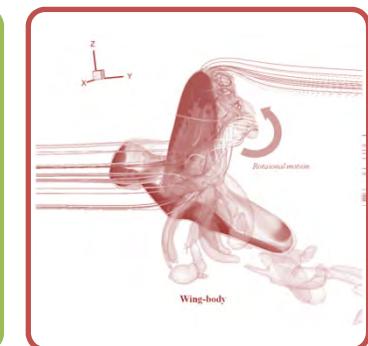
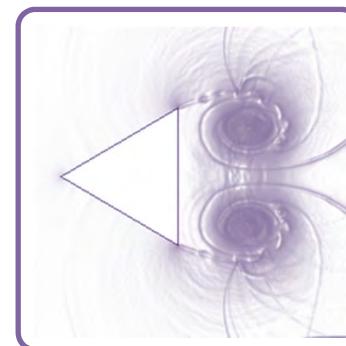
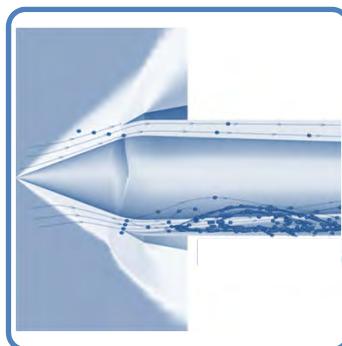


Progress in Limiting Strategy: From One-dimensional Concepts to Multi-dimensional Limiting Process



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11, November, 2013



Computational Fluid Dynamics



- **A Discipline to Solve Conservation Laws of Fluid Dynamics Numerically.**
 - **Flow information**
 - pressure, velocity, temperature...
 - force (lift, drag) & moment or heat flux distribution
 - **Multidisciplinary design optimization and Flow control**
 - performance improvement
 - **Input data**
 - Structural design
 - Engine performance
 - S&C, and so on



Finite Volume Method



- Numerical Discretization
 - FDM, FEM, and FVM
- Finite Volume Discretization
 - Directly apply the integral form of the conservation laws to each computational cell in physical domain.
 - One-step conservative finite volume discretization for 1-D Euler eqns.

$$\mathbf{U}_t + \mathbf{F}(\mathbf{U})_x = \mathbf{0} \text{ with } \mathbf{U} = \begin{bmatrix} \rho \\ \rho u \\ \rho E \end{bmatrix}, \mathbf{F}(\mathbf{U}) = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ \rho u H \end{bmatrix}, \begin{bmatrix} E = e + \frac{u^2}{2} \\ H = h + \frac{u^2}{2} \end{bmatrix}, p = (\gamma - 1)\rho \left[E - \frac{u^2}{2} \right], \gamma = 1.4$$

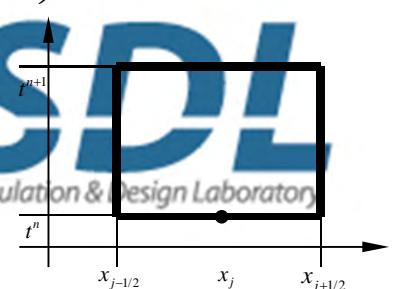
From $\int_{t^n}^{t^{n+1}} \int_{x_{j-1/2}}^{x_{j+1/2}} (\mathbf{U}_t + \mathbf{F}(\mathbf{U})_x) dx dt = \mathbf{0}$, we have

$$\int_{x_{j-1/2}}^{x_{j+1/2}} \mathbf{U}(x, t^{n+1}) dx - \int_{x_{j-1/2}}^{x_{j+1/2}} \mathbf{U}(x, t^n) dx = - \left(\int_{t^n}^{t^{n+1}} \mathbf{F}[\mathbf{U}(x_{j+1/2}, t)] dt - \int_{t^n}^{t^{n+1}} \mathbf{F}[\mathbf{U}(x_{j-1/2}, t)] dt \right)$$

Seeking a solution in an average sense over $\Delta x = [x_{j-1/2}, x_{j+1/2}]$, $\Delta t = [t^n, t^{n+1}]$

by introducing $\mathbf{U}_j^n = \frac{1}{\Delta x} \int_{x_{j-1/2}}^{x_{j+1/2}} \mathbf{U}(x, t^n) dx$, $\hat{\mathbf{F}}_{j+1/2}^n = \frac{1}{\Delta t} \int_{t^n}^{t^{n+1}} \mathbf{F}[\mathbf{U}(x_{j+1/2}, t)] dt$,

We have $\mathbf{U}_j^{n+1} = \mathbf{U}_j^n - \frac{\Delta t}{\Delta x} (\hat{\mathbf{F}}_{j+1/2}^n - \hat{\mathbf{F}}_{j-1/2}^n)$ with $\hat{\mathbf{F}}_{j+1/2}^n = \hat{\mathbf{F}}_{j+1/2}^n(\mathbf{U}_j^n, \mathbf{U}_{j+1}^n)$.





Monotone Scheme

- Some conditions for an entropy-satisfying weak solution

- Consistency

- $\hat{f}(u, u) = f(u)$
- Lipschitz continuity

$$|\hat{f}(u_{j-1}, u_j) - f(u)| \leq L \max(|u_{j-1} - u|, |u_j - u|)$$

- Conservation

- $\hat{f}_{j+1/2,L} = \hat{f}_{j+1/2,R}$

- Monotonicity

- Monotone scheme (Harten-Hyman-Lax, 1987)

$$u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x} (\hat{f}_{j+1/2} - \hat{f}_{j-1/2}) = H_j(u_{j+k}, \dots, u_j, \dots, u_{j-k}),$$

$$\frac{\partial H_j}{\partial u_{j+l}} \geq 0, \quad \forall |l| \leq k.$$

- Consider 3-point scheme

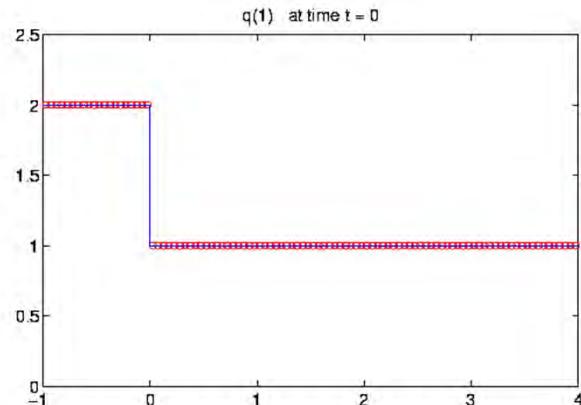
$$u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x} (\hat{f}_{j+1/2} - \hat{f}_{j-1/2}) = H_j(u_{j-1}, u_j, u_{j+1}) \text{ with } \hat{f}_{j+1/2} = \hat{f}_{j+1/2}(u_j, u_{j+1})$$

$\frac{\partial H_j}{\partial u_{j+l}} \geq 0$ gives $\hat{f}_{j+1/2}(u, v)$ is monotonically increasing/decreasing with respect to u / v .

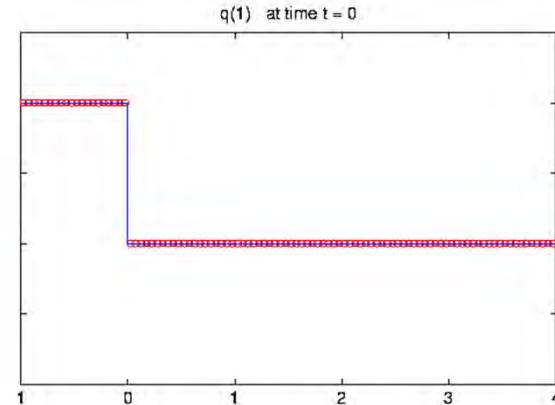
Monotone Scheme

- Computed results

- Conservative and non-conservative scheme

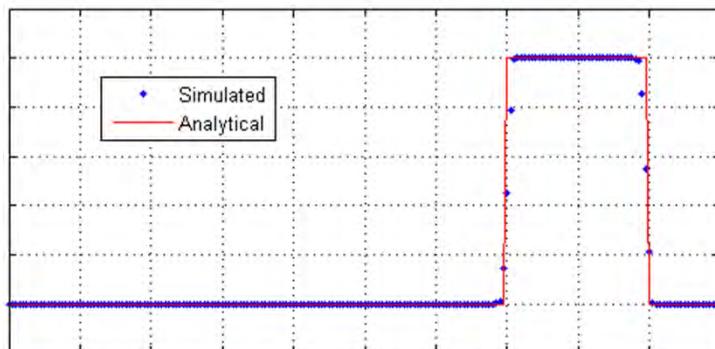


< Conservative scheme >

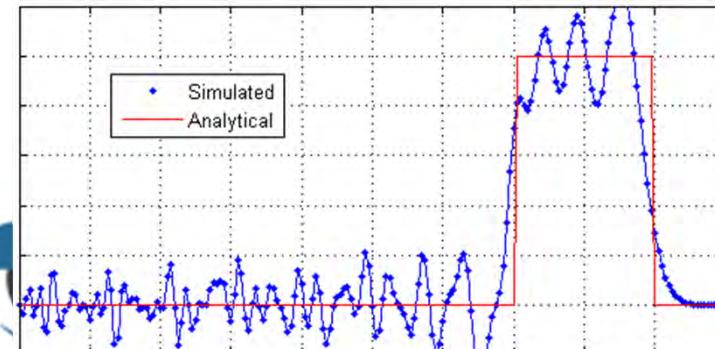


< Non-conservative scheme >

- Monotone and non-monotone scheme



< Monotone scheme >



< Non-monotone scheme >



Non-linear Stability and Hyperbolic PDE

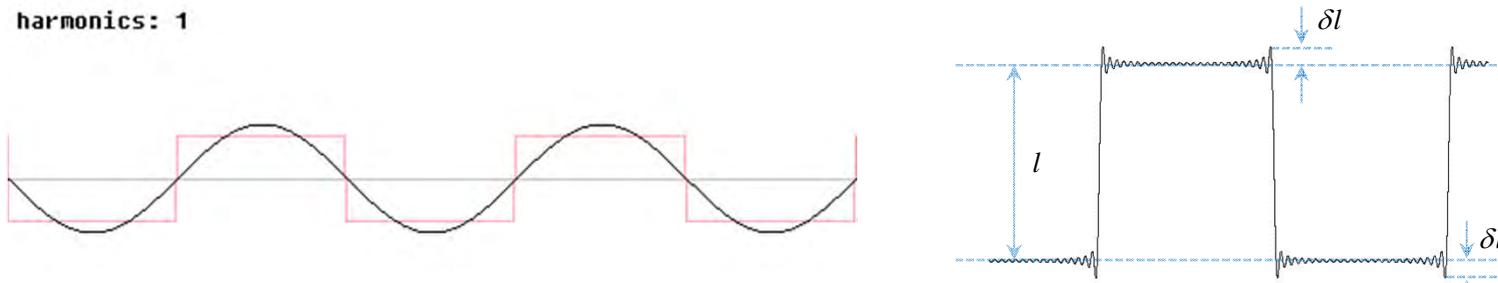
- **Gibb's Phenomenon and Monotonicity**
- **Non-linear Stability**
- **Oscillation Control Strategies**



Gibb's Phenomenon and Monotonicity



- **Gibb's Phenomenon (1899)**
 - Approximation of a profile including discontinuity by Fourier Series (or interpolating techniques based on global basis functions)
 - Numerical oscillations across discontinuity with $O(1)$ → It never dies out even if the number of basis function is increasing.



< Animation of the Gibbs phenomenon >

- Magnitude of overshoot/undershoot $(\delta l / l) \sim \pm 14\%$
- Locally converge (or L_1, L_2 convergence) but not uniformly (L_∞ convergence)
→ warning to naive capturing of discontinuities by increasing the number of interpolating function or mesh point

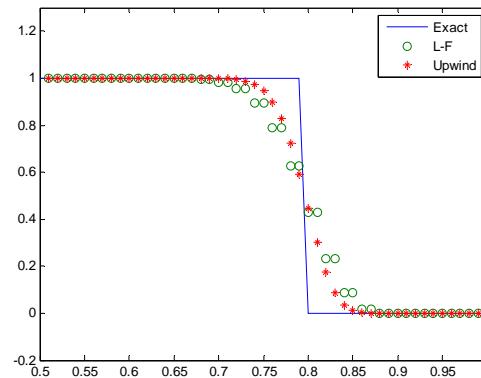


- First-order Scheme and Numerical Diffusion

- For $u_t + au_x = 0$ with Upwind or L-F scheme
- Modified equation

$$u_t + au_x = \begin{cases} \frac{\Delta x^2}{2\Delta t} (1 - \sigma^2) u_{xx} & \text{for L-F} \\ \frac{a\Delta x}{2} (1 - \sigma^2) u_{xx} & \text{for upwind} \end{cases}$$

$$\Rightarrow u_t + au_x = c_1(\Delta x, \Delta t) u_{xx} \text{ with } \sigma = a \frac{\Delta t}{\Delta x}$$



- Leading error term is numerically dissipative → a smooth transition without oscillations

- Excessive numerical dissipation
 - Unacceptable loss of accuracy → Too many grid points
 - Viscous computation and resolution of boundary layer

Gibb's Phenomenon and Monotonicity (Cont')



- Second-order Scheme and Numerical Dispersion

- Modified equation of L-W and B-W scheme

$$u_t + au_x = \begin{cases} -\frac{a\Delta x^2}{6}(1-\sigma^2)u_{xxx} & \text{for L-W} \\ \frac{a\Delta x^2}{6}(2-3\sigma+\sigma^2)u_{xxx} & \text{for B-W} \end{cases}$$

$$\Rightarrow u_t + au_x = c_1 u_{xxx} \xrightarrow[u=\hat{u}(\omega)e^{i\omega x}]{} \hat{u}(\omega, t)_t + i(a\omega + c_1\omega^3)\hat{u}(\omega, t) = 0$$

$$\Rightarrow \hat{u}(\omega, t) \sim e^{-i(a\omega + c_1\omega^3)t} \quad \text{vs. } \hat{u}_{ex}(\omega, t) \sim e^{-ia\omega t}$$

- Numerical dispersion relation: $a(\omega) = a\omega + c_1\omega^3$

For each Fourier component with ω , group velocity $a_g(\omega) \equiv da(\omega)/d\omega$

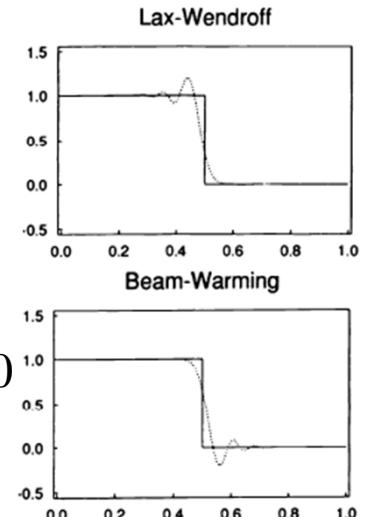
$a_g(\omega) = a + 3c_1\omega^2 \approx a$ for small wave number (long wave length)

$\neq a$ for large wave number (short wave length)

if $a > 0, c_1 < 0$ for L-W \Rightarrow lagging error

$c_1 > 0$ for B-W \Rightarrow leading error

- Numerical oscillations across discontinuity regardless of central differencing or upwinding once the order of accuracy is greater than one.





