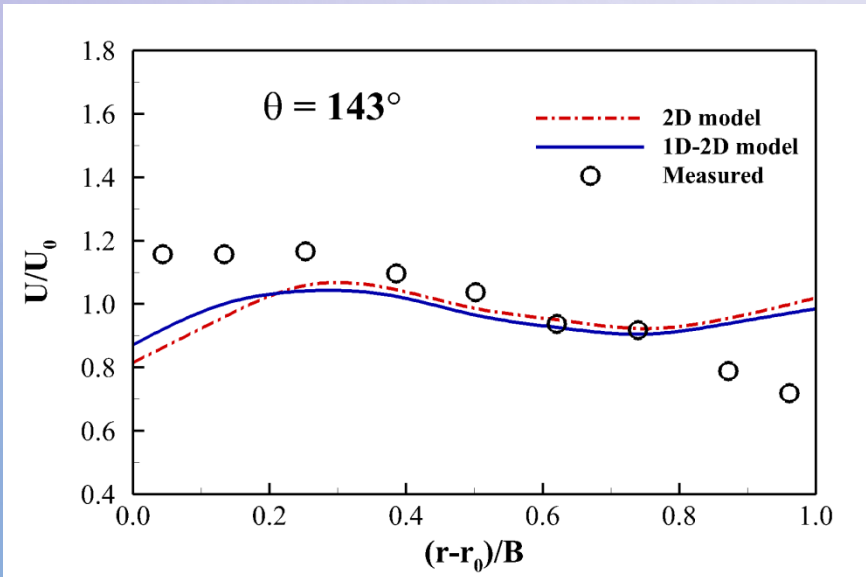
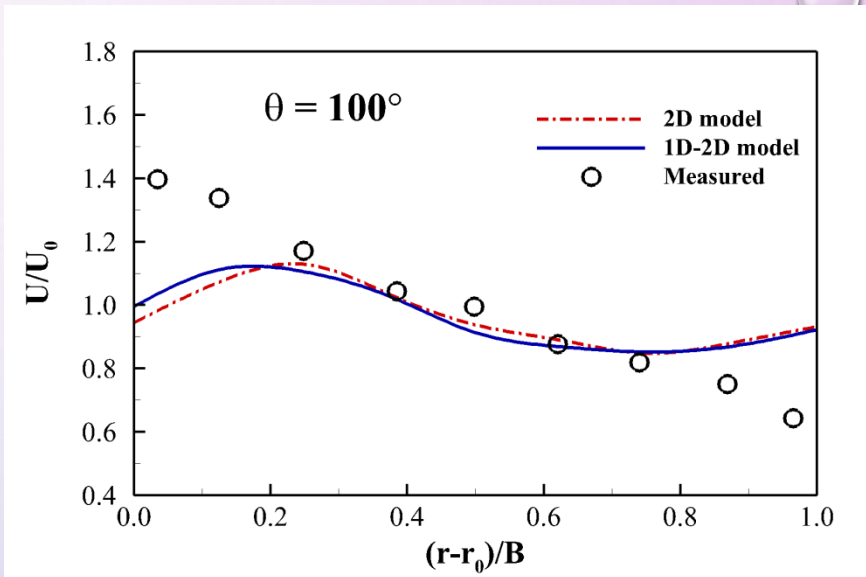
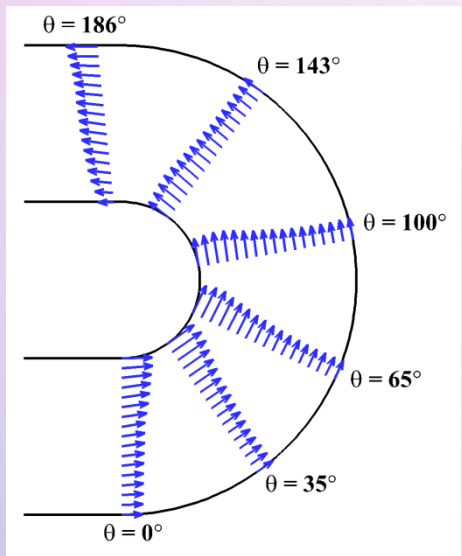




