

Universality of random matrices, Dyson Brownian Motion and Quantum Unique Ergodicity

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Eugene Wigner (1956): Successive energy levels of large nuclei have universal spacing statistics that can be described by random matrices.

— a grand vision

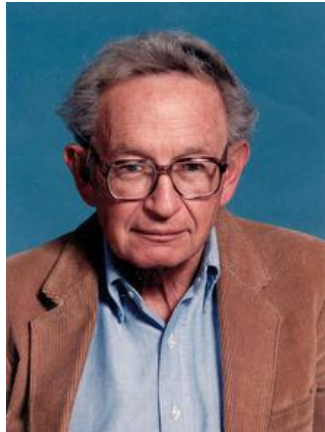
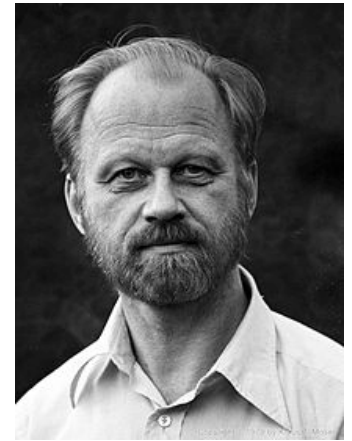


Freeman Dyson (1962): A gas of particles with a logarithmic interaction reaches local equilibrium very fast.

— a seminal idea

De Giorgi, Nash, Moser
(1957-60):

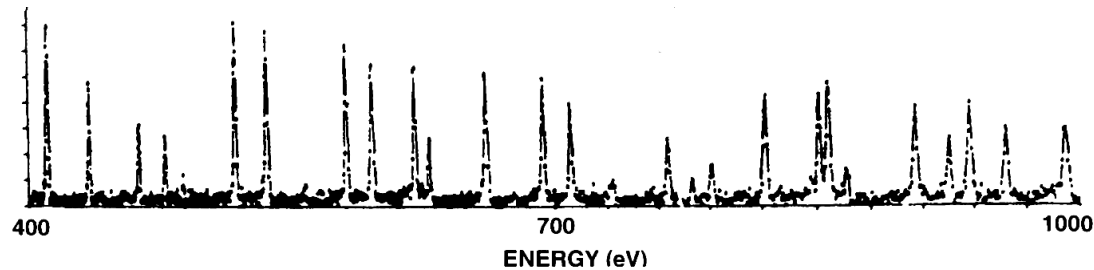
Regularity theory for
parabolic equations



Quantum Chaos conjecture (1977-84)
and
Anderson Model (1958)

Rudnick-Sarnak (1991) eigenfunctions on negative curved manifolds are
“flat” (Quantum unique ergodicity).

Experimental data for excitation spectra of heavy nuclei:

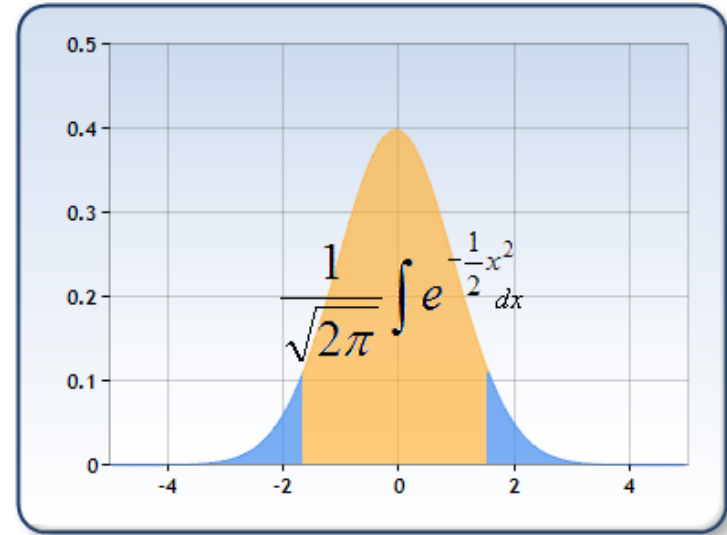


Wigner (1955): This pattern can be modeled by the **spacing distribution** of eigenvalues of random matrices.