- **Godunov Barrier Theorem on Monotonicity**

- (Godunov, 1959) *For the fully discretized ‘linear’ scheme of $u_j^{n+1} = Lu_j^n \equiv \sum_q c_{jq} u_{j+q}^n$,*

it can not be better than first-order accurate if the scheme is maximum-norm bounded.

- (Positivity condition) If L is stable in maximum-norm, c_{jq} should be non-negative (or $\frac{du_j^{n+1}}{du_{j+q}^n} \geq 0$).
- The monotonicity condition by HHL can be recovered.

- **Observation**

- **To obtain more than 2nd-order monotone scheme, the scheme should be non-linear even for linear equation.**

- $c_{jq}(\Delta x, \Delta t) \Rightarrow c_{jq}(\Delta x, \Delta t, u_{j+q}^n)$

- **Limiting strategy to realize higher-order accuracy and oscillation-free profile is essential.**

- **Non-linear Stability Criteria**

- **Stability using L_1 norm**

- Total Variation Diminishing (TVD) and Total Variation Bounded (TVB)

- **Stability using L_∞ norm**

- Discrete maximum principle

- Local Extremum Diminishing (LED)



- Total Variation Stability

- A way to connect non-linear stability with the convergence of a computed solution

$$TV(u) \equiv \int_{-\infty}^{\infty} |u'(x)| dx \quad \text{or} \quad TV(u) \equiv \sum_{j=-\infty}^{\infty} |u_j - u_{j-1}| = 2(\sum \text{maxima} - \sum \text{minima})$$

- A useful tool to measure local oscillation
- A monotone scheme is TVD, and a TVD scheme is monotonicity-preserving.
 - 3-point TVD formulation of $TV(u^{n+1}) \leq TV(u^n)$ (Harten, 1983)
- TVB allows oscillation only if it does not grow unboundedly.
 - $TV(u^{n+1}) \leq (1 + \alpha \Delta t) TV(u^n)$
 - TVB may avoid clipping at extrema but badly influences convergence.

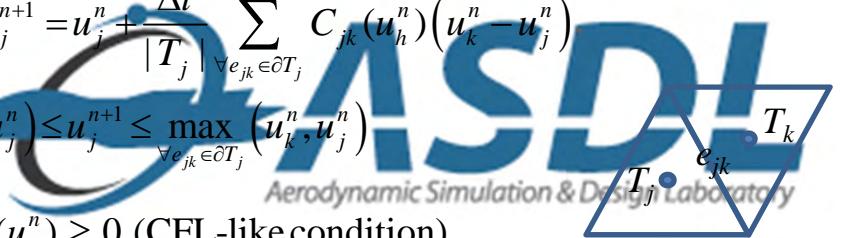
- Discrete Maximum Principle

- Edge-based local space-time discrete maximum principle

- (Barth, 2003) For a fully-discrete scheme of $u_j^{n+1} = u_j^n + \frac{\Delta t}{|T_j|} \sum_{\forall e_{jk} \in \partial T_j} C_{jk}(u_h^n)(u_k^n - u_j^n)$

$$\min_{\forall e_{jk} \in \partial T_j} (u_k^n, u_j^n) \leq u_j^n \leq \max_{\forall e_{jk} \in \partial T_j} (u_k^n, u_j^n) \Rightarrow \min_{\forall e_{jk} \in \partial T_j} (u_k^n, u_j^n) \leq u_j^{n+1} \leq \max_{\forall e_{jk} \in \partial T_j} (u_k^n, u_j^n)$$

If $C_{jk}(u_h^n) \geq 0$ for $\forall e_{jk} \in \partial T_j$, and $1 - \frac{\Delta t}{|T_j|} \sum_{\forall e_{jk} \in \partial T_j} C_{jk}(u_h^n) \geq 0$ (CFL-like condition)





Oscillation Control Strategies



- Remarkable Oscillation Control Strategies since 1970s
 - Flux Corrected Transport (FCT)
 - Boris (1973), Zalesak (1984)
 - MUSCL and Geometric Subcell Reconstruction
 - Van Leer (1977, 1979)
 - Total Variation Diminishing (TVD) / Total Variation Bounded (TVB)
 - Harten (1983), Sweby (1984), Shu (1987)
 - Essentially Non-Oscillatory (ENO) / Weighted ENO (WENO)
 - Harten and Shu et al. (1987), Liu et al. (1994)
 - Spekreijse's Monotonic Concept
 - Spekreijse (1989)
 - Multi-dimensional Reconstruction with Slope Limiter
 - Barth (1989, 1990)
 - Local Extremum Diminishing (LED)
 - A. Jameson (1993)
 - Adaptive Stencil Reconstruction (ENO/WENO)
 - Abgrall (1993), Shu (1999)



One-dimensional Limiting Strategies

- Flux Corrected Transport
- TVD Scheme using Flux Limiter
- MUSCL and Slope Limiter
- ENO/WENO Schemes





- Flux Corrected Transport (FCT) Method and Flux Limiter
 - The first algorithm that recognized the consequence of Godunov's theorem and presented a non-linear way of limiting the cell-interface flux
 - For $u_t + au_x = 0$ with $u_j^{n+1} = u_j^n - \frac{\Delta t}{\Delta x} (F_{j+1/2} - F_{j-1/2})$

Design $F_{j+1/2}$ s.t. $F_{j+1/2} = \begin{cases} \text{2nd-order in smooth region} \\ \text{1st-order across local extrema} \end{cases}$

⇒ Let $\begin{cases} F_{j+1/2}^H : \text{a 2nd-order 'non-monotonic' flux} \\ F_{j+1/2}^L : \text{a 1st-order 'monotonic' flux} \end{cases}$

$$F_{j+1/2} = F_{j+1/2}^L + \alpha_{j+1/2} (F_{j+1/2}^H - F_{j+1/2}^L) \quad \text{with } \alpha_{j+1/2} \begin{cases} \approx 1 & \text{for smooth region} \\ \approx 0 & \text{near local extrema} \end{cases}$$

● Procedure

- S1) Compute $F_{j+1/2}^L$, $F_{j+1/2}^H$ from u_j^n
- S2) Define 'anti-diffusive' flux as $\tilde{F}_{j+1/2} = F_{j+1/2}^H - F_{j+1/2}^L = d_{j+1/2}^L - d_{j+1/2}^H = \varepsilon \Delta u_{j+1/2}^n$ ($\varepsilon = \varepsilon^L - \varepsilon^H > 0$)
- S3) Obtain the intermediate lower-order (or 1st-order) 'monotonic' solution

$$\bar{u}_j = u_j^n - \frac{\Delta t}{\Delta x} (F_{j+1/2}^L - F_{j-1/2}^L)$$

- S4) Correct $\tilde{F}_{j+1/2}$ s.t. the final updated solution (u_j^{n+1}) is free of extrema not found in \bar{u}_j^n or u_j^n

$$F_{j+1/2}^c \equiv \alpha_{j+1/2} \tilde{F}_{j+1/2} = \alpha_{j+1/2} (d_{j+1/2}^L - d_{j+1/2}^H) \quad \text{with } 0 \leq \alpha_{j+1/2} \leq 1$$





Flux Corrected Transport



- S5) Update the final solution with the corrected flux $F_{j+1/2}^c$

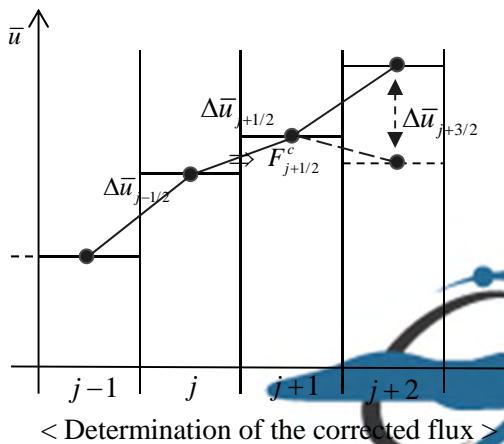
$$u_j^{n+1} = \bar{u}_j - \frac{\Delta t}{\Delta x} (F_{j+1/2}^c - F_{j-1/2}^c) = u_j^n - \frac{\Delta t}{\Delta x} \left\{ \left[F_{j+1/2}^L + \alpha_{j+1/2} (F_{j+1/2}^H - F_{j+1/2}^L) \right]_{j+1/2} - \left[\dots \right]_{j-1/2} \right\}$$

- S6) The corrected flux $F_{j+1/2}^c$ is designed as

$$F_{j+1/2}^c = \min \text{mod} \left(\frac{\Delta x}{\Delta t} \Delta \bar{u}_{j-1/2}, \tilde{F}_{j+1/2}, \frac{\Delta x}{\Delta t} \Delta \bar{u}_{j+3/2} \right)$$

- The anti-diffusive flux is controlled such that it does not create new local extrema.
- Updated soln. satisfies the monotonic constraint in terms of the intermediate distribution

$$\min(\bar{u}_{j-1}, \bar{u}_j, \bar{u}_{j+1}) \leq u_j^{n+1} \leq \max(\bar{u}_{j-1}, \bar{u}_j, \bar{u}_{j+1})$$

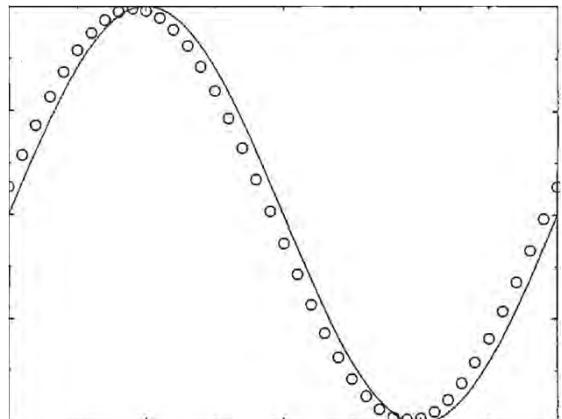




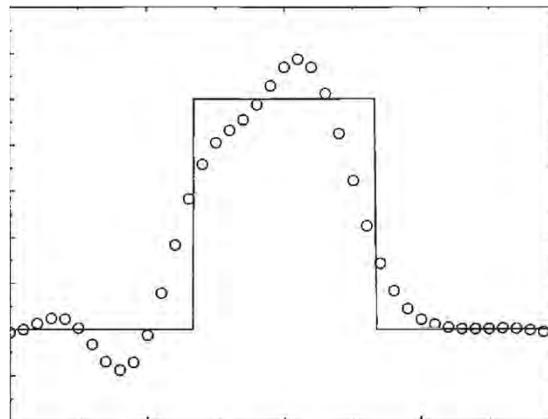
Flux Corrected Transport (Cont')

- Example

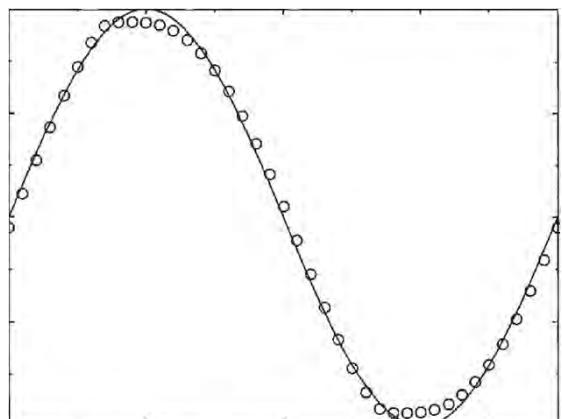
- Linear advection problem with smooth and discontinuous profiles



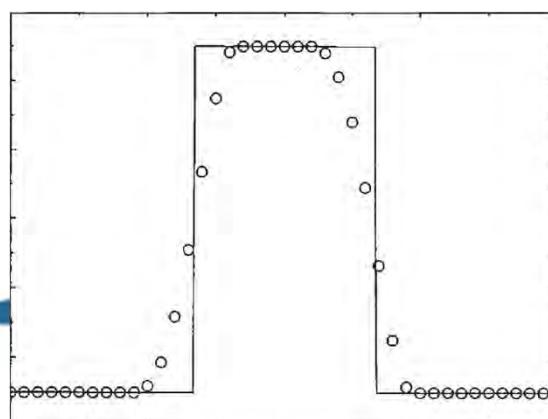
Lax-Wendroff method for test case 1



Lax-Wendroff method for test case 2



FCT method for test case 1



FCT method for test case 2



- TVD Schemes and One-step Flux Limiting Function
 - A Class of one-step scheme using refined form of flux limiters and satisfying TVD property

- The limited flux form is assumed as $F_{j+1/2} = F_{j+1/2}^L + \phi_j (F_{j+1/2}^H - F_{j+1/2}^L)$.
 ϕ_j is a limiter function monitoring the behavior of local solution u_j^n .
 - For $u_t + au_x = 0$ with $a > 0$, consider $F_{j+1/2}^H$ as the L-W flux, and $F_{j+1/2}^L$ as the upwind flux

$$\begin{aligned}
 u_j^{n+1} &= u_j^n - \frac{\sigma}{2} (u_{j+1}^n - u_{j-1}^n) + \frac{\sigma^2}{2} (u_{j+1}^n - 2u_j^n + u_{j-1}^n) \\
 &= u_j^n - \sigma (u_j^n - u_{j-1}^n) - \frac{\sigma(1-\sigma)}{2} (u_{j+1}^n - 2u_j^n + u_{j-1}^n) \\
 &= u_j^n - \frac{\Delta t}{\Delta x} (F_{j+1/2} - F_{j-1/2}) \\
 \Rightarrow F_{j+1/2} &= \underbrace{au_j^n}_{\text{upwind}} + \underbrace{\frac{a(1-\sigma)}{2} \Delta u_{j+1/2}^n}_{\text{Lax-Wendroff correction}}
 \end{aligned}$$

Thus, the limited flux form is $F_{j+1/2} = au_j + \frac{a(1-\sigma)}{2} \Delta u_{j+1/2} \phi_j$ with $\phi_j \geq 0$.