# ➤ Comparison between 2D and 1D-2D SPH-SWEs models

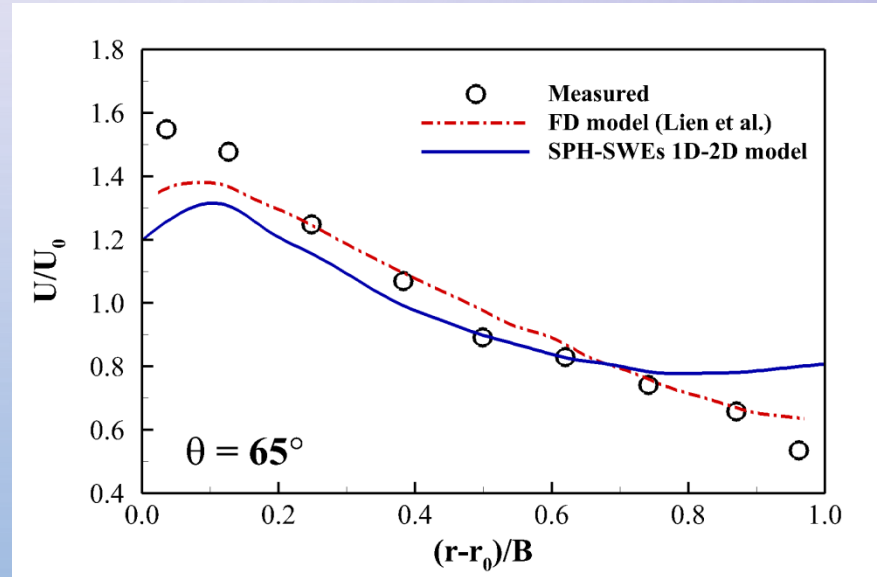
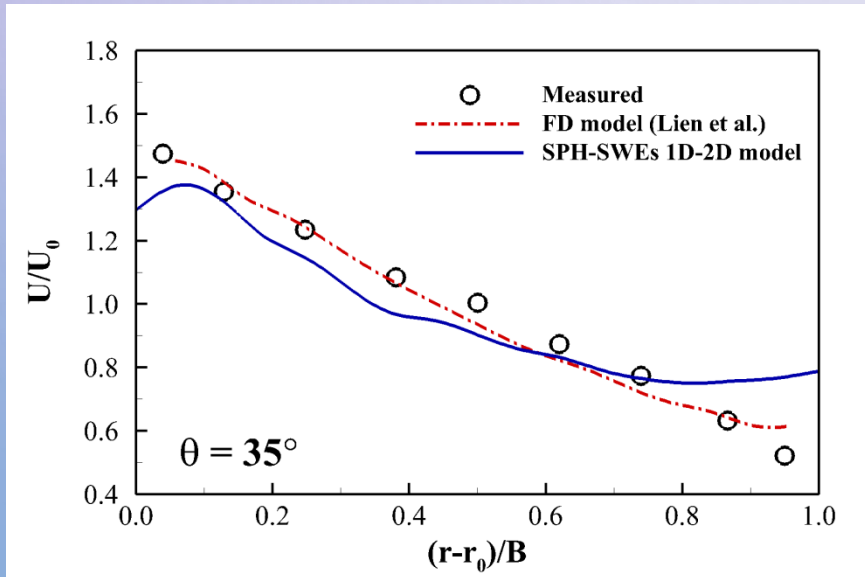
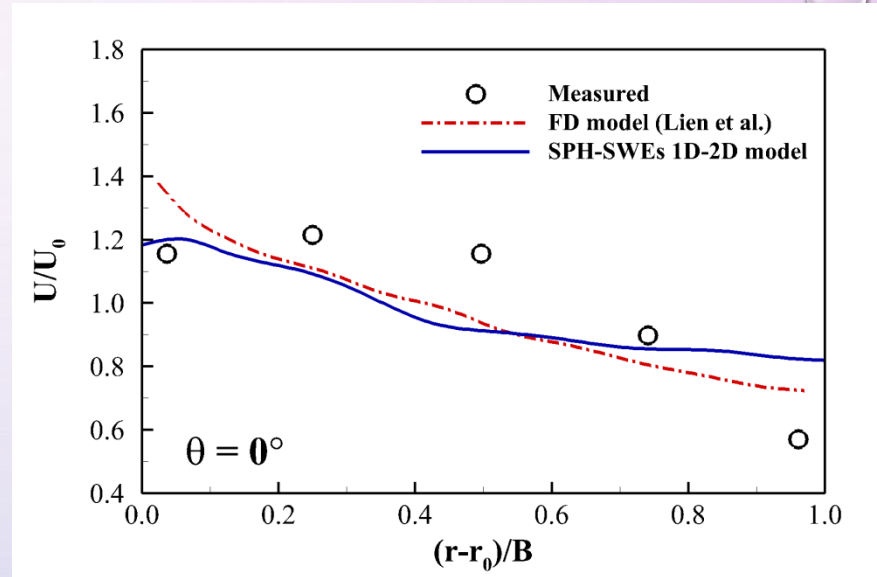
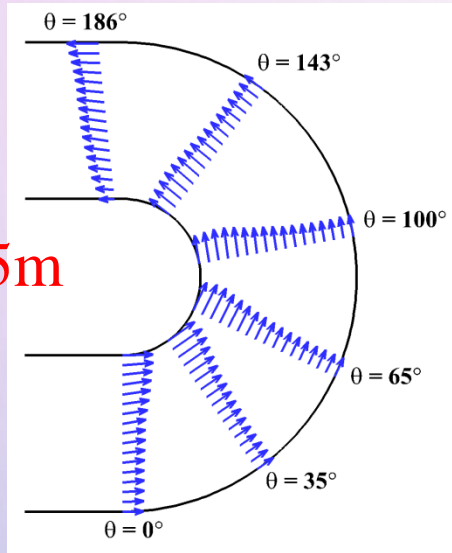
$\Delta x_0 = 0.025\text{m}$