Perhaps I am now too courageous when I try to guess the distribution of the distances between successive levels (of energies of heavy nuclei). Theoretically, the situation is quite simple if one attacks the problem **in a simpleminded fashion**. The question is simply what are the **distances** of the characteristic values of a **symmetric matrix with random coefficients**. — E. Wigner

Gaussian Orthogonal Ensemble (GOE):

$$H = \begin{pmatrix} h_{11} & h_{12} & \dots & h_{1N} \\ h_{21} & h_{22} & \dots & h_{2N} \\ \vdots & \vdots & & \vdots \\ h_{N1} & h_{N2} & \dots & h_{NN} \end{pmatrix}$$



$h_{jk} = h_{kj}$ (for $j < k$) are real independent normal random variables

$$\mathbb{E}h_{jk} = 0, \quad \mathbb{E}|h_{jk}|^2 = \frac{1}{N}$$

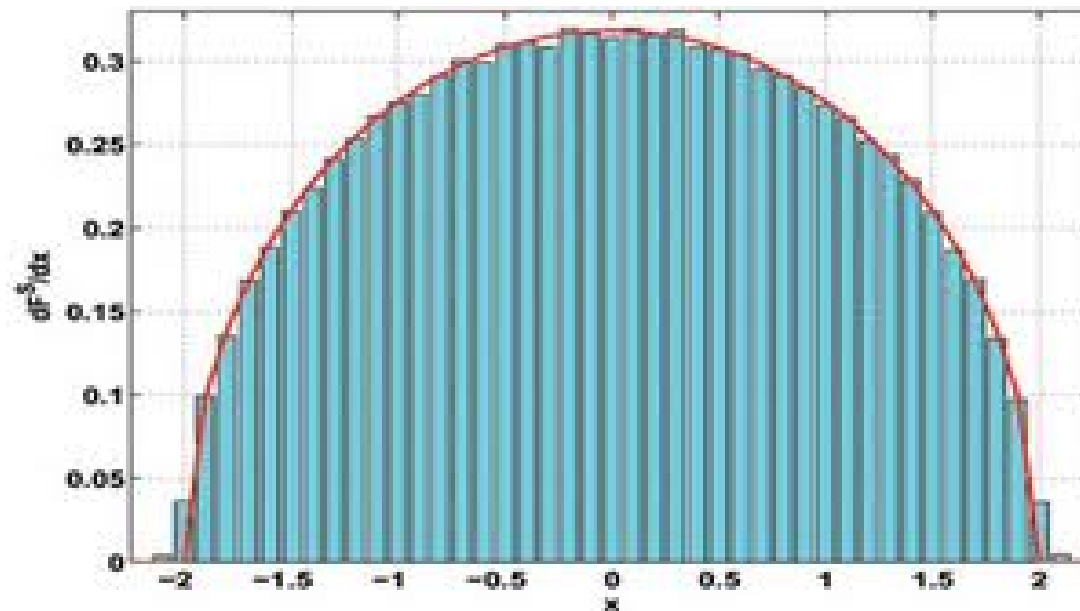
The eigenvalues $\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_N$ are of order one.

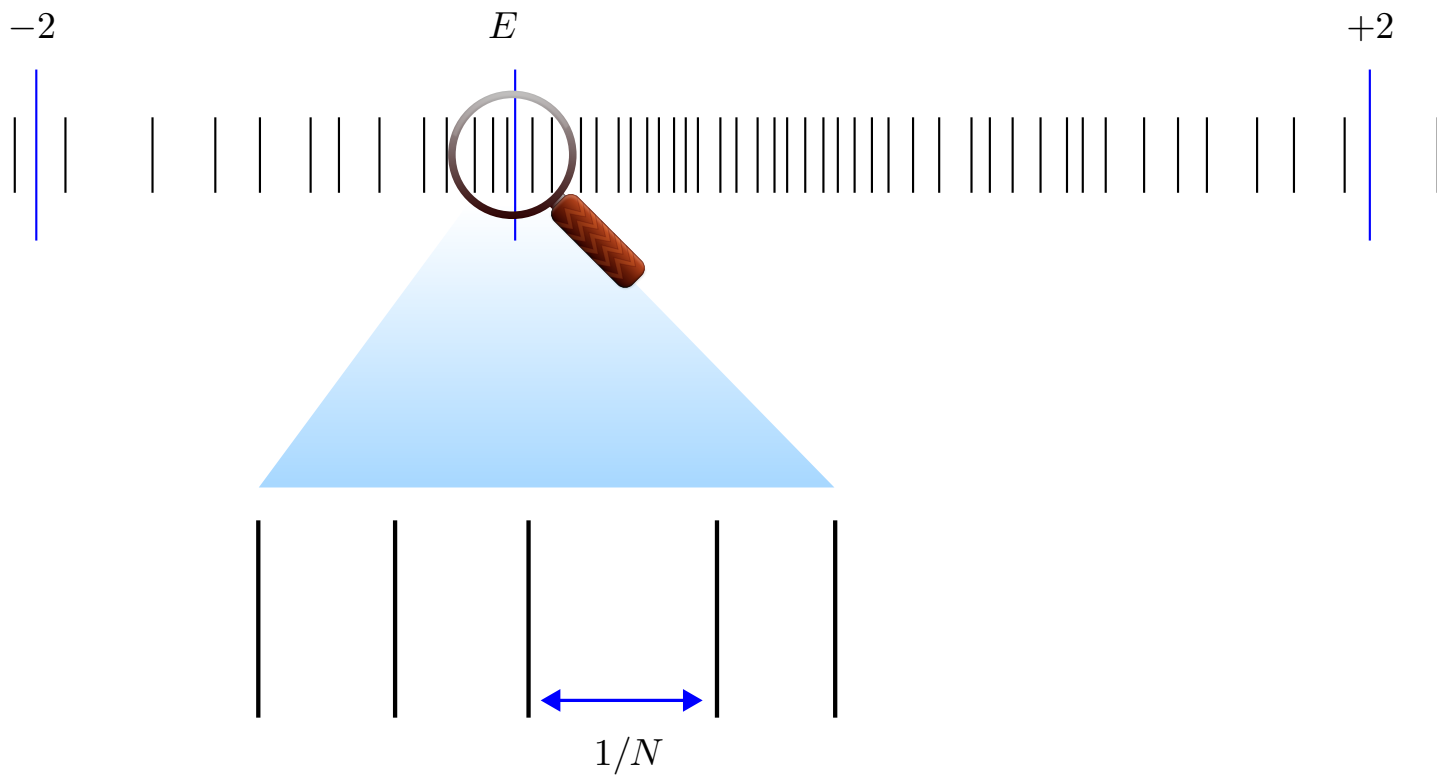
Also, Hermitian or quaternion self-dual (GUE, GSE) ensemble: classical ensembles by **Dyson's classification**.