- Define $r_j = \Delta u_{j-1/2} / \Delta u_{j+1/2}$ to measure the change of local slope (or total variation) and design $\phi_j = \phi(r_j)$ with the TVD principle



TVD Scheme using Flux Limiter

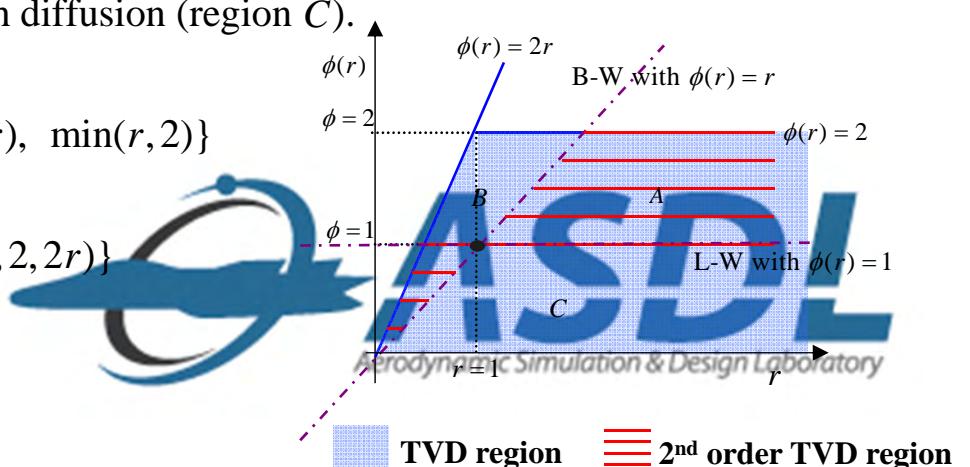
- **Three-point TVD condition**

- The three-point scheme of the form $u_j^{n+1} = u_j^n - C_{j-1/2} \Delta u_{j-1/2}^n + D_{j+1/2} \Delta u_{j+1/2}^n$ is TVD if $C_{j-1/2}, D_{j+1/2} \geq 0$, and $C_{j+1/2} + D_{j+1/2} \leq 1$ for $\forall j$.
- The flux limited form of L-W scheme satisfies the TVD condition is satisfied if

we have $\phi(r)$ s.t. $\begin{cases} 0 \leq \frac{\phi(r)}{r} \leq 2 \text{ and } 0 \leq \phi(r) \leq 2, \text{ if } r \geq 0 \\ \phi(r) = 0, \text{ if } r < 0 \text{ (to prevent the accentuation of local extrema)} \end{cases}$

- **TVD region and limiters**

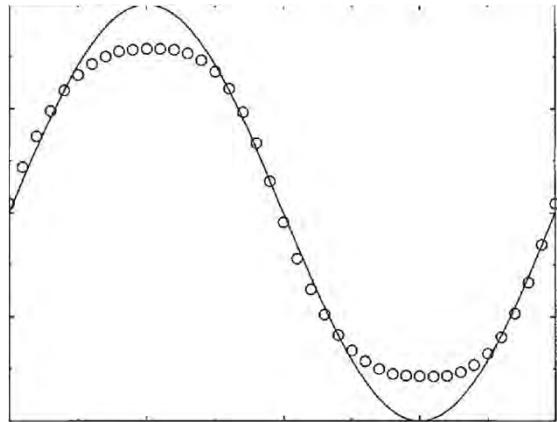
- $\phi(r)_{r=1}=1$ to get a smooth transition with second-order accuracy
- Convex combination of L-W and B-W is desirable to avoid too much compression (region B) or too much diffusion (region C).
- Some TVD limiters
 - superbee limiter: $\phi(r) = \max\{0, \min(1, 2r), \min(r, 2)\}$
 - van Leer limiter: $\phi(r) = (|r|+r)/(1+|r|)$
 - MC limiter: $\phi(r) = \max\{0, \min((1+r)/2, 2, 2r)\}$
 - minmod limiter: $\phi(r) = \min(1, r)$



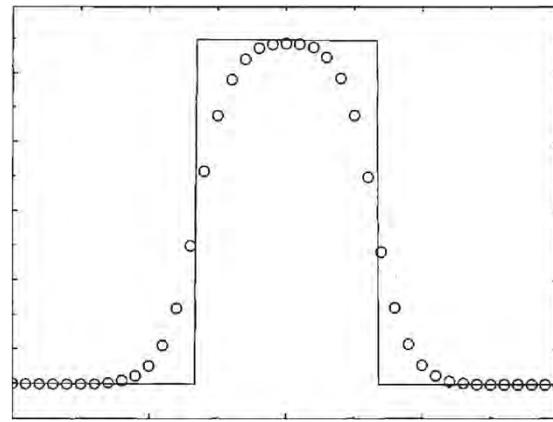


- Example

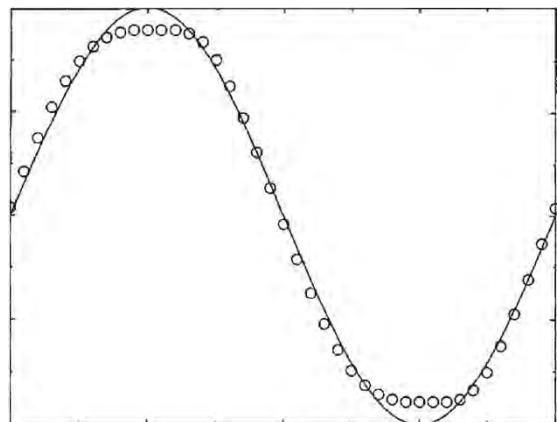
- Linear advection problem with smooth and discontinuous profiles



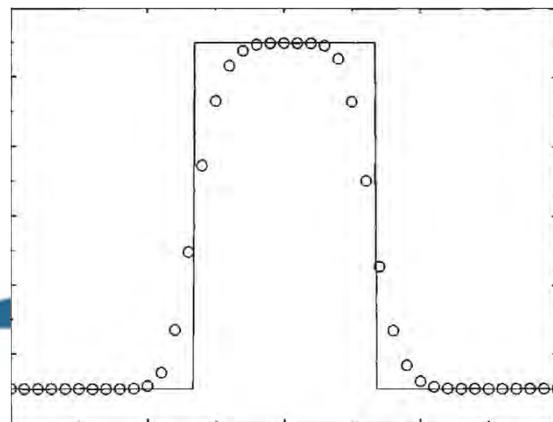
Flux-limited method with minmod limiter



Flux-limited method with minmod limiter



Flux-limited method with van Leer limiter

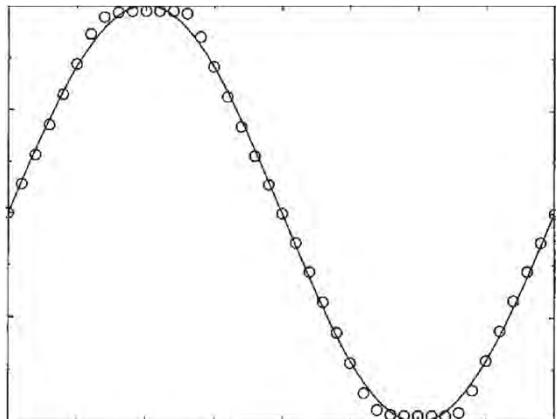


Flux-limited method with van Leer limiter

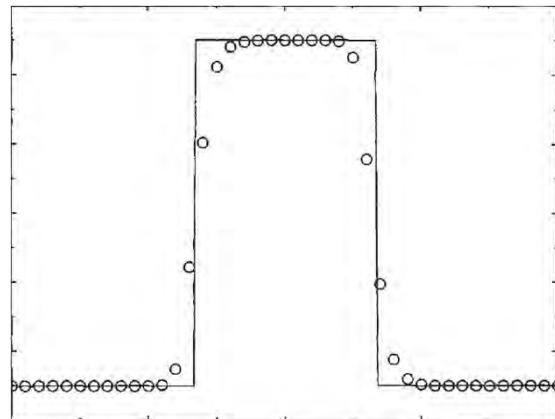


● Example

- Linear advection problem with smooth and discontinuous profiles



Flux-limited method with superbee limiter



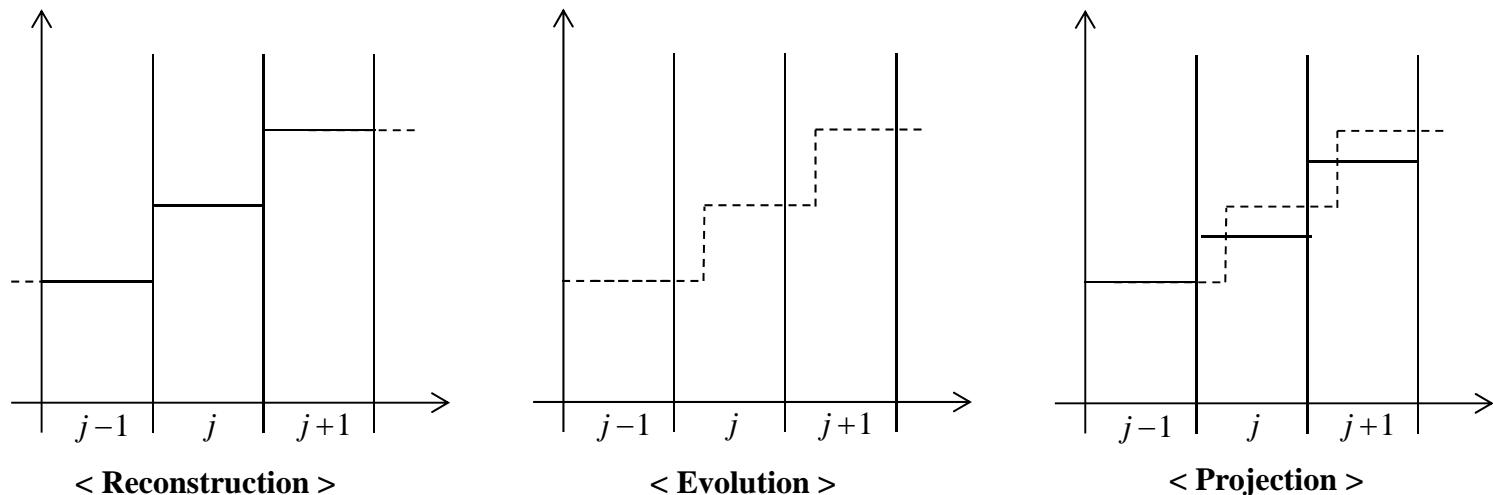
Flux-limited method with superbee limiter

- Accuracy preservation in both smooth and discontinuous region



MUSCL and Slope Limiter

- Geometric Approach for Monotonic Schemes
 - Numerical scheme is analyzed in terms of reconstruction, evolution and projection stage.
 - Ex) Flow physics of first-order upwind scheme for $u_t + au_x = 0, a > 0$



- Reconstruction stage: approximation of exact initial distribution from cell-averaged values
- Evolution stage: reconstructed initial profile is convected $\Delta s_x = a \cdot \Delta t$ by $u_t + au_x = 0$.
- Projection stage: projection of the convected profile to update a new cell-averaged solution at $t^{n+1} = t^n + \Delta t$

MUSCL and Slope Limiter (Cont')



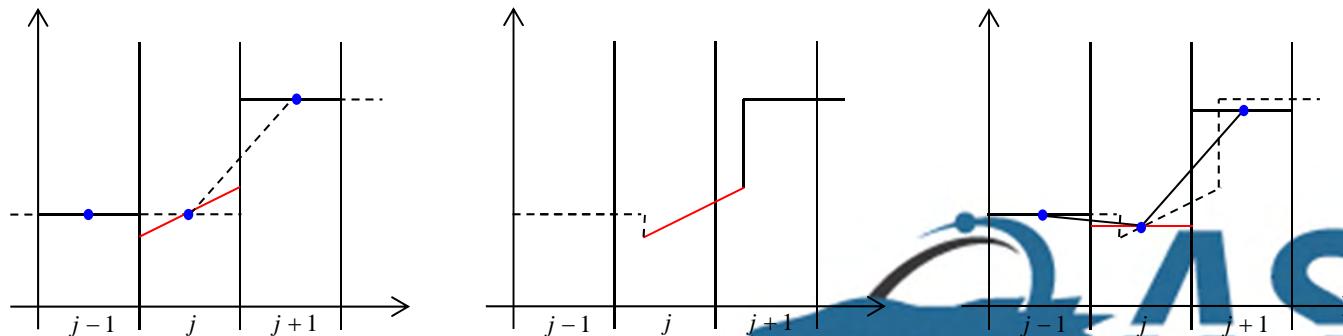
- Evolution and projection stages are not handled numerically.
- A higher-order monotonic scheme via a higher-order reconstruction of the initial data with the monotonic constraint
- Sub-cell distribution for second-order schemes

$$u(x) = u_j + \left(\frac{\partial u}{\partial x} \right)_j (x - x_j) + \left(\frac{\partial^2 u}{\partial x^2} \right)_j \frac{(x - x_j)^2}{2} + O(\Delta x^3)$$

$$\approx u_j + \frac{\Delta u_j}{\Delta x} (x - x_j), \quad x_{j-1/2} \leq x_j \leq x_{j+1/2} \quad \text{with } u_j = \int_{x_{j-1/2}}^{x_{j+1/2}} u(x) dx$$

- Estimation of the local slope Δu_j

Second-order upwind scheme with monotonicity



- Non-monotonic estimation of the local slope (Δu_j) creates a new (artificial) local maxima/minima.

MUSCL and Slope Limiter (Cont')



- If $TV(u(x)) \leq TV(u^n)$ in reconstruction stage, the whole stages are TVD, since evolution and projection stages are intrinsically total variation not increasing.
- Slope limiters to satisfy the monotonic subcell distribution
 - $\Delta u_j|_{\min} = \min(\Delta u_{j+1/2}, \Delta u_{j-1/2})$ • $\Delta u_j|_{sb} = \max(\Delta u_{i1}, \Delta u_{i2})$
 - with $\Delta u_{i1} = \min(\Delta u_{j+1/2}, 2\Delta u_{j-1/2}), \Delta u_{i2} = \min(\Delta u_{j-1/2}, 2\Delta u_{j+1/2})$
 - $\Delta u_j|_{MC} = \min\left(\frac{\Delta u_{j+1/2} + \Delta u_{j-1/2}}{2}, 2\Delta u_{j+1/2}, 2\Delta u_{j-1/2}\right)$

● MUSCL schemes with slope limiters

- First-order upwinding: $\hat{f}_{j+1/2} = \frac{1}{2}(f_j + f_{j+1}) - \frac{|a_{j+1/2}|}{2} \Delta u_{j+1/2}$
- Sub-cell linear distribution using slope limiters: $u(x) = u_j + \frac{\Delta u_j}{\Delta x}(x - x_j), \quad x_{j-1/2} \leq x \leq x_{j+1/2}$
 $\rightarrow u_{jR} = u(x_{j+1/2}) = u_j + \frac{\Delta u_j}{2}$ and $u_{jL} = u(x_{j-1/2}) = u_j - \frac{\Delta u_j}{2}$ to replace u_j and u_{j+1}
- Interface flux using the interpolated values: $\hat{f}_{j+1/2} = \frac{1}{2}(f_L + f_R) - \frac{|a_{j+1/2, L/R}|}{2} \Delta u_{j+1/2, L/R}$



- **Essentially Non-Oscillatory Higher-order Interpolation Procedure**

- **TVD vs. ENO**

- (TVD scheme) Locally first-order accurate across all extrema to strictly enforce monotonicity → accuracy loss across smooth extrema with excessive diffusion or clipping
- (TVD scheme) A fixed stencil (3-point TVD scheme with $u_j, u_{j\pm 1}$)
→ difficult to obtain a higher-order TVD scheme
- (ENO scheme) Allow the increase of local extrema up to the order of truncation error to achieve higher-order accuracy across smooth extrema
- (ENO scheme) Locally adaptive smooth stencil for higher-order interpolation

- **Procedure to construct ENO stencil (or locally adaptive stencil)**

- Cell-averaged value or flux function

- **ENO polynomial reconstruction using cell-averaged value, u_i**

- Step 1. Introduce Newton's divided difference based on cell-averaged values

$$u[x_i] \equiv u_i, u[x_i, x_{i+1}] \equiv \{u[x_{i+1}] - u[x_i]\} / \Delta x = (u_{i+1} - u_i) / \Delta x, \dots,$$

$$u[x_i, x_{i+1}, \dots, x_{i+k}] \equiv \{u[x_{i+1}, x_{i+2}, \dots, x_{i+k}] - u[x_i, x_{i+1}, \dots, x_{i+k-1}]\} / (k \Delta x)$$

If flow field is sufficiently smooth, $u[x_i, x_{i+1}, \dots, x_{i+k}] = \frac{1}{k!} \left. \frac{d^k u}{dx^k} \right|_{\xi} \quad x_i \leq \xi \leq x_{i+k}$