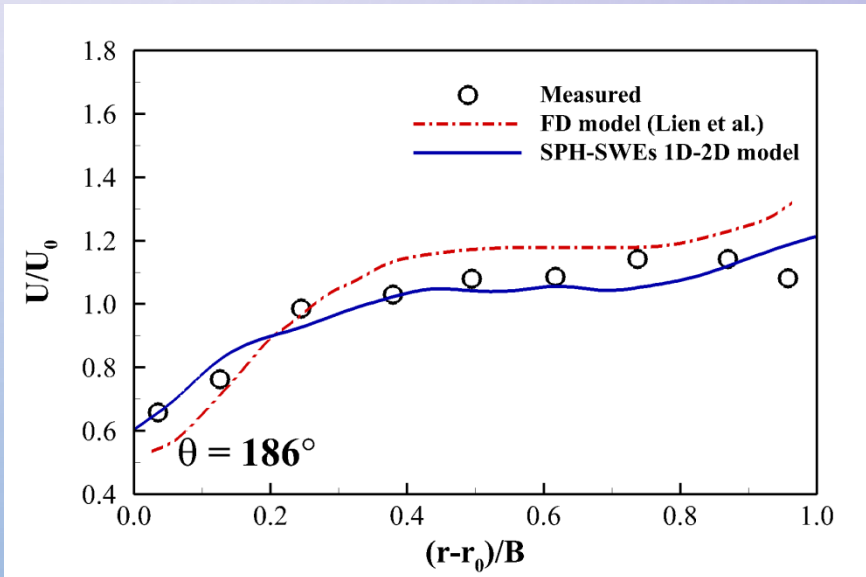
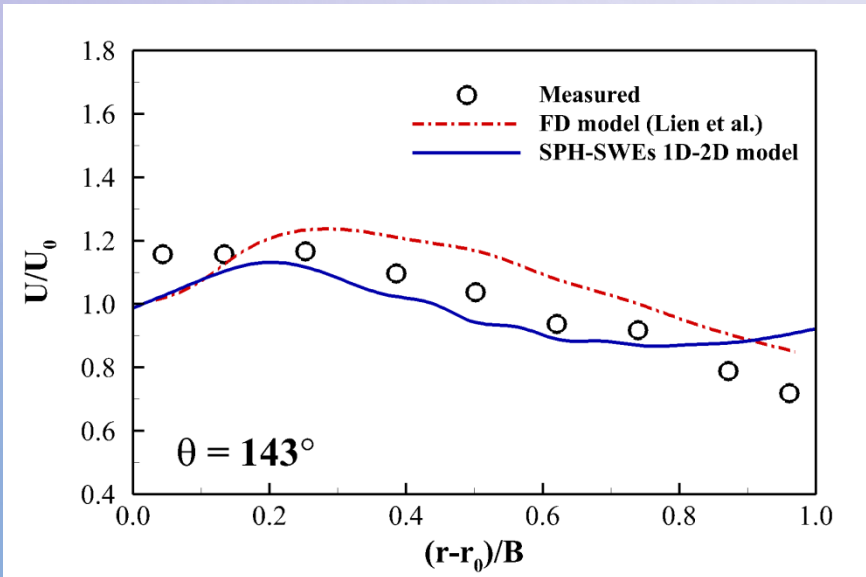
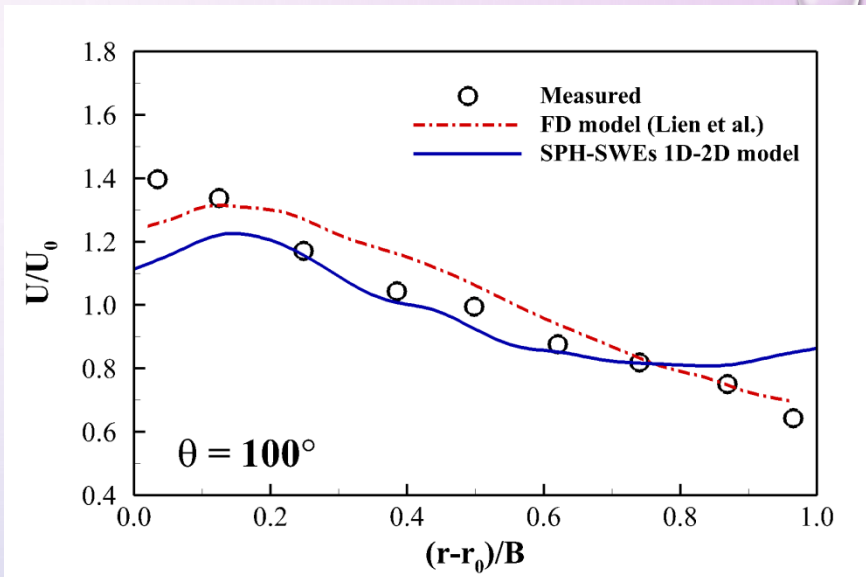
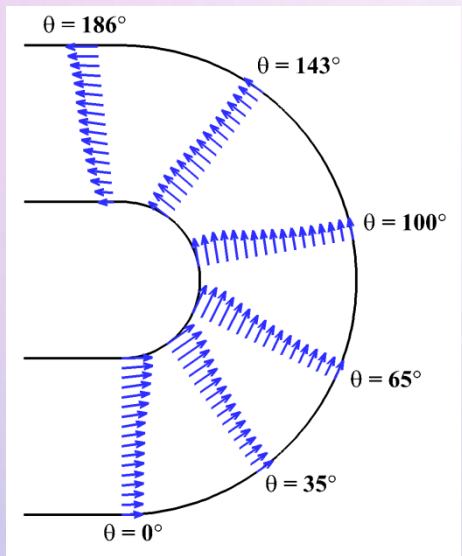