Wigner ensembles: h_{ij} are just independent (not necessarily normal) distributions.

Wigner semicircle law: Density of eigenvalues

$$\rho(x) \sim \sqrt{4 - x^2}$$

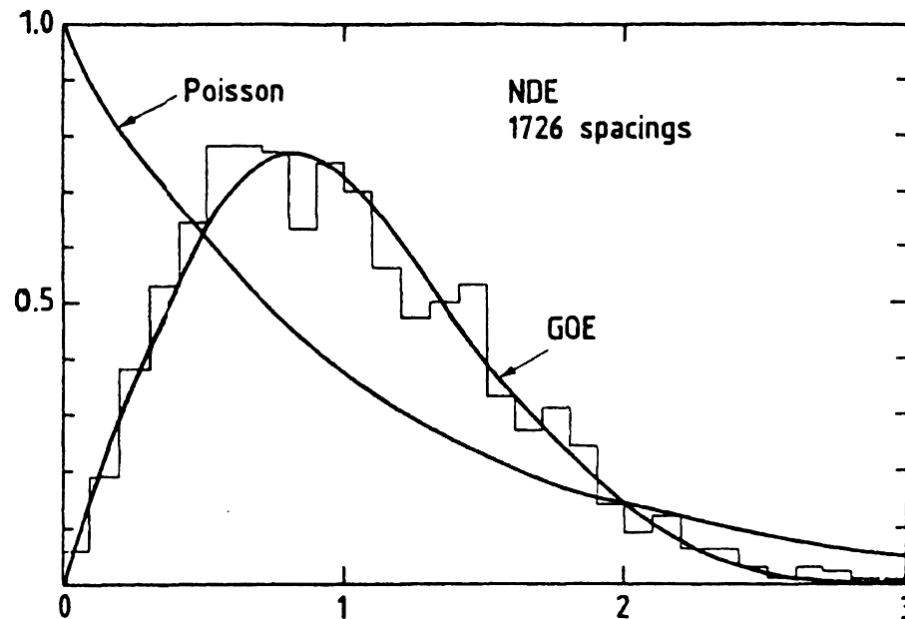




$$\mathbb{P}\left(\lambda_1 \sim E + \frac{x_1}{N}, \lambda_2 \sim E + \frac{x_2}{N}\right) \sim 1 - \left(\frac{\sin \pi(x_1 - x_2)}{\pi(x_1 - x_2)}\right)^2$$

Computation by Gaudin-Mehta and Dyson

By 1970 we had decided that random matrix theory was a beautiful piece of pure mathematics having nothing to do with physics. **Random matrix theory went temporarily to sleep.** — Dyson



Nearest-neighbor spacing distribution for the "Nuclear Data Ensemble" comprising 1726 spacings histogram (1983)

Is this exact computation for Gaussian models valid for general systems?