ENO/WENO Schemes (Cont')



- Step 1. Introduce Newton's divided difference based on cell-averaged values

And, if $u(x)$ has a jump discontinuity of the l -th derivative at $x = x_p$,

$$u[x_i, x_{i+1}, \dots, x_{i+k}] = O\left(\frac{1}{\Delta x^{k-l}} \left(\frac{d^l u(x_p^+)}{dx^l} - \frac{d^l u(x_p^-)}{dx^l} \right)\right), \quad x_i \leq \xi \leq x_{i+k}$$

- Step 2. Compare each divided difference to determine ENO stencil successively

Starting from the cell i , take $\begin{cases} i+1 & \text{if } |u[x_i, x_{i+1}]| \leq |u[x_{i-1}, x_i]| \\ i-1 & \text{if } |u[x_i, x_{i+1}]| > |u[x_{i-1}, x_i]| \end{cases}$

$$\Rightarrow u(x) = u_i + u[x_j, x_{j+1}](x - x_i), \quad x_{j-1/2} \leq x \leq x_{j+1/2} \quad (\text{linear reconstruction})$$

$$\text{with } j = \begin{cases} i & \text{if } |u[x_i, x_{i+1}]| \leq |u[x_{i-1}, x_i]| \\ i-1 & \text{if } |u[x_i, x_{i+1}]| > |u[x_{i-1}, x_i]| \end{cases}$$

- Step 3. ENO stencil in a recursive manner

For $m = 0, 1, 2, \dots, n-1$ with $l_0(i) = i$,

$$l_{m+1}(i) = \begin{cases} l_m(i) & \text{if } |u[x_{l_m(i)}, x_{l_m(i)+1}, \dots, x_{l_m(i)+m+1}]| \leq |u[x_{l_m(i)-1}, x_{l_m(i)}, \dots, x_{l_m(i)+m}]| \\ l_m(i)-1 & \text{if } |u[x_{l_m(i)}, x_{l_m(i)+1}, \dots, x_{l_m(i)+m+1}]| > |u[x_{l_m(i)-1}, x_{l_m(i)}, \dots, x_{l_m(i)+m}]| \end{cases}$$



- Step 3. ENO stencil in a recursive manner

From $[x_{l_m(i)}, x_{l_m(i)+1}, \dots, x_{l_m(i)+m}]$ with $0 \leq m \leq n$, construct a n -th order ENO polynomial as

$$u_n(x) = u[x_{l_n(i)}] + u[x_{l_n(i)}, x_{l_n(i)+1}] (x - x_{l_n(i)}) + u[x_{l_n(i)}, x_{l_n(i)+1}, x_{l_n(i)+2}] (x - x_{l_n(i)}) (x - x_{l_n(i)+1}) + \dots \\ + u[x_{l_n(i)}, x_{l_n(i)+1}, \dots, x_{l_n(i)+n}] (x - x_{l_n(i)}) (x - x_{l_n(i)+1}) \dots (x - x_{l_n(i)+n-1})$$

$$\Rightarrow u(x) = \sum_{j=0}^n u[x_{l_n(i)}, x_{l_n(i)+1}, \dots, x_{l_n(i)+j}] \prod_{k=0}^{j-1} (x - x_{l_n(i)+k}), \quad x_{j-1/2} \leq x \leq x_{j+1/2}$$

- With the conservation constraint of $\int_{x_{i-1/2}}^{x_{i+1/2}} u(x) dx = u_i$, it can be shown that

$$TV(u^{n+1}) \leq TV(u^n) + O(\Delta x^r) \text{ for } r\text{-th order ENO interpolation,}$$

if there are at least $(r+1)$ smooth points between local smooth extrema.

ENO interpolation using flux function, f_i

- Step 1. Start from first-order flux (mostly, Lax-Friedrich flux)

$$f_i = f_i^+ + f_i^- \text{ with } f_i^+ \equiv \frac{1}{2}(f_i + \alpha u_i), f_i^- \equiv \frac{1}{2}(f_i - \alpha u_i), \alpha \geq \max_{u \in [u_i, u_{i+1}]} |\partial f / \partial u|$$

the cell-interface flux is obtained by $\hat{f}_{i+1/2} = f_i^+ + f_{i+1}^- = [(f_i + f_{i+1}) - \alpha \Delta u_{i+1/2}] / 2$.

- Step 2. Carry out the same piecewise polynomial reconstruction using f_i^\pm

$$u[x_i, x_{i+1}, \dots, x_{i+k}] \rightarrow f^\pm[x_i, x_{i+1}, \dots, x_{i+k}] \Rightarrow \hat{f}_{i+1/2} = (f_{iL}^+ + f_{(i+1)R}^-) / 2$$



- Characteristics of ENO scheme

- Pure interpolation not limiting technique → accuracy preservation across smooth extrema
- Adaptive smoothest stencil of ENO scheme could be changed by a small perturbation at round-off level → convergence and accuracy problem

- Weighted ENO (WENO) Scheme

- Smoothness indicator is introduced for non-linear weighting coefficients.
- WENO polynomial reconstruction using cell-averaged value, u_i

- Step 1. Interpolation on global/local stencil

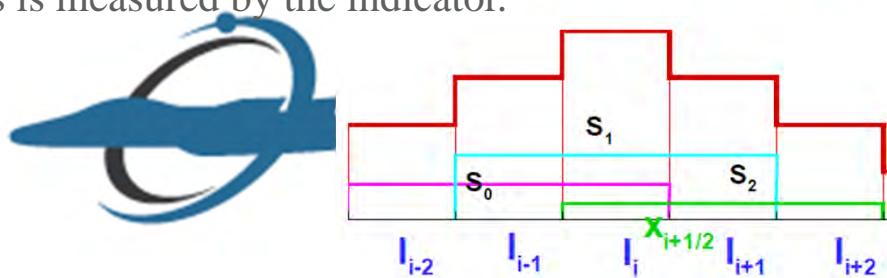
Reconstruct $k - th$ degree polynomial $p_j(x)$ on S_j and $2k - th$ degree polynomial $Q(x)$ on $T = \sum_j S_j$ as

$$\bar{u}_{i+l} = \int_{I_{i+l}} p_j(x) dx / \Delta x_{i+l}, \quad (l = -k + j, \dots, j) \text{ and } \bar{u}_{i+l} = \int_{I_{i+l}} Q(x) dx / \Delta x_{i+l}, \quad (l = -k, \dots, k)$$

Find linear weights such that $Q(x_{i+1/2}) = \sum_{j=0}^k \gamma_j p_j(x_{i+1/2})$

- Step 2. For each cell S_j , local smoothness is measured by the indicator.

$$\beta_j = \sum_{l=1}^k \int_{x_{i-1/2}}^{x_{i+1/2}} \Delta x^{2l-1} \left(\frac{\partial^l}{\partial x^l} p_j(x) \right)^2 dx$$



ENO/WENO Scheme (Cont')



- Step 3. Non-linear weights to satisfy the ENO property
 - If the stencil S_j is in smooth region: $\omega_j = O(1)$
 - If the stencil S_j is in non-smooth region: $\omega_j \leq O(h^k)$
- Compute the nonlinear weights based on the smoothness indicator

$$\bar{\omega}_j = \frac{\gamma_j}{(\varepsilon + \beta_j)^2} \quad \text{and} \quad \omega_j = \frac{\bar{\omega}_j}{\sum_l \bar{\omega}_l}$$

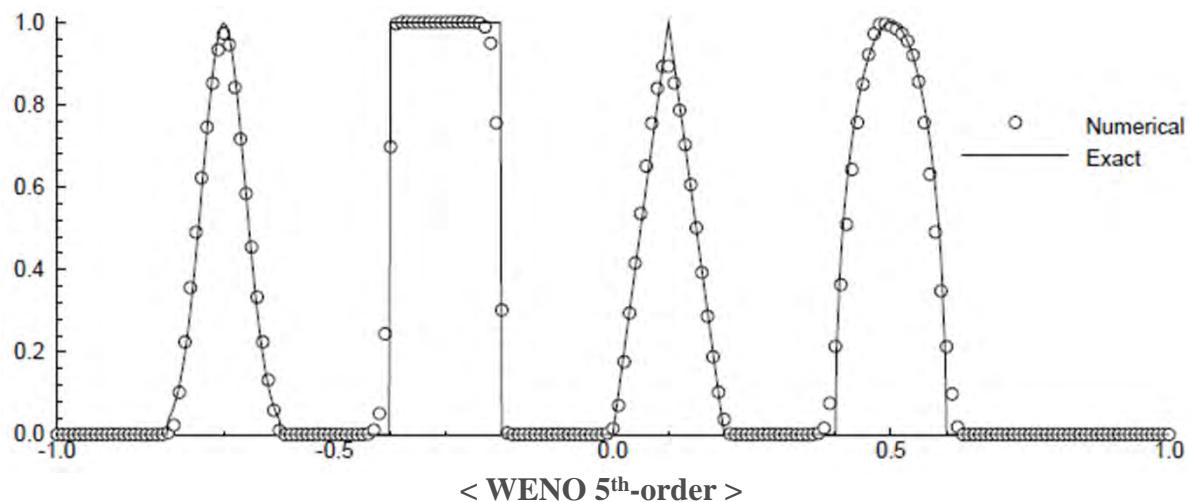
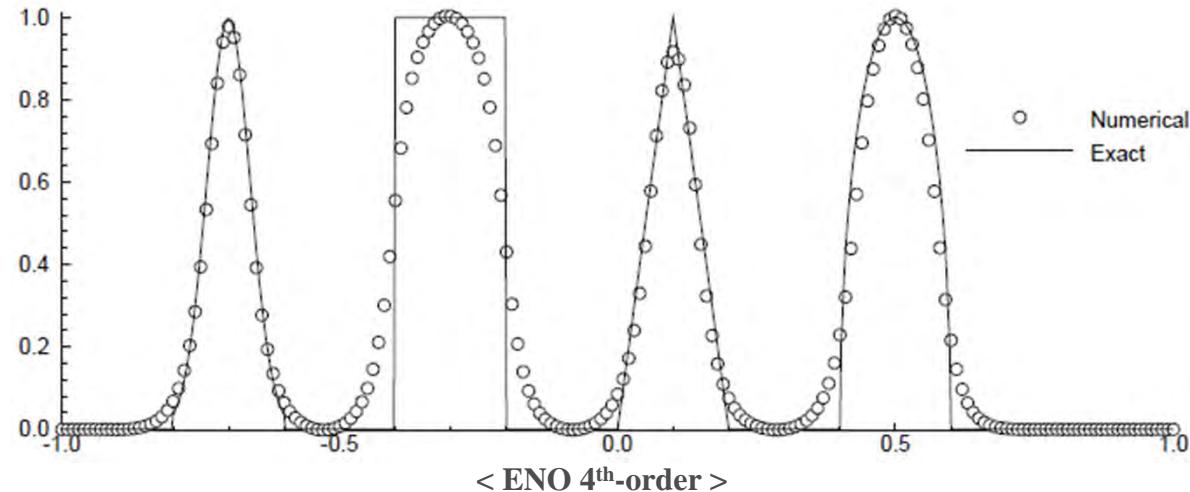
- Step 4. $(2k+1)$ -th degree WENO reconstruction via convex combination with the non-linear weights

$$R_i(x) = \sum_{j=0}^k \omega_j p_j(x) \Rightarrow u_{iR} = R_i(x_{i+1/2}), \quad u_{iL} = R_i(x_{i-1/2})$$





- Example





Multi-dimensional Limiting Process

- Problems in Multi-dimensional Extension
- Limiting Condition on Multiple Dimensions
- Formulation of MLP on Structured Grids
- Maximum Principle and MLP Condition
- Numerical Results



Problems in Multi-dimensional Extension



- Critical Survey

- Analyses based on one-dimensional flow physics

- 1-D SCL of $\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} = 0$ with $f(u) = au, \frac{u^2}{2}$ $\Rightarrow \frac{d\bar{u}_j}{dt} = -\frac{1}{\Delta x}(\hat{f}_{j+1/2} - \hat{f}_{j-1/2})$
- Monotonicity constraint on $u_{j+1/2}$: $\min(\bar{u}_j, \bar{u}_{j+1}) \leq u_{j+1/2} \leq \max(\bar{u}_j, \bar{u}_{j+1})$

- Multi-dimensional extension

- 2-D SCL of $\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} + \frac{\partial g(u)}{\partial y} = 0 \Rightarrow \frac{d\bar{u}_{i,j}}{dt} = -\frac{1}{\Delta x}(\hat{f}_{i+1/2,j} - \hat{f}_{i-1/2,j}) - \frac{1}{\Delta y}(\hat{g}_{i,j+1/2} - \hat{g}_{i,j-1/2})$
- Monotonicity constraint by dimensional splitting:

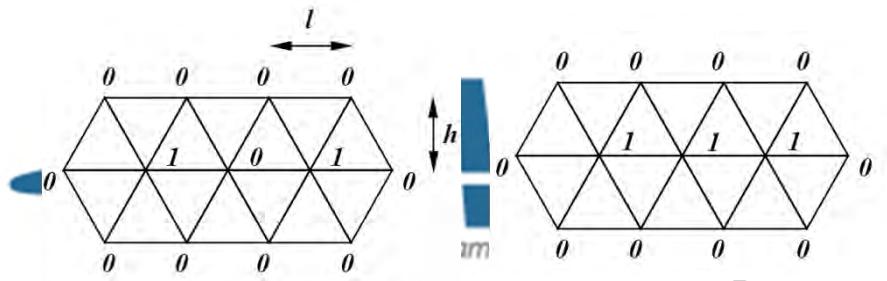
$$\left. \begin{array}{l} (x\text{-dir}) \quad \min(\bar{u}_{i,j}, \bar{u}_{i+1,j}) \leq u_{i+1/2,j} \leq \max(\bar{u}_{i,j}, \bar{u}_{i+1,j}) \\ (y\text{-dir}) \quad \min(\bar{u}_{i,j}, \bar{u}_{i,j+1}) \leq u_{i,j+1/2} \leq \max(\bar{u}_{i,j}, \bar{u}_{i,j+1}) \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} \text{Is it good enough to handle multi-} \\ \text{dimensional flow situation?} \end{array} \right.$$

- Two-peaks and one-ridge problem (Jameson, 1995)