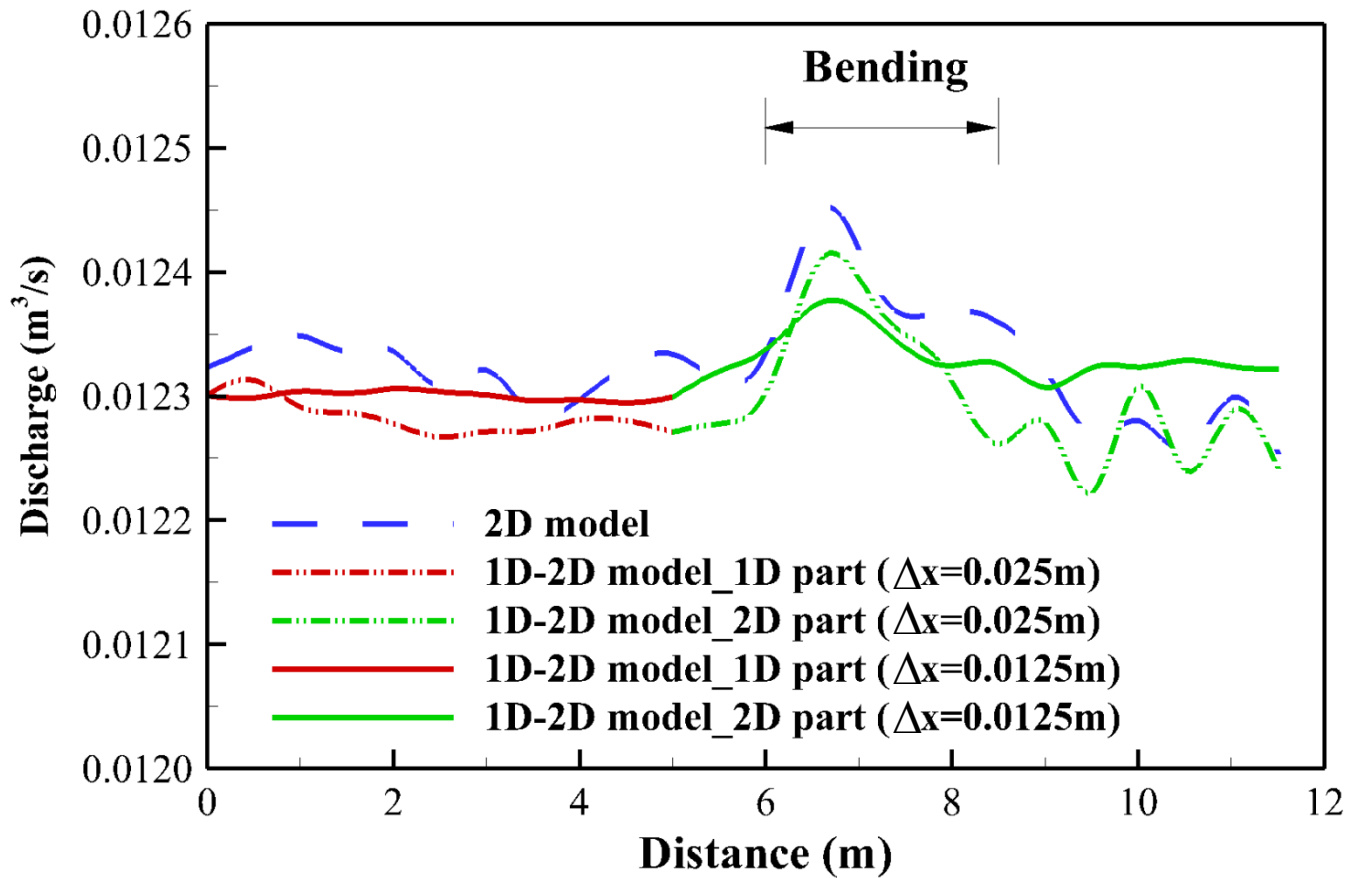
# ➤ Comparison among measured data, FD model and 1D-2D SPH-SWEs model

$\Delta x_0 = 0.0125\text{m}$



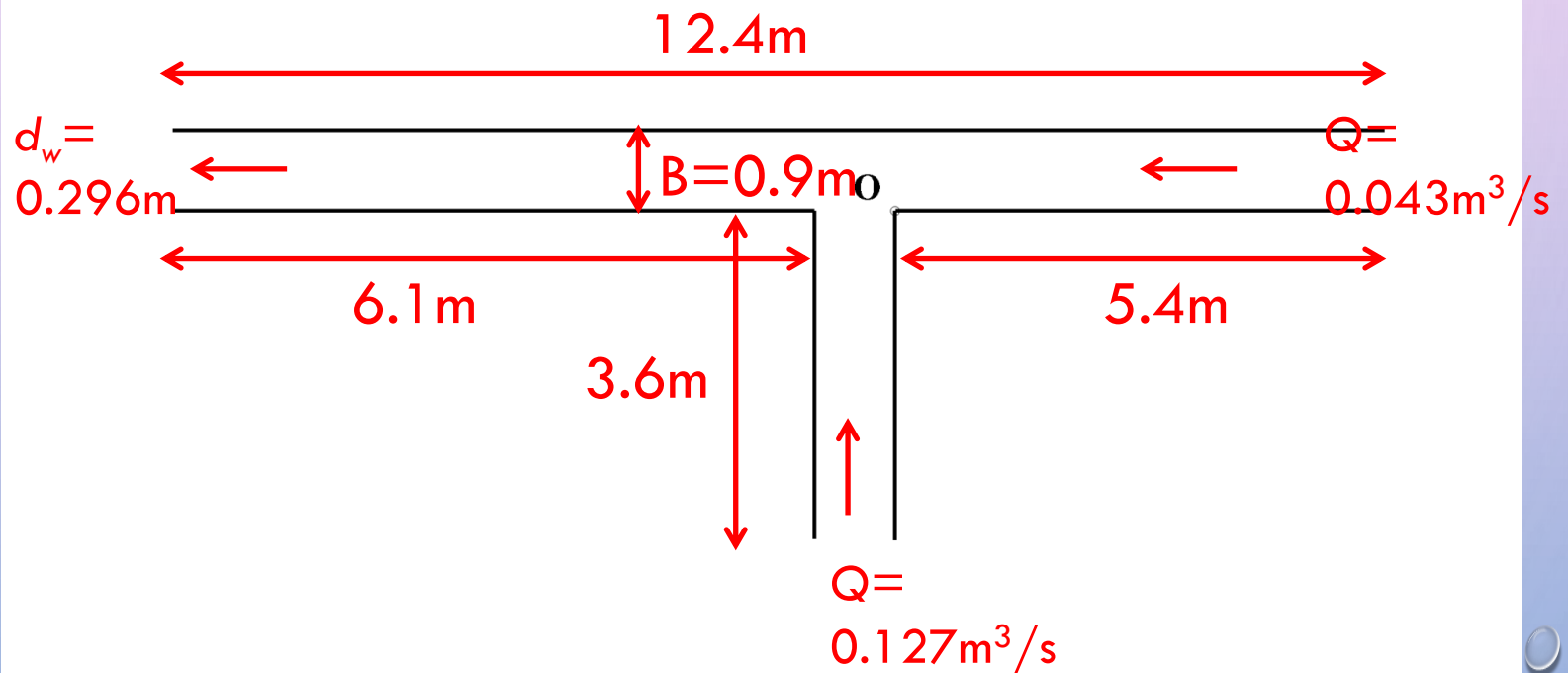


# ➤ Comparison of mass conservation between 2D and 1D-2D SPH-SWEs models



# CASE II: 90 JUNCTION CHANNEL