- For $TV(\bar{u}) = \int_V \|\nabla \bar{u}\|_p ds$ with $p=1, 2, \infty$

$$TV_{\text{two-peaks}} < TV_{\text{one-ridge}}$$

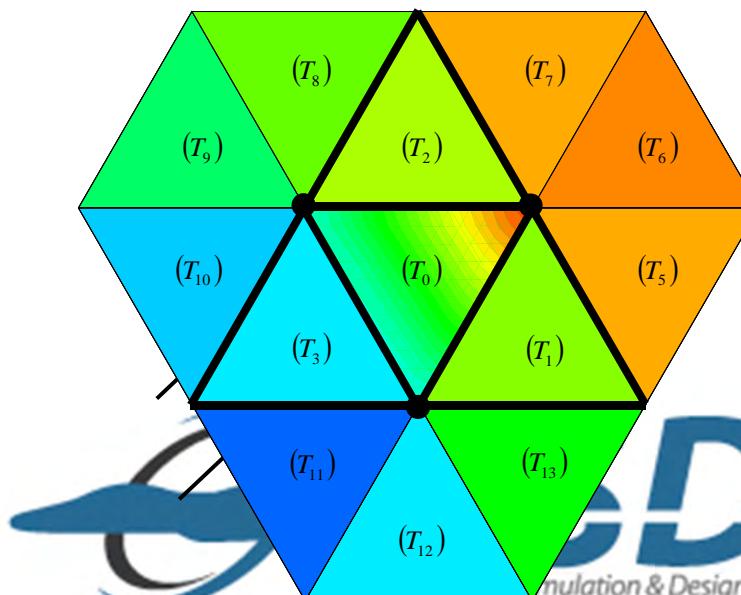
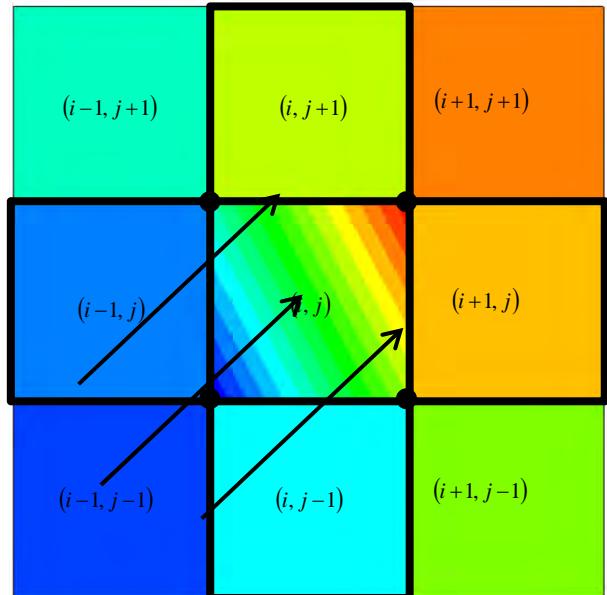
- 2-D TVD scheme is at most first-order accurate (Goodman and LeVeque, 1985)



Problems in Multi-dimensional Extension (Cont')



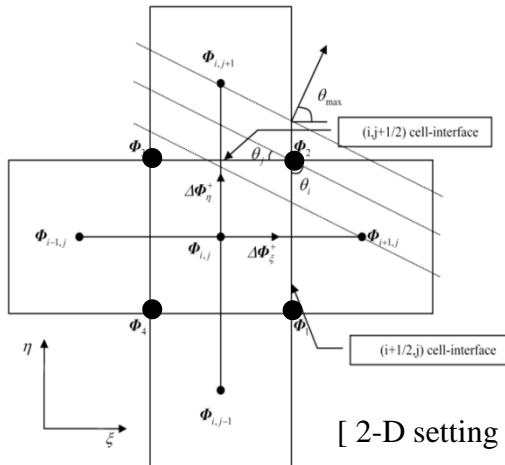
- Dimensional splitting extension is not sufficient to control numerical oscillations in multiple dimensions.
 - Dimensional splitting TVD schemes do not guarantee monotonicity at vertex (Kim & Kim, 2005 / Yoon & Kim, 2008)
- Limiting strategy in Non Grid-aligned Distributions
 - Dimensional splitting approach cannot handle local extrema at vertex



Limiting Condition on Multiple Dimensions



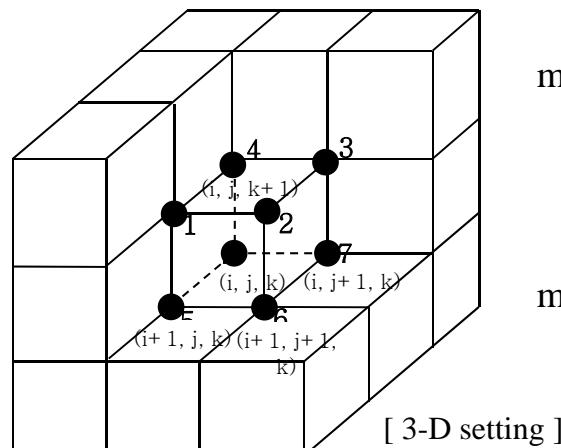
- **Multiple Dimensional Limiting Condition:** $\bar{u}_{neighbor}^{\min} \leq u \leq \bar{u}_{neighbor}^{\max}$
- **MLP Condition on 2-D and 3-D Structured Grids**



$$\min(\bar{u}_{i,j}, \bar{u}_{i+1,j}, \bar{u}_{i,j-1}, \bar{u}_{i+1,j-1}) \leq u_1 \leq \max(\bar{u}_{i,j}, \bar{u}_{i+1,j}, \bar{u}_{i,j-1}, \bar{u}_{i+1,j-1})$$

...

$$\min(\bar{u}_{i,j}, \bar{u}_{i-1,j}, \bar{u}_{i,j+1}, \bar{u}_{i-1,j+1}) \leq u_4 \leq \max(\bar{u}_{i,j}, \bar{u}_{i-1,j}, \bar{u}_{i,j+1}, \bar{u}_{i-1,j+1})$$



$$\min\left(\bar{u}_{i,j,k}, \bar{u}_{i+1,j,k}, \bar{u}_{i,j-1,k}, \bar{u}_{i+1,j-1,k}, \bar{u}_{i,j,k+1}, \bar{u}_{i+1,j,k+1}, \bar{u}_{i,j-1,k+1}, \bar{u}_{i+1,j-1,k+1}\right) \leq u_1 \leq \max\left(\bar{u}_{i,j,k}, \bar{u}_{i+1,j,k}, \bar{u}_{i,j-1,k}, \bar{u}_{i+1,j-1,k}, \bar{u}_{i,j,k+1}, \bar{u}_{i+1,j,k+1}, \bar{u}_{i,j-1,k+1}, \bar{u}_{i+1,j-1,k+1}\right)$$

...

$$\min\left(\bar{u}_{i,j,k}, \bar{u}_{i-1,j,k}, \bar{u}_{i,j+1,k}, \bar{u}_{i-1,j+1,k}, \bar{u}_{i,j,k+1}, \bar{u}_{i-1,j,k+1}, \bar{u}_{i,j+1,k+1}, \bar{u}_{i-1,j+1,k+1}\right) \leq u_8 \leq \max\left(\bar{u}_{i,j,k}, \bar{u}_{i-1,j,k}, \bar{u}_{i,j+1,k}, \bar{u}_{i-1,j+1,k}, \bar{u}_{i,j,k+1}, \bar{u}_{i-1,j,k+1}, \bar{u}_{i,j+1,k+1}, \bar{u}_{i-1,j+1,k+1}\right)$$

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Formulation of MLP on Structured Grids

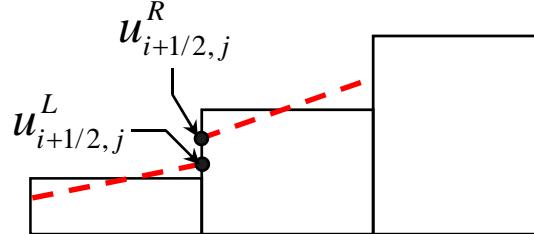


- TVD-MUSCL vs. MLP on Structured Mesh

$$u_{i+1/2,j}^L = \bar{u}_{i,j} + 0.5\phi(r_L)\Delta\bar{u}_{i-1/2,j}$$

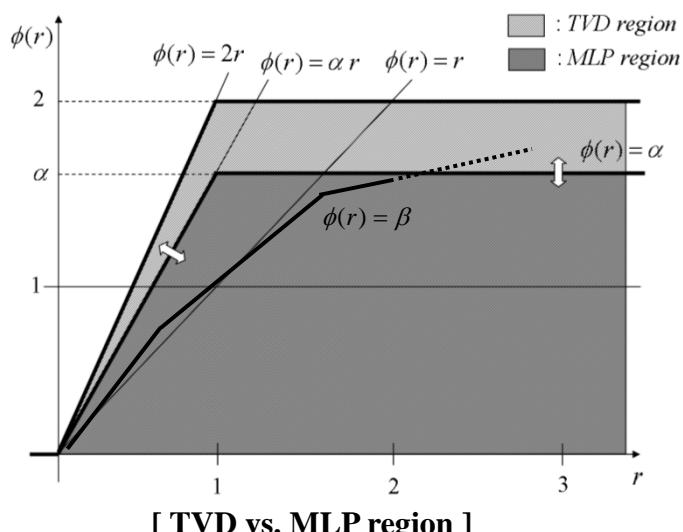
$$u_{i+1/2,j}^R = \bar{u}_{i+1,j,k} - 0.5\phi(r_R)\Delta\bar{u}_{i+3/2,j}$$

$$\text{with } r_L = \frac{\Delta\bar{u}_{i+1/2,j}}{\Delta\bar{u}_{i-1/2,j}}, r_R = \frac{\Delta\bar{u}_{i+1/2,j}}{\Delta\bar{u}_{i+3/2,j}}$$



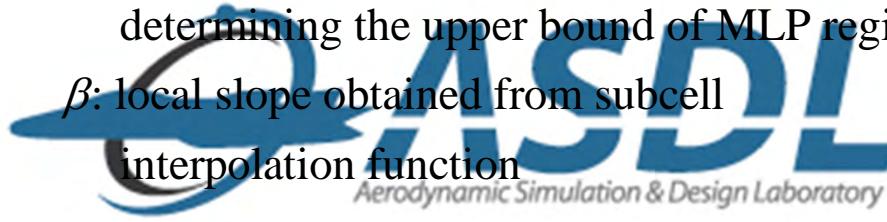
- Limiting region

- $\phi(r) = \max(0, \min(2, 2r))$: TVD limiting region based on 1-D analysis
- $\phi(r) = \max(0, \min(\alpha, \alpha r))$: MLP limiting region for multi-dimensional flows



- MLP limiting function

- $\phi(r) = \max(0, \min(\alpha, \alpha r, \beta))$ with
 - α : multi-dimensional restriction coefficient determining the upper bound of MLP region
 - β : local slope obtained from subcell interpolation function



Formulation of MLP on Structured Grids (Cont')



- Determination of Multi-dimensional Restriction Coefficient, α

- Step 1. Assume a cell-vertex value estimated by

$$u_{i+1/2,j+1/2} = \bar{u}_{i,j} + \Delta u_{i+1/2,j}^x + \Delta u_{i,j+1/2}^y = \bar{u}_{i,j} + (1 + r_{xy}) \Delta u_{i+1/2,j}^x \text{ with } r_{xy} = \Delta u_{i,j+1/2}^y / \Delta u_{i+1/2,j}^x$$

- Note that we only need to check the upper bound of the maximum vertex value and the lower bound of the minimum vertex values. Thus, r_{xy} is always positive.

- Step 2. Obtain the neighboring minimum and maximum values by checking all cell-averaged values sharing the same vertex point ($i + 1/2, j + 1/2$)

$$u_{i+1/2,j+1/2}^{\min} = \min(\bar{u}_{i,j}, \bar{u}_{i+1,j}, \bar{u}_{i,j+1}, \bar{u}_{i+1,j+1}), \quad u_{i+1/2,j+1/2}^{\max} = \max(\bar{u}_{i,j}, \bar{u}_{i+1,j}, \bar{u}_{i,j+1}, \bar{u}_{i+1,j+1})$$

- Step 3. Enforce the MLP limiting condition $\phi_{MLP}(r) = \min(\alpha, \alpha r)$ onto $u_{i+1/2,j+1/2}$

$$\bar{u}_{i+1/2,j+1/2}^{\min} \leq u_{i+1/2,j+1/2} \leq \bar{u}_{i+1/2,j+1/2}^{\max} \quad \text{or} \quad \bar{u}_{i+1/2,j+1/2}^{\min} \leq \bar{u}_{i,j} + (1 + r_{xy}) \Delta u_{i+1/2,j}^x \leq \bar{u}_{i+1/2,j+1/2}^{\max}$$

$$\frac{\bar{u}_{i+1/2,j+1/2}^{\min} - \bar{u}_{i,j}}{1 + r_{xy}} \leq \Delta u_{i+1/2,j}^x \leq \frac{\bar{u}_{i+1/2,j+1/2}^{\max} - \bar{u}_{i,j}}{1 + r_{xy}} \quad \text{or}$$

$$\frac{\bar{u}_{i+1/2,j+1/2}^{\min} - \bar{u}_{i,j}}{1 + r_{xy}} \leq 0.5 \phi_{MLP}(r_x) \Delta \bar{u}_{i-1/2,j} \leq \frac{\bar{u}_{i+1/2,j+1/2}^{\max} - \bar{u}_{i,j}}{1 + r_{xy}} \quad \text{with } r_x = \Delta \bar{u}_{i+1/2,j} / \Delta \bar{u}_{i-1/2,j}$$

Formulation of MLP on Structured Grids (Cont')



- Determination of Multi-dimensional Restriction Coefficient, α (cont'd)
 - (S4) Obtain the range of α

If $\Delta u_{i+1/2,j}^x > 0$ (local maximum),

$$0 \leq 0.5\phi(r_x)\Delta\bar{u}_{i-1/2,j} \leq \frac{\bar{u}_{i+1/2,j+1/2}^{\max} - \bar{u}_{i,j}}{1 + r_{xy}} \Rightarrow 0 \leq \alpha_{i+1/2,j+1/2} \leq \frac{2 \max(1, r_x)}{(1 + r_{xy})\Delta\bar{u}_{i+1/2,j}} (\bar{u}_{i+1/2,j+1/2}^{\max} - \bar{u}_{i,j})$$

If $\Delta u_{i+1/2,j}^x < 0$ (local minimum),

$$\frac{\bar{u}_{i+1/2,j+1/2}^{\min} - \bar{u}_{i,j}}{1 + r_{xy}} \leq 0.5\phi(r_x)\Delta\bar{u}_{i-1/2,j} \leq 0 \Rightarrow 0 \leq \alpha_{i+1/2,j+1/2} \leq \frac{2 \max(1, r_x)}{(1 + r_{xy})\Delta\bar{u}_{i+1/2,j}} (\bar{u}_{i+1/2,j+1/2}^{\min} - \bar{u}_{i,j})$$

Thus, we have

$$0 \leq \alpha_{i+1/2,j+1/2} \leq \left| \frac{2 \max(1, r_x)}{(1 + r_{xy})\Delta\bar{u}_{i+1/2,j}} \right| \min \left(\left| \bar{u}_{i+1/2,j+1/2}^{\max} - \bar{u}_{i,j} \right|, \left| \bar{u}_{i+1/2,j+1/2}^{\min} - \bar{u}_{i,j} \right| \right)$$

By checking four vertex points $(i \pm 1/2, j \pm 1/2)$, we finally have

$$0 \leq \alpha \leq \left| \frac{2 \max(1, r_x)}{(1 + r_{xy})\Delta\bar{u}_{i+1/2,j}} \right| \min \left(\left| \bar{u}_{i \pm 1/2, j \pm 1/2}^{\max} - \bar{u}_{i,j} \right|, \left| \bar{u}_{i \pm 1/2, j \pm 1/2}^{\min} - \bar{u}_{i,j} \right| \right) \text{ with } r_{xy} = \Delta u_{i,j+1/2}^y / \Delta u_{i \pm 1/2,j}^x$$

- Note that if $r_{xy} = 0$, α recovers the 1-D TVD form.