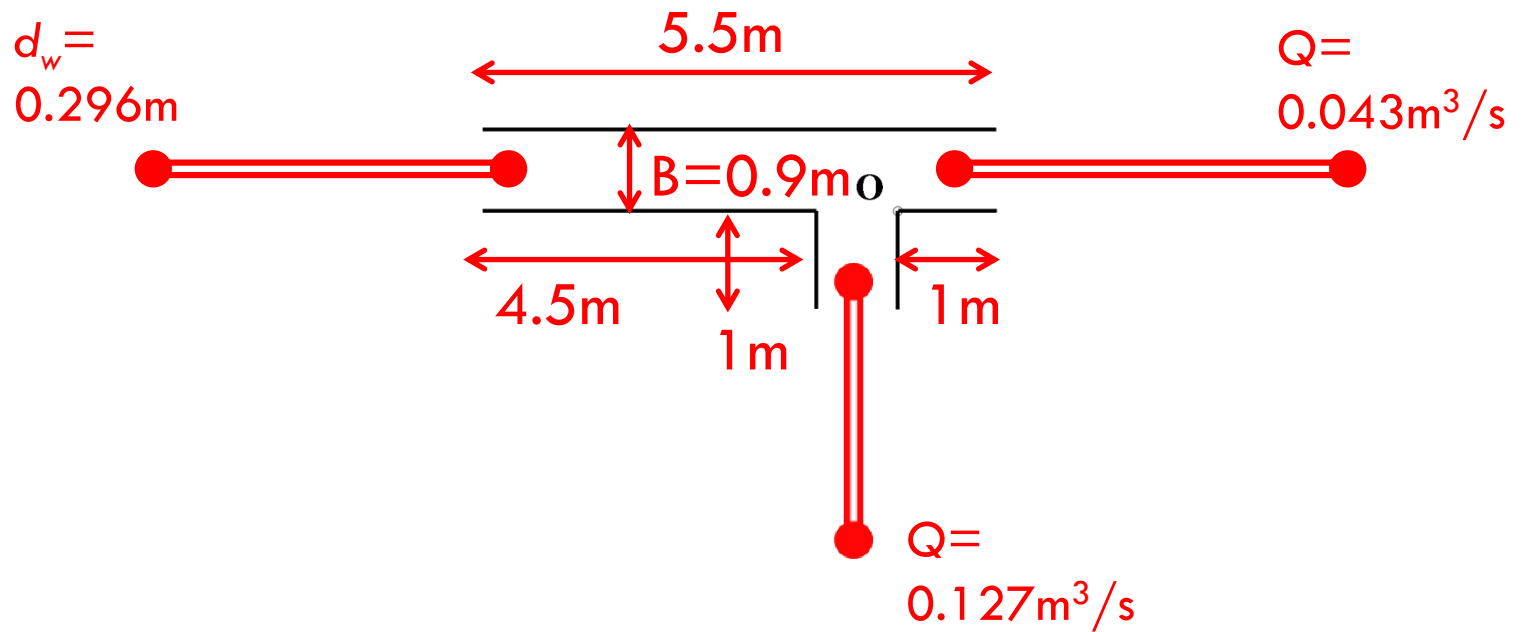
## 2D domain

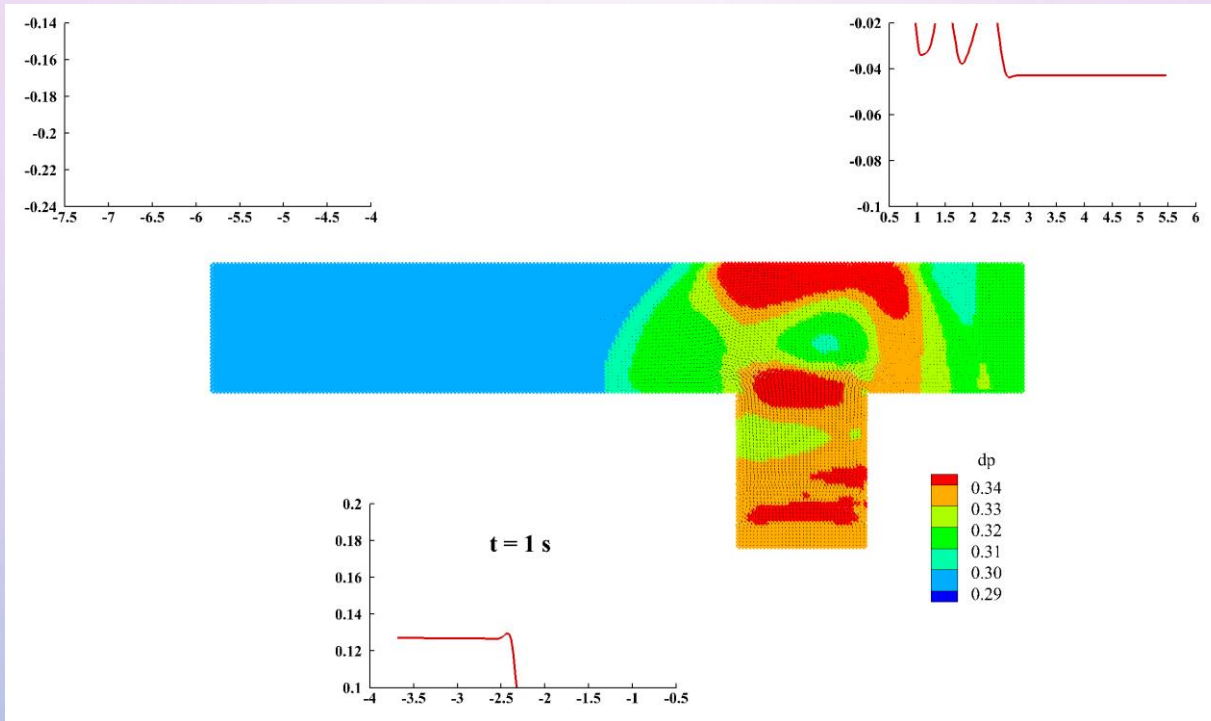




# CASE II: 90 JUNCTION CHANNEL

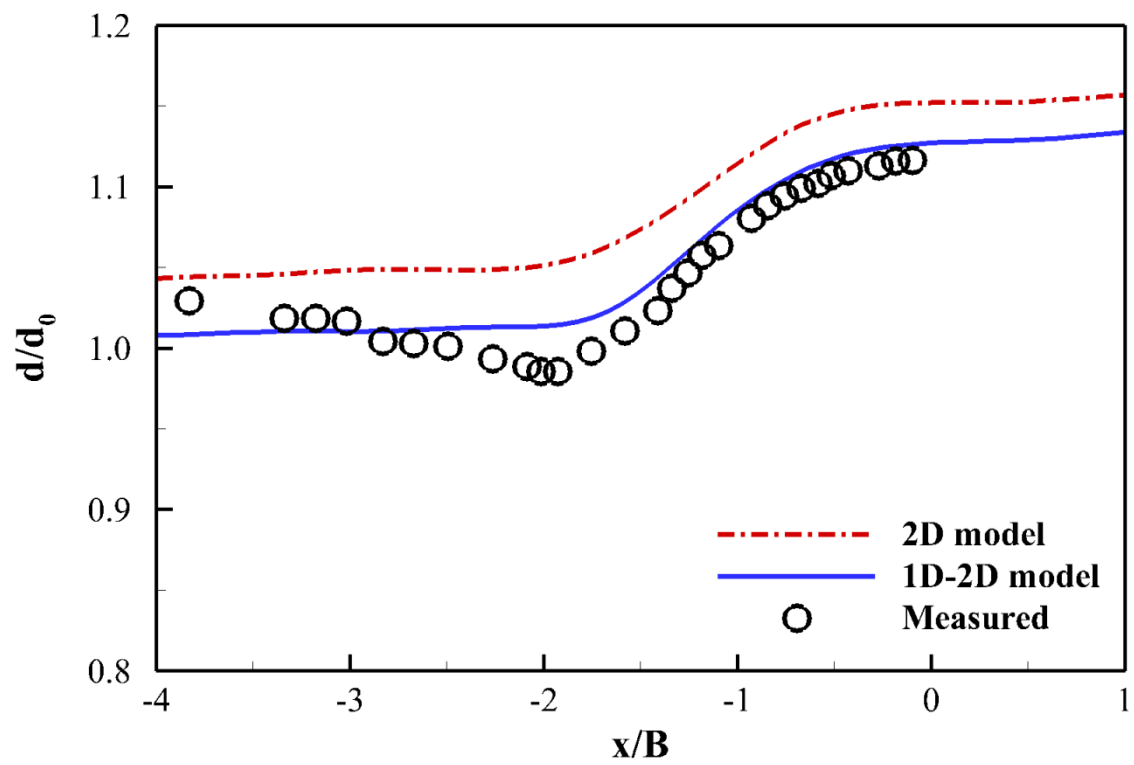
## 1D-2D domain





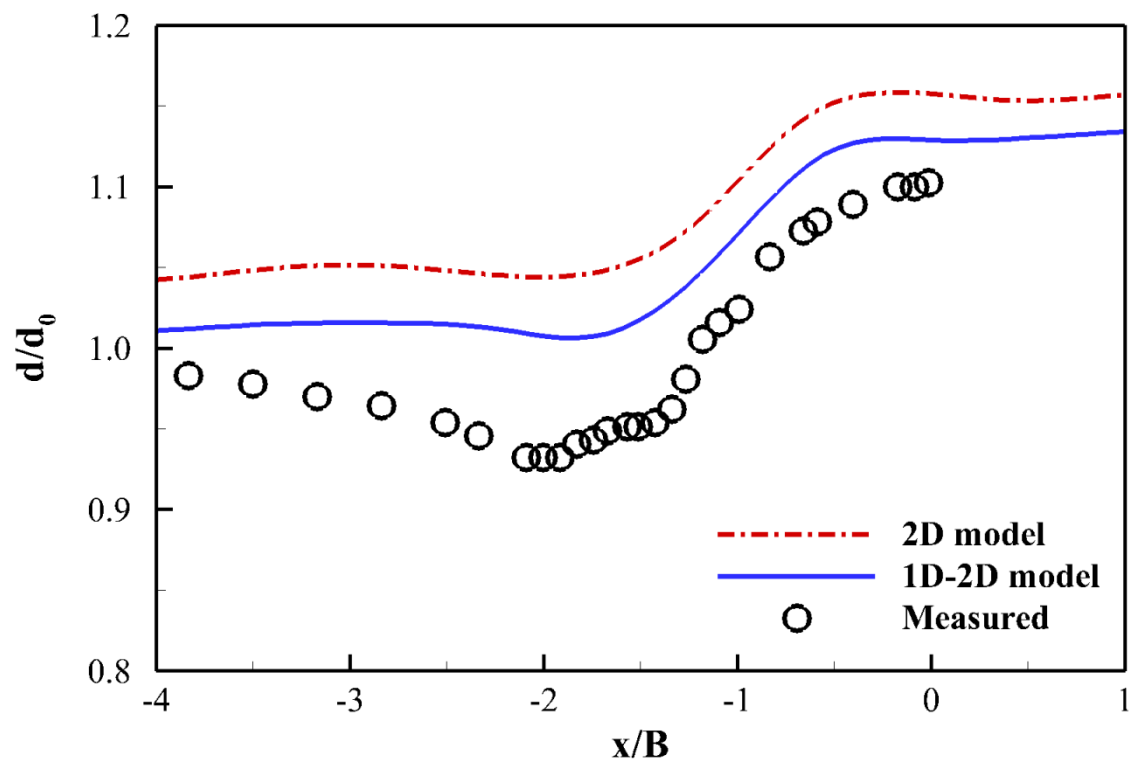
-----  $y/B=0.833$

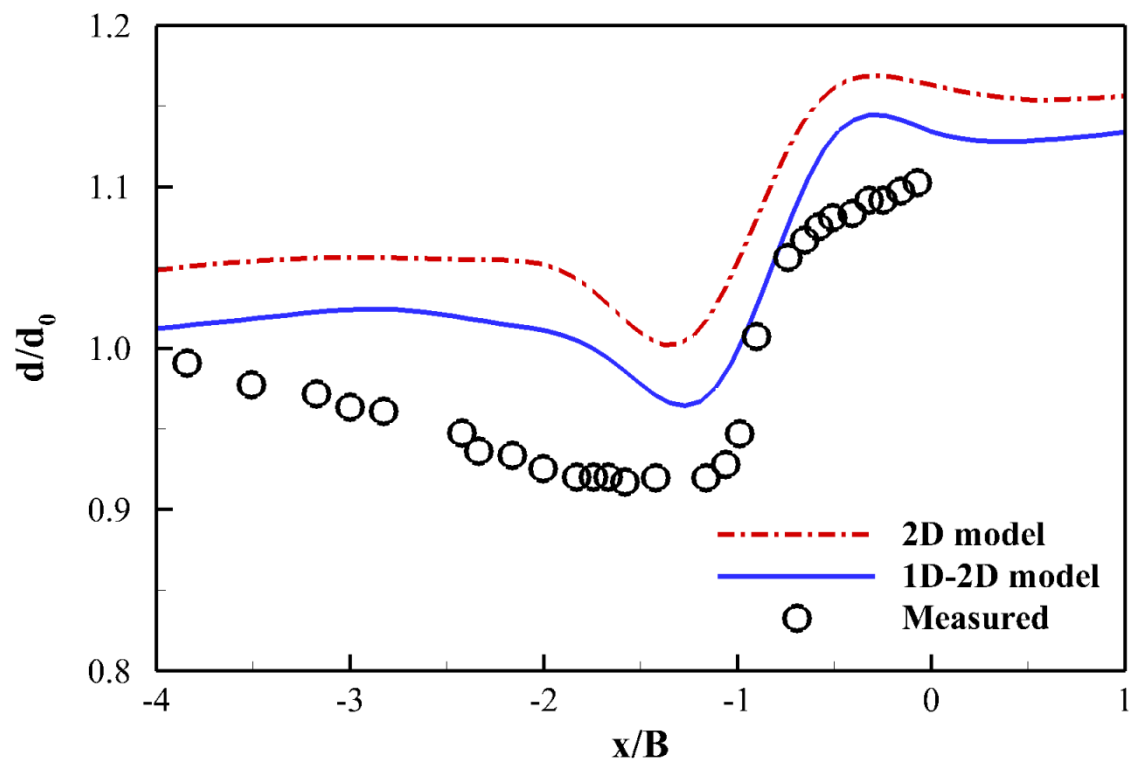
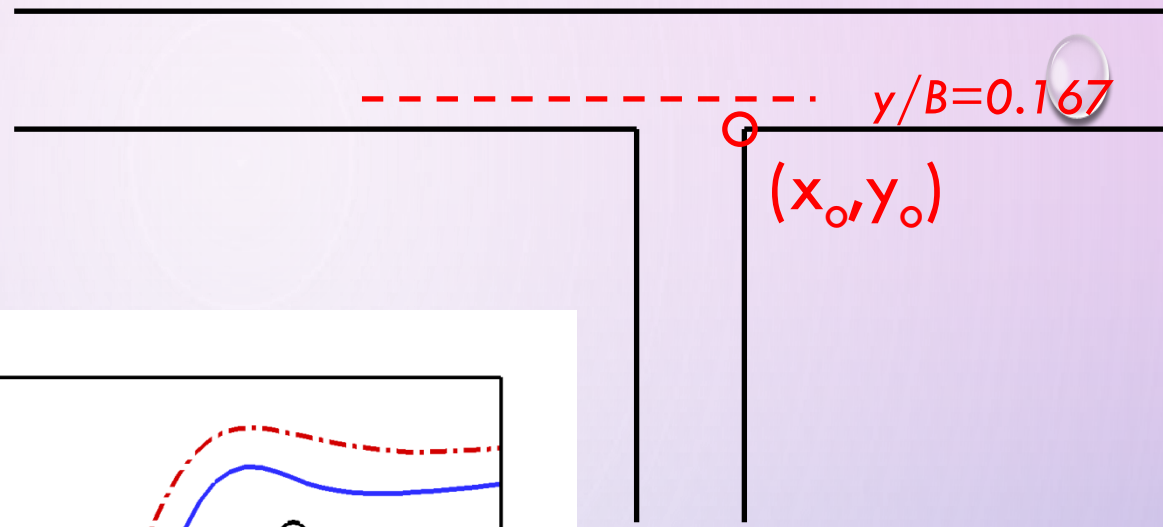
○  
 $(x_o, y_o)$