Formulation of MLP on Structured Grids (Cont')



- MLP Slope Limiters

- Cell-interface values in the x -direction

$$u_{i+1/2,j}^L = \bar{u}_{i,j} + 0.5\phi(r_L)\Delta\bar{u}_{i-1/2,j} = \bar{u}_{i,j} + 0.5 \max(0, \min(\alpha_L, \alpha_L r_L, \beta_L)) \Delta\bar{u}_{i-1/2,j}$$

$$u_{i+1/2,j}^R = \bar{u}_{i+1,j} - 0.5\phi(r_R)\Delta\bar{u}_{i+3/2,j} = \bar{u}_{i+1,j} - 0.5 \max(0, \min(\alpha_R, \alpha_R r_R, \beta_R)) \Delta\bar{u}_{i+3/2,j}$$

$$\alpha_L = \left| \frac{2 \max(1, r_L)}{(1+r_{xy}) \Delta\bar{u}_{i+1/2,j}} \right| \min \left(\left| \bar{u}_{i\pm 1/2,j\pm 1/2}^{\max} - \bar{u}_{i,j} \right|, \left| \bar{u}_{i\pm 1/2,j\pm 1/2}^{\min} - \bar{u}_{i,j} \right| \right), \quad \alpha_R = \alpha_L|_{i \rightarrow i+1}$$

with $r_L = \Delta\bar{u}_{i+1/2,j} / \Delta\bar{u}_{i-1/2,j}$ and $r_R = \Delta\bar{u}_{i+1/2,j} / \Delta\bar{u}_{i+3/2,j}$

- Choice of local slopes (β)

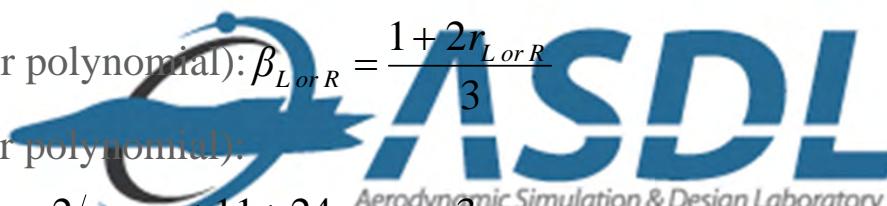
- MLP-van Leer: $\beta_{L \text{ or } R} = \frac{2r_{L \text{ or } R}}{1 + |r_{L \text{ or } R}|}$

- MLP-superbee: $\beta_{L \text{ or } R} = \begin{cases} \min(2r_{L \text{ or } R(i-1)}, 1), & \text{if } 0 < r_{L \text{ or } R(i-1)} < 1 \\ \min(r_{L \text{ or } R(i-1)}, 2), & \text{if } r_{L \text{ or } R(i-1)} > 1 \end{cases}$ with $r_{R(i-1)} = \frac{\Delta\bar{u}_{i-1/2,j}}{\Delta\bar{u}_{i+1/2,j}}$

- MLP3 (MLP slope limiting with 3rd-order polynomial): $\beta_{L \text{ or } R} = \frac{1+2r_{L \text{ or } R}}{3}$

- MLP5 (MLP slope limiting with 5th-order polynomial):

$$\beta_L = \frac{-2/r_{L(i-1)} + 11 + 24r_L - 3r_L r_{L(i+1)}}{30}, \quad \beta_R = \frac{-2/r_{R(i+2)} + 11 + 24r_{R(i+1)} - 3r_{R(i+1)} r_R}{30}$$





- Maximum Principle

- Maximum principle is well-established in parabolic and elliptic PDEs as a tool for L_∞ stability.
- TVD condition, which is a sufficient condition for the maximum principle, is not available in multi-dimensional situation.
- A condition ensuring multi-dimensional monotonicity
 - Cockburn *et al.* (1990), Liu (1993), Barth (2003), Kim *et al.* (2008, 2010, 2012)

- L_∞ Stability of MLP Limiting

- *For a multi-dimensional hyperbolic scalar conservation law of*

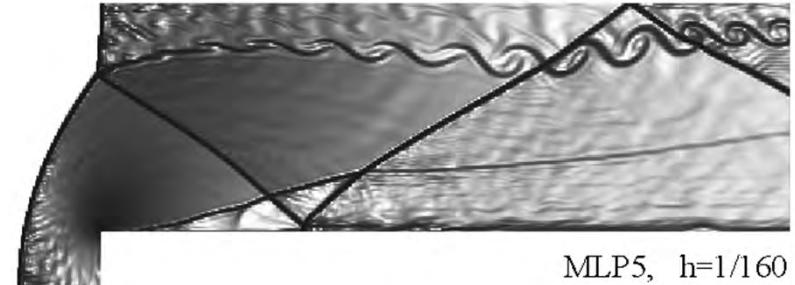
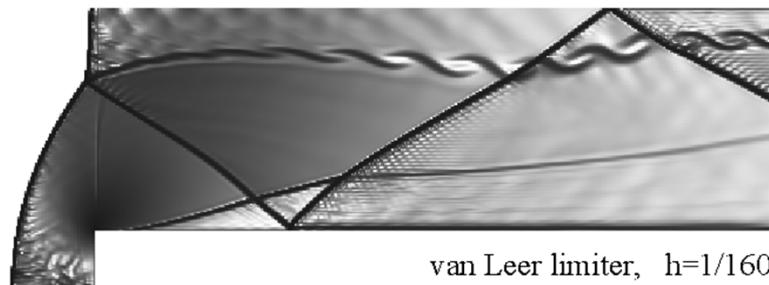
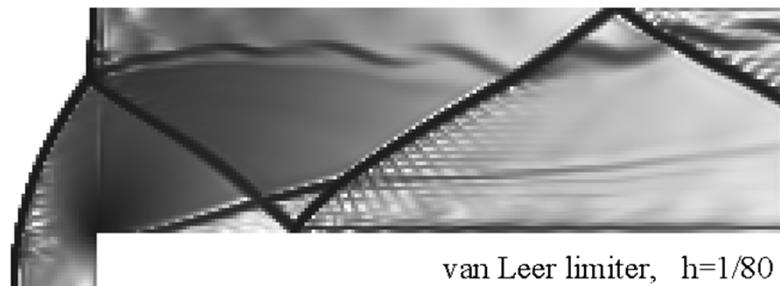
$$\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} + \frac{\partial g(u)}{\partial y} = 0,$$

the fully discrete scheme using the MLP limiting satisfies the local maximum principle under a suitable CFL condition.

If $\bar{u}_{neighbor}^{\min,n} \leq \bar{u}_{i,j}^n \leq \bar{u}_{neighbor}^{\max,n}$, then $\bar{u}_{neighbor}^{\min,n} \leq \bar{u}_{i,j}^{n+1} \leq \bar{u}_{neighbor}^{\max,n}$.

Numerical Results

- A Mach 3 Wind Tunnel with a Step
 - A standard test case for high resolution schemes
 - Advantage in resolving the slip lines



< Numerical Schlieren of a Mach 3 wind tunnel with a step problem >



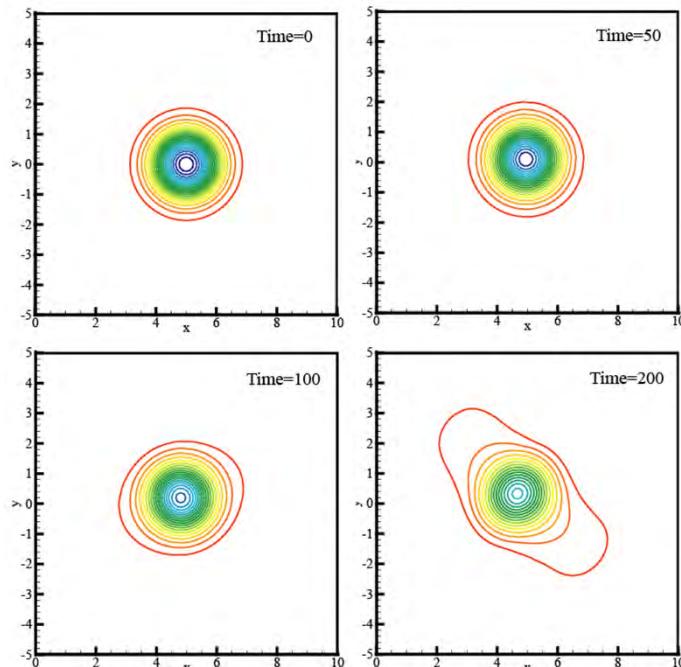
Numerical Results (Cont')



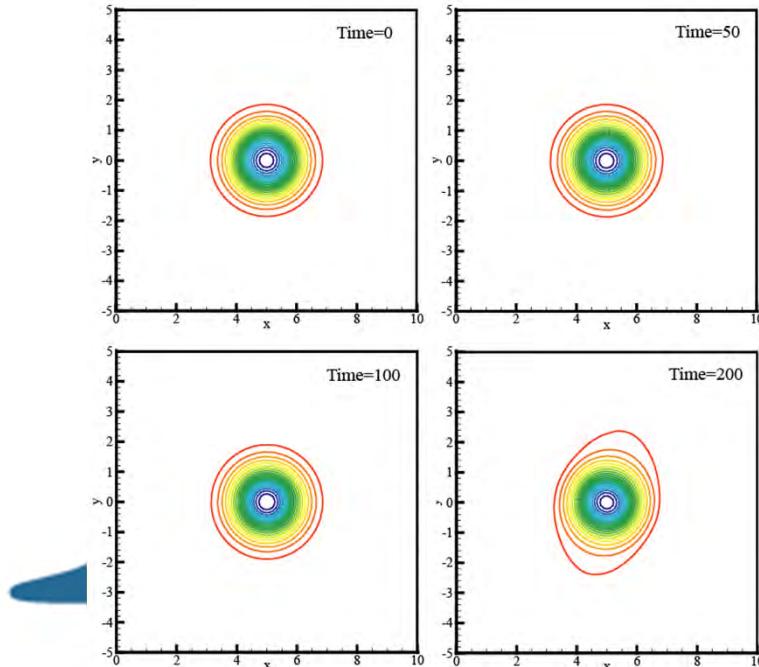
- Isentropic Vortex Advection
 - Initial mean flow and perturbation values

$u_\infty = 1, v_\infty = 1, p_\infty = \rho_\infty = T_\infty = 1$ with

$$(\delta u, \delta v) = \frac{\beta}{2\pi} e^{(1-r^2)/2} (-\bar{y}, \bar{x}), \delta T = -\frac{(\gamma-1)\beta^2}{8\gamma\pi^2} e^{1-r^2}.$$



< van Leer limiter >



< MLP5 >





Numerical Results (Cont')



| Scheme | Size | L_1 error | L_1 order | L_∞ error | L_∞ order |
|------------------------|-----------|-------------|-------------|------------------|------------------|
| van Leer limiter | 50 × 50 | 4.6165E−03 | — | 8.9412E−02 | — |
| | 100 × 100 | 1.0489E−03 | 2.14 | 2.2579E−02 | 1.99 |
| | 150 × 150 | 4.5167E−04 | 2.08 | 1.0169E−02 | 1.97 |
| | 200 × 200 | 2.5121E−04 | 2.04 | 5.6879E−03 | 2.02 |
| MLP-van Leer limiter | 50 × 50 | 3.7699E−03 | — | 8.0859E−02 | — |
| | 100 × 100 | 7.1064E−04 | 2.41 | 2.0432E−02 | 1.98 |
| | 150 × 150 | 2.9575E−04 | 2.16 | 9.2804E−03 | 1.95 |
| | 200 × 200 | 1.5657E−04 | 2.21 | 5.1476E−03 | 2.05 |
| MLP3 | 50 × 50 | 2.0713E−03 | — | 3.9473E−02 | — |
| | 100 × 100 | 2.9620E−04 | 2.81 | 7.2746E−03 | 2.44 |
| | 150 × 150 | 9.8557E−05 | 2.71 | 2.9668E−03 | 2.21 |
| | 200 × 200 | 4.3257E−05 | 2.86 | 1.4262E−03 | 2.55 |
| MLP5 | 50 × 50 | 1.1441E−03 | — | 2.4834E−02 | — |
| | 100 × 100 | 2.2427E−04 | 2.35 | 5.5389E−03 | 2.16 |
| | 150 × 150 | 7.6479E−05 | 2.65 | 1.9767E−03 | 2.54 |
| | 200 × 200 | 4.0756E−05 | 2.19 | 1.0291E−03 | 2.27 |
| Third order polynomial | 50 × 50 | 1.5736E−03 | — | 4.1150E−02 | — |
| | 100 × 100 | 2.6245E−04 | 2.58 | 6.3480E−03 | 2.70 |
| | 150 × 150 | 9.6728E−05 | 2.46 | 2.0090E−03 | 2.84 |
| | 200 × 200 | 4.9086E−05 | 2.36 | 8.9188E−04 | 2.82 |

< Grid refinement test >

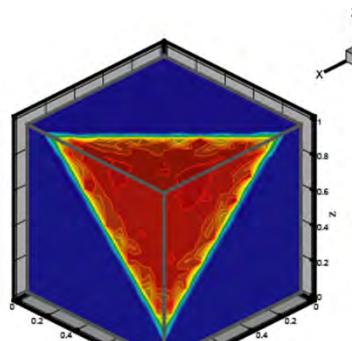
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Numerical Results (Cont')



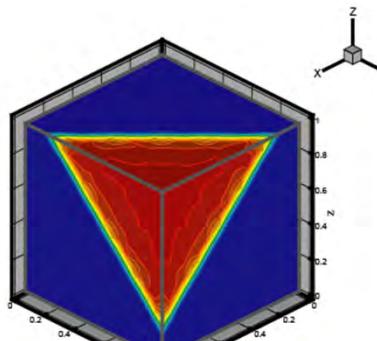
- Three-dimensional Normal Shock Discontinuity
 - Examine the shock-capturing characteristics of TVD and MLP
 - Normal shock discontinuity inclined by 45 degree angle to each cell-surface

< Density contour >

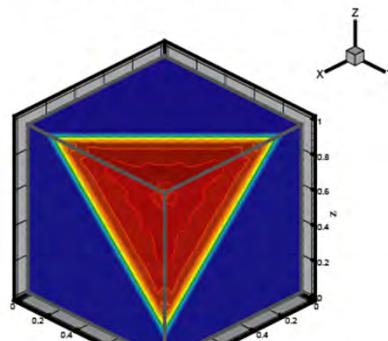


a. Roe's FDS with superbee limiter

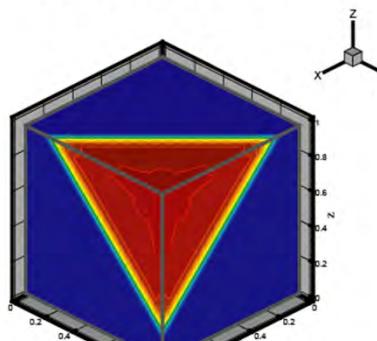
< Density contour >



b. Roe's FDS with van Leer limiter

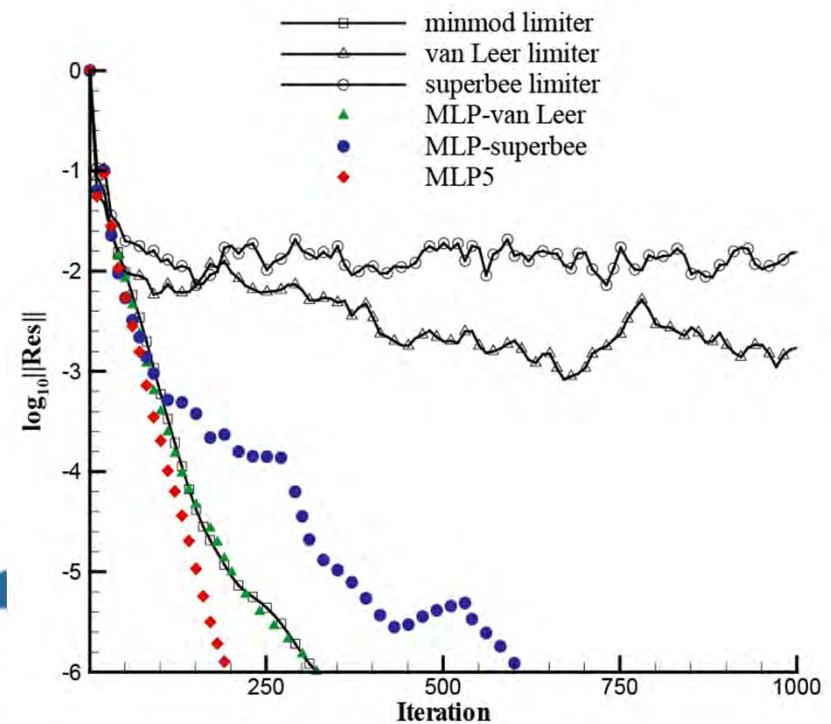


c. Roe's FDS with MLP-superbee



d. Roe's FDS with MLP-van Leer

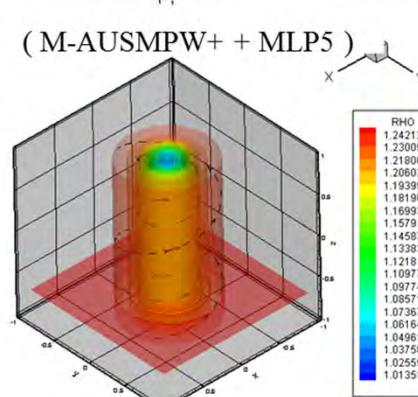
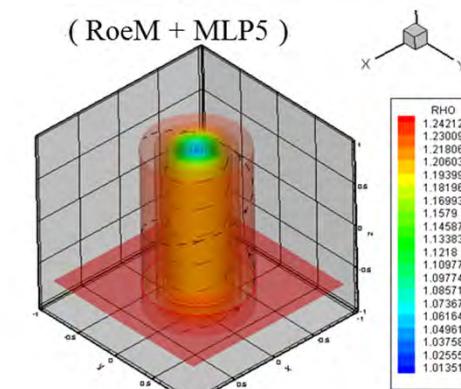
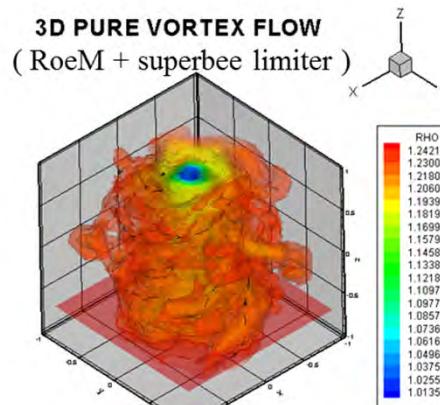
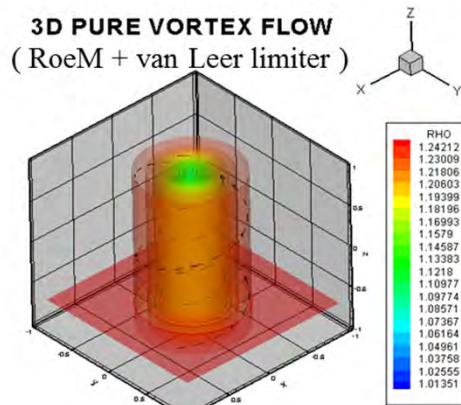
< Error history >



Numerical Results (Cont')

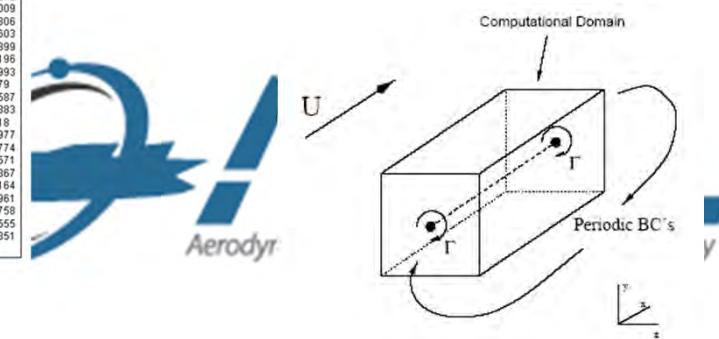
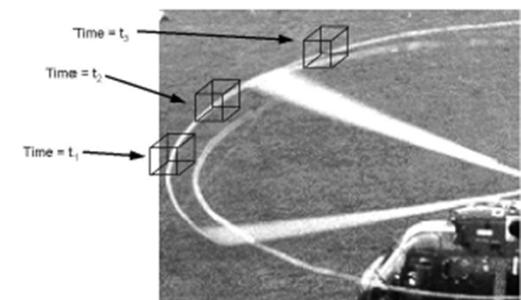


- Three-dimensional Pure Vortex Flow
 - A simple model of a tip vortex element released from a rotor blade.
 - Thomson-Rankine vortex model with $50 \times 50 \times 50$ grids



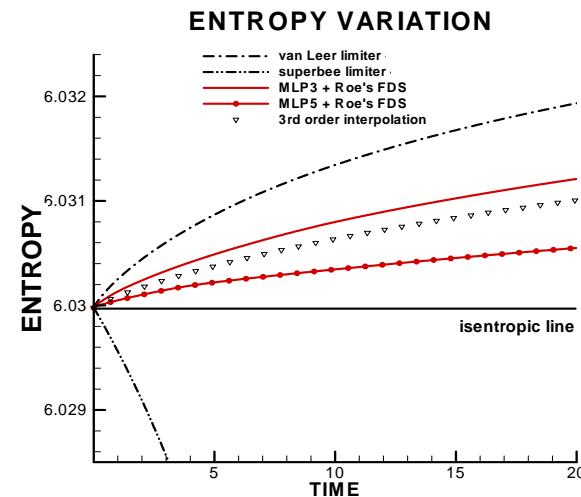
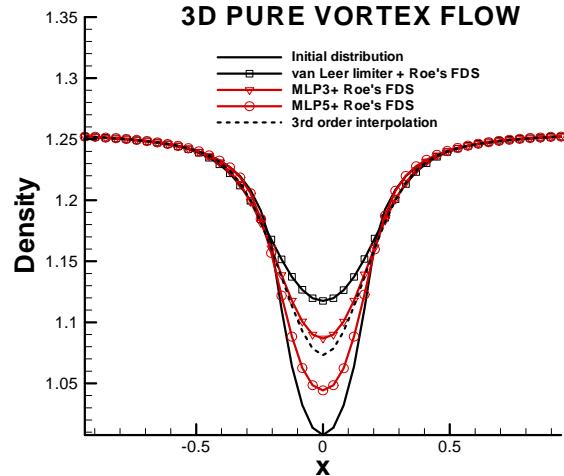
< Density contours at non-dimensional time 20.0 >

$$v_\theta = \begin{cases} \omega r, & r < a \\ \frac{\omega a^2}{r}, & r > a \end{cases} \text{ with } \omega = \text{const}$$

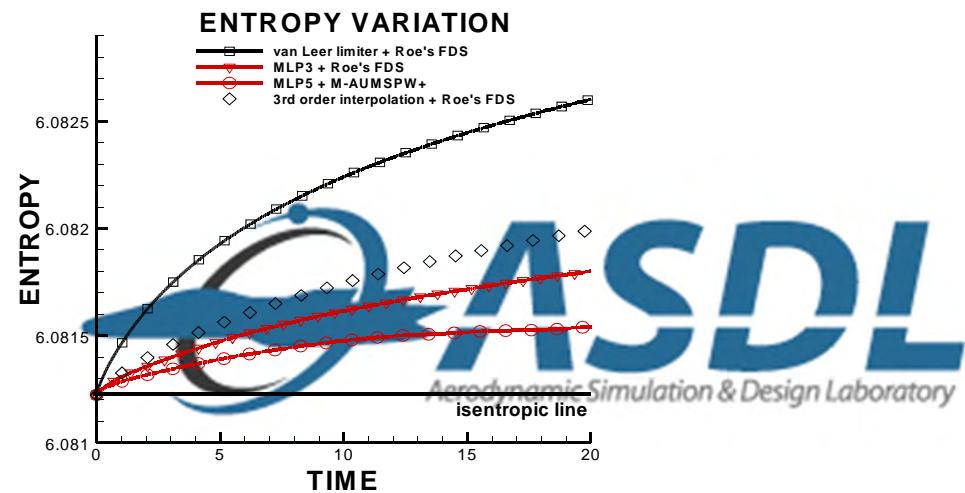
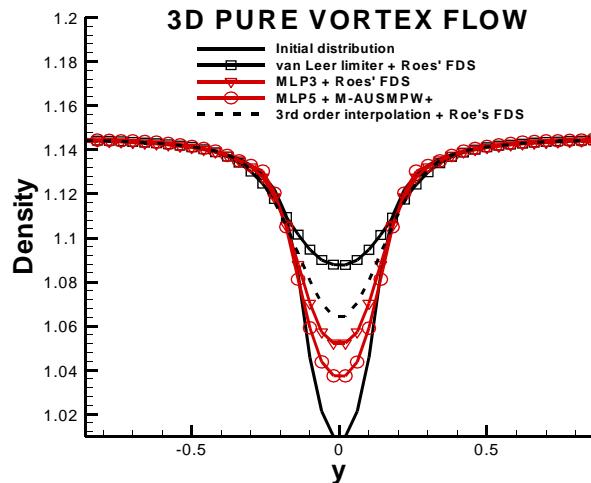


Numerical Results (Cont')

- The Case of Vortex Core Aligned with z-axis



- The Case of Vortex Core Non-aligned with z-axis

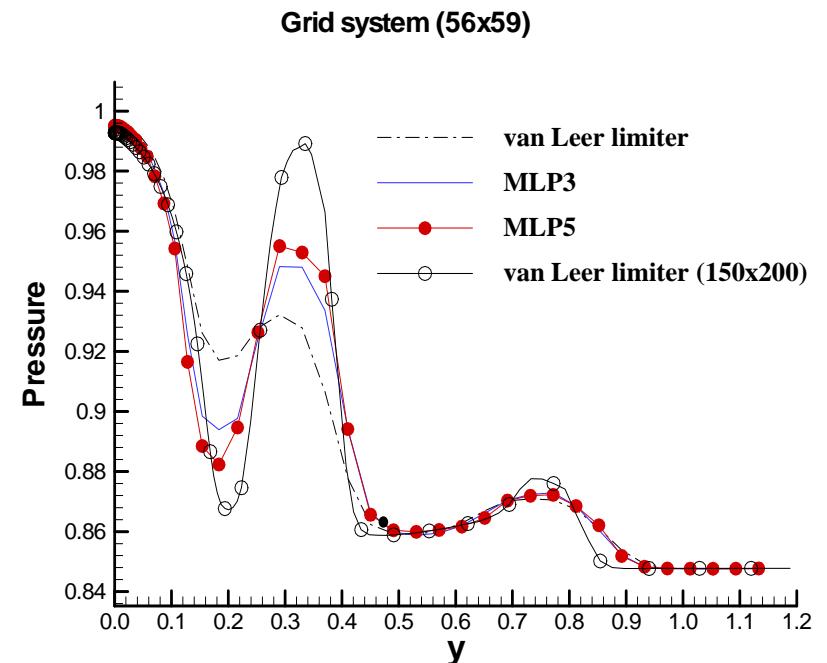
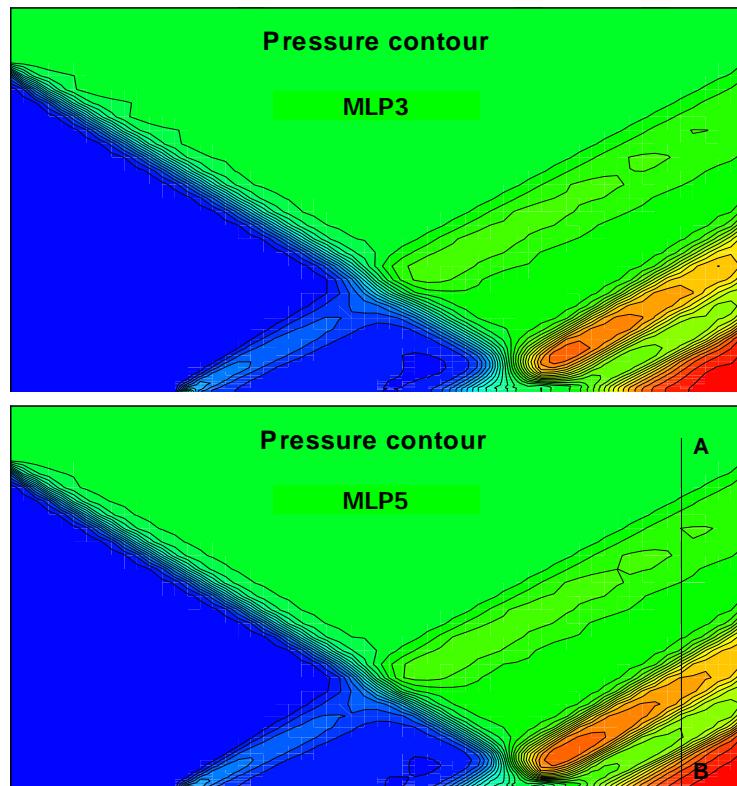


Numerical Results (Cont')