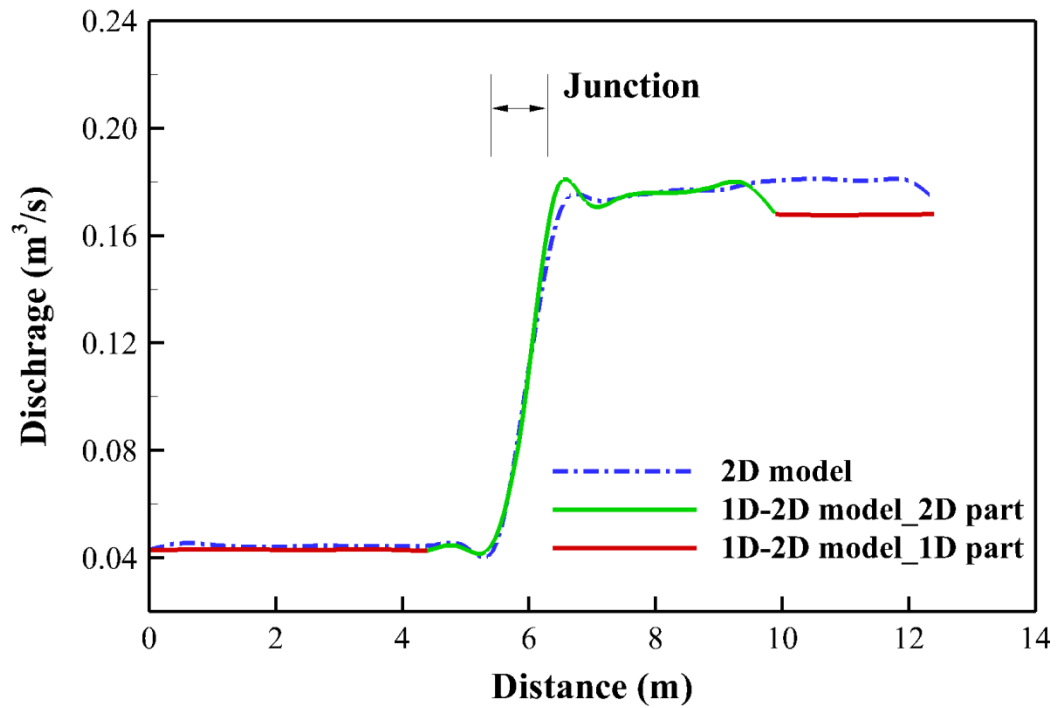


-----  $y/B=0.5$

$(x_o, y_o)$







# **SUMMARY**

### Case I (150s)

	No. of particles	CPU time (sec)	Speed up
2D model	15238	1.13E+05	-
1D-2D model	2D: 8999 1D: 296	6.90E+04	1.63

### Case II (130s)

	No. of particles	CPU time (sec)	Speed up
2D model	23472	6.83E+04	-
1D-2D model	2D: 9792 1D: 416	2.34E+04	2.9



The background features a light purple-to-blue gradient. Numerous realistic water droplets of various sizes are scattered across the surface, some with highlights and shadows. In the upper center, there is a faint, circular logo or watermark that is difficult to discern.

**THANKS FOR LISTENING**