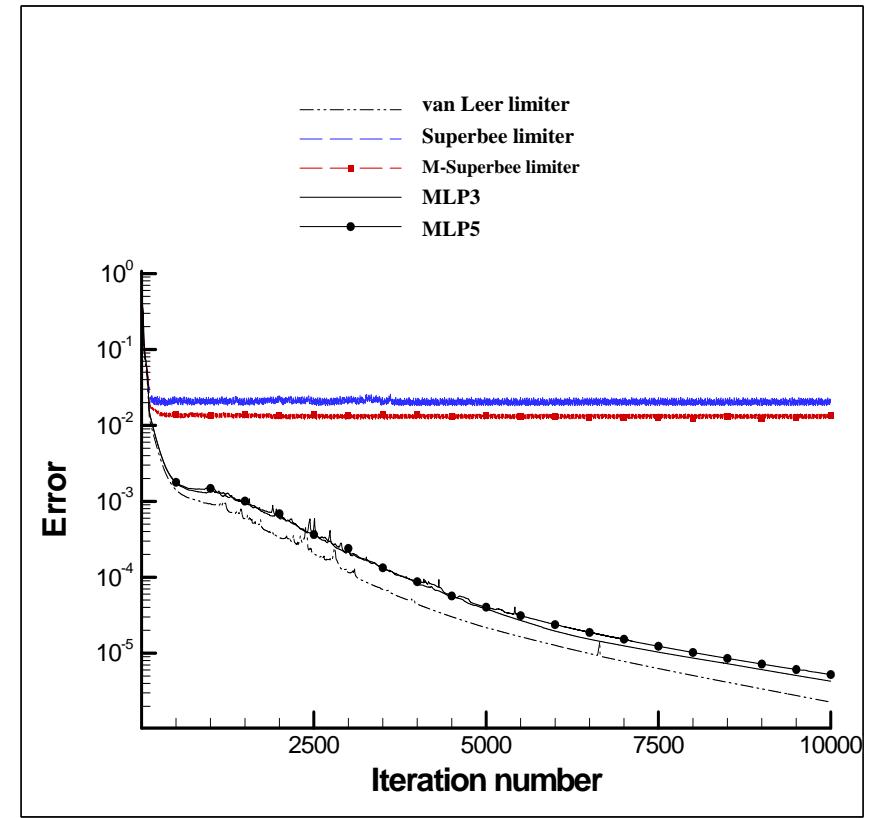
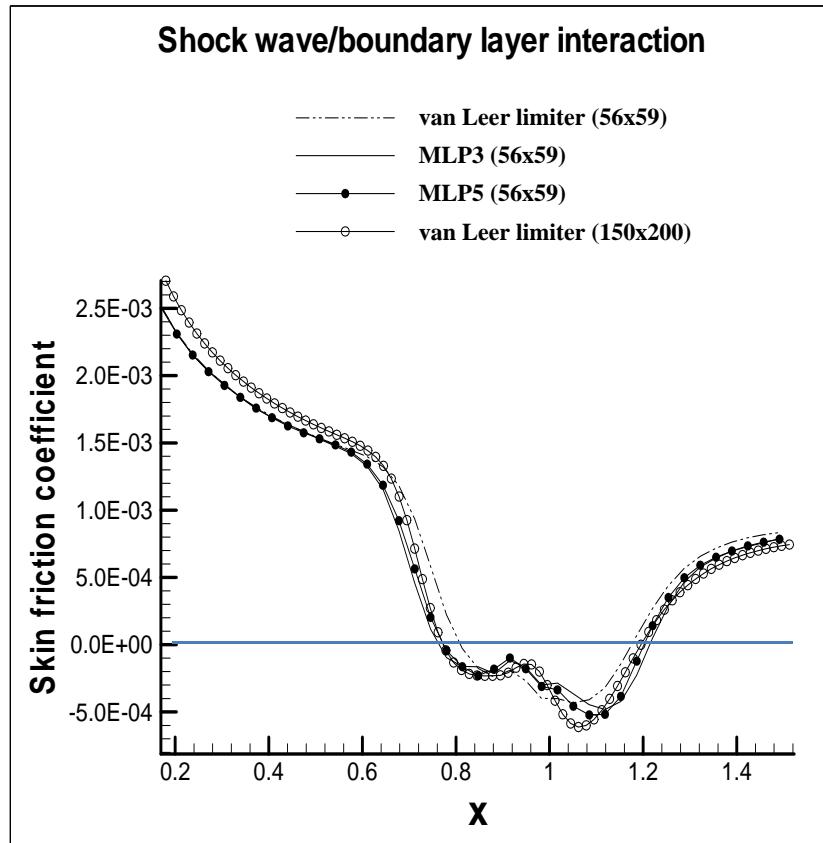
- Laminar Shock-Boundary Layer Interaction

- Free stream Mach number : $M_\infty = 2.0$, shock impinging angle : $\theta = 32.585$ degree
- Reynolds number : $Re = 2.96 \times 10^5$
- Grid system : $56 \times 59, 150 \times 200$



< Comparison of pressure distribution along line AB >

Numerical Results (Cont')



< Skin friction coefficient and error history >

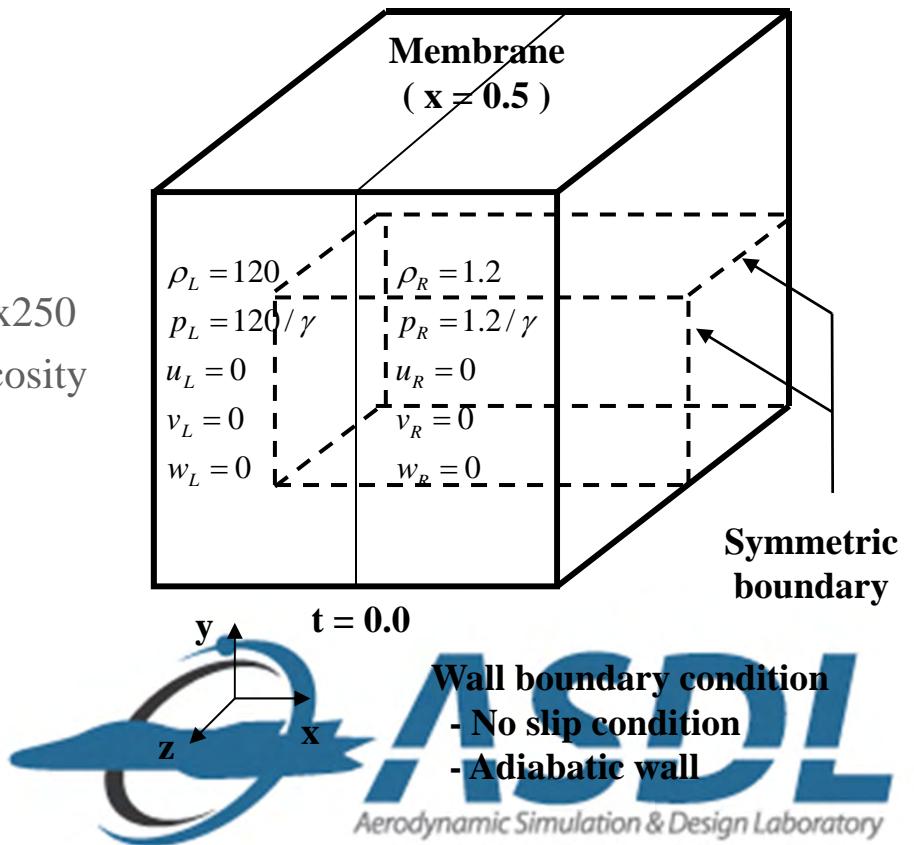




● 3-D Viscous Shock Tube Problem

● Computed conditions

- Spatial discretization
 - AUSMPW+, M-AUSMPW+
- Time integration
 - 3rd-order TVD Runge-Kutta
- Grid system
 - 250x125x125, 350x175x175, 500x250x250
- Reynolds number: 200 with constant viscosity
- Boundary values
 - No slip B.C (x=0, x=1, y=0, z=0)
 - Symmetric B.C (y=0.5, z=0.5)

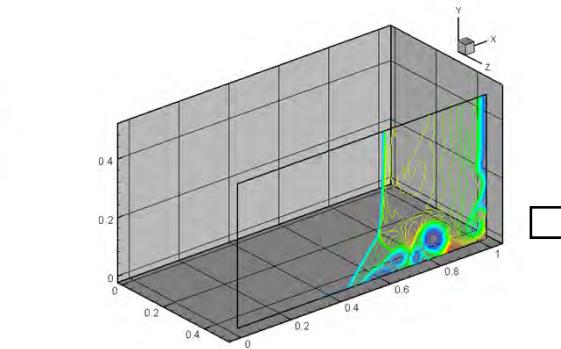
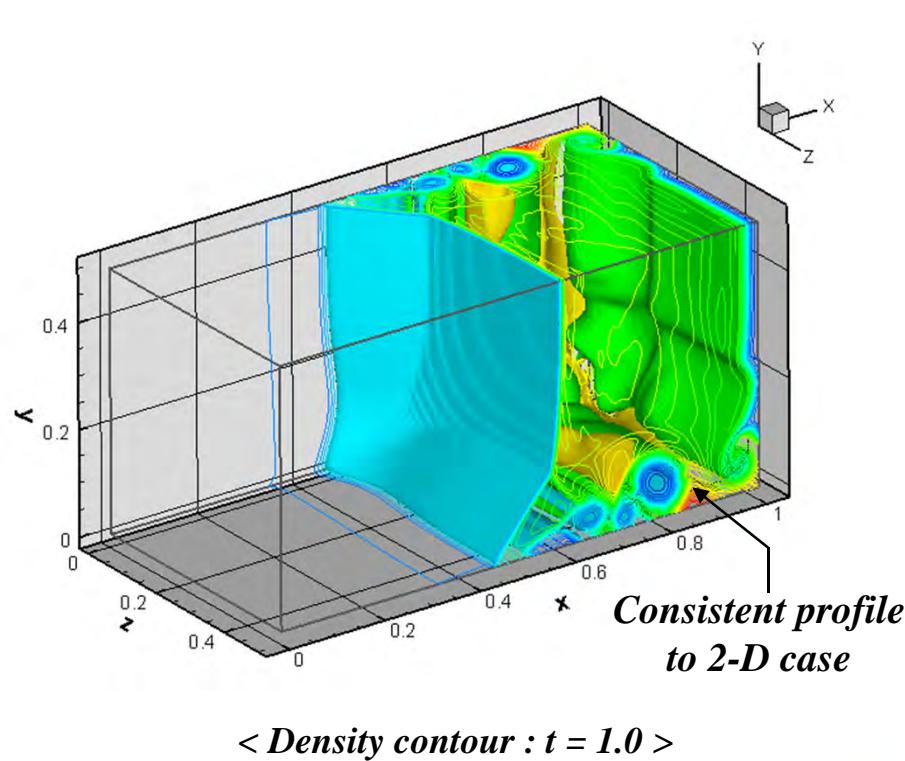


Numerical Results (Cont')

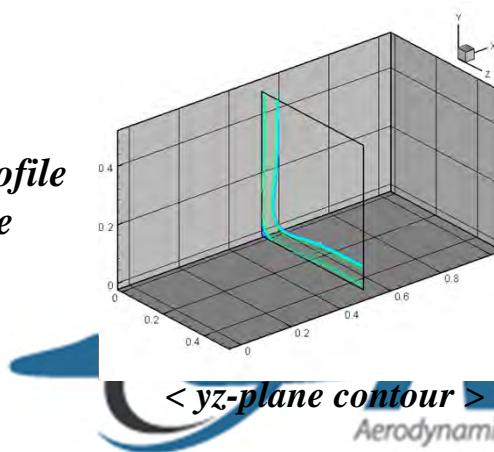


● Profiles of 3-D Viscous Shock Tube Problem

- Complicated shock-boundary interaction, and vortical structure



*Near the wall,
distinct 3-D vortex
structure*

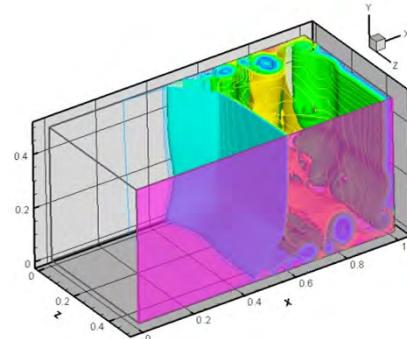


*Similar to corner
flow, built-up swirl
flow*

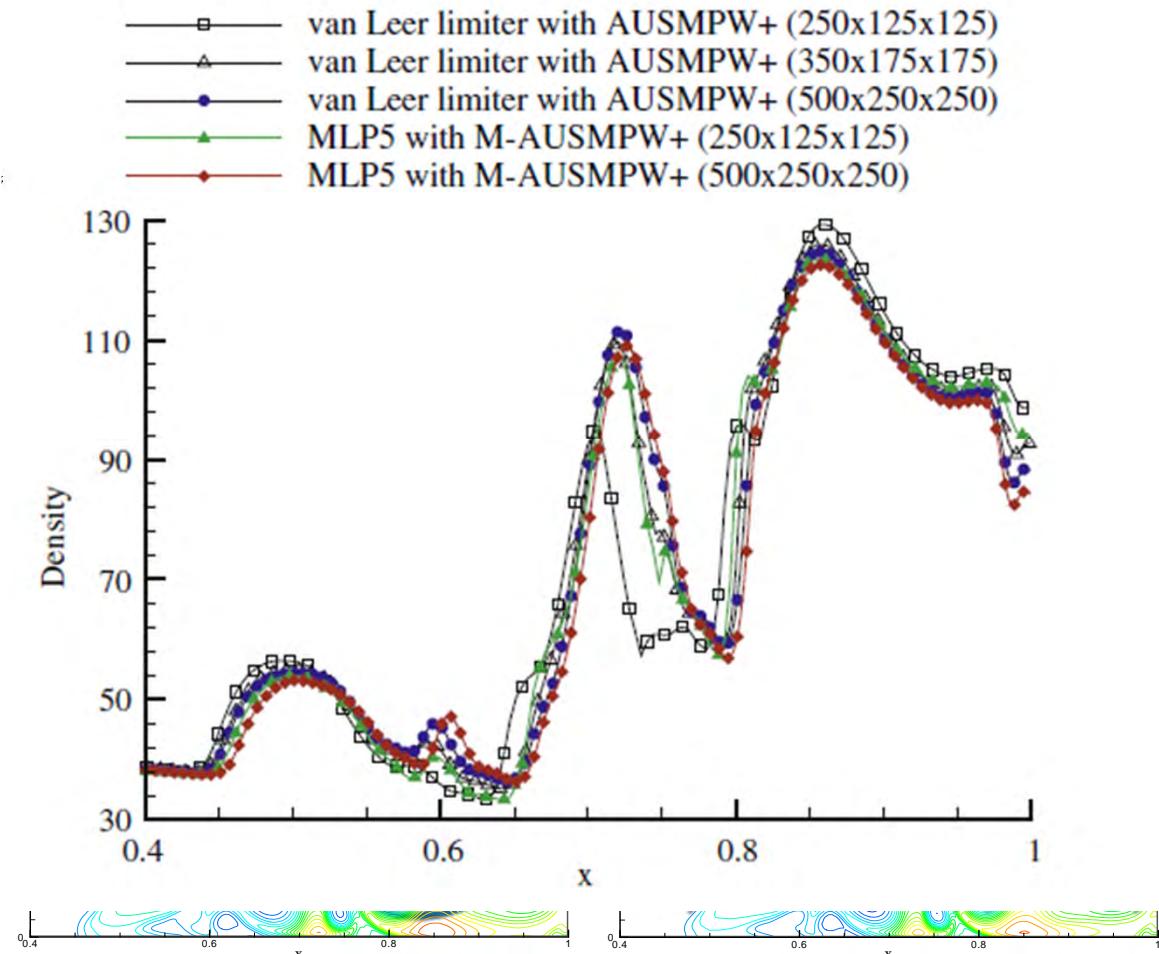
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Numerical Results (Cont')

Comparison :
*profile of
separated vortex*



| | No. of grid pts. |
|-------------|------------------------------------|
| Coarse mesh | 3,906,250 (250x125x125) |
| Medium mesh | 10,718,750 (350x175x175) |
| Fine mesh | 31,250,000 (500x250x250) |





Concluding Remarks



- In order to obtain high-resolution flow structure, higher-order numerical scheme is essential.
 - Treatment of discontinuous and smooth region in multiple dimensions
- From one- to multi-dimensional flows, a limiting process must include to the treatment of vertex to efficiently include a missing flow physics at 1-D based limiting strategy.
 - Overcome the weakness of dimensional splitting approach, and eliminate spurious oscillations caused by multi-dimensional discontinuity.
- Characteristics of MLP
 - Multi-dimensional monotonicity guided by the MLP condition and the discrete maximum principle → L_∞ stability
 - Slopes based on higher-order polynomial reconstruction
 - Computational efficiency comparable to conventional TVD limiters
 - Enhanced accuracy and convergence characteristics