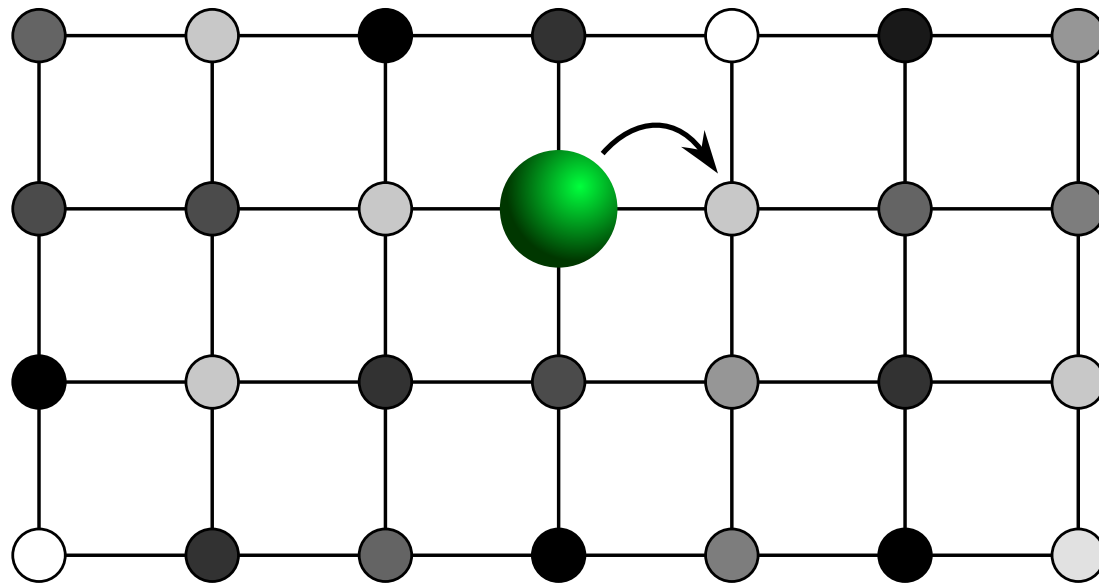
Fundamental belief of universality: Random matrix statistics are a new class of universal laws for highly correlated systems.

A simple manifestation (**Wigner-Dyson-Mehta conjecture**) : If h_{ij} are independent, then the local eigenvalue statistics are the same as those of the Gaussian ensembles.

Only symmetry types of the ensembles matter.

Quantum Chaos and Anderson Model

Anderson (1958): V_ω random potential on \mathbb{R}^d or \mathbb{Z}^d .
 $H = -\Delta + \lambda V_\omega$. Local eigenvalue statistics:



Depending on λ and d , there are two distinct regimes.

GOE in the **delocalization regime** (e.g., small λ in 3 dimension)

Poisson statistics (**Minami**) in the **localization regime** (**Frohlich-Spencer, Aizenman**).

Similarly for $-\Delta$ on a domain with chaotic classical dynamics.
Quantum Chaos conjecture: **Bohigas-Giannoni-Schmit (1984)**

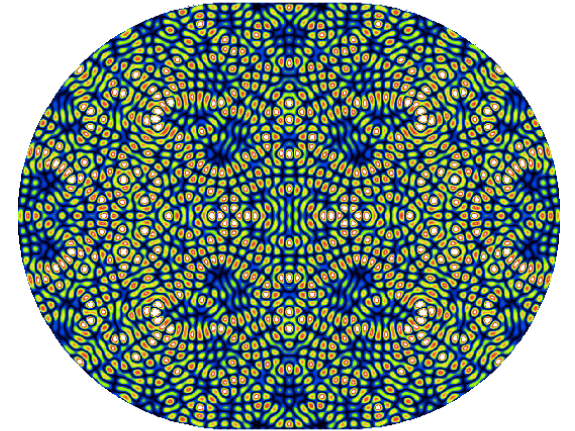
For integrable dynamics: **Berry-Tabor (1977)**

Quantum unique ergodicity conjecture
(QUE) [Rudnick-Sarnak]:

$(\psi_k)_{k \geq 1}$ orthogonal eigenfunctions of M
with negative curvature:

$$\int a(x) |\psi_j(x)|^2 d\text{Vol}(x) \rightarrow \int a(x) d\text{Vol}(x)$$

for any large energy limit.



Quantum Ergodicity Theorem: Averaged version of QUE [Snirelman, Colin de Verdiere, Zelditch] Anantharaman. Arithmetic QUE (Luo-Sarnak, Bourgain, Lindenstrauss, Soundararajan)

Question: Does **QUE hold for Wigner matrices?** Yes for GUE/GOE because the eigenvectors are uniformly distributed on $U(N)/O(N)$.

Quantum (unique) ergodicity/delocalization

\iff **random matrix statistics** ?

Two basic mathematical questions:

1. Universality conjecture for Wigner matrices
2. QUE for Wigner matrices.

Key ingredient in solving these problems: **Dyson Brownian motion.**

Generalized Wigner Ensembles

$$H = (h_{ij})_{1 \leq i, j \leq N}, \quad \bar{h}_{ji} = h_{ij} \quad \text{independent}$$

$$\mathbb{E}h_{ij} = 0, \quad \mathbb{E}|h_{ij}|^2 = s_{ij}, \quad \sum_i s_{ij} = 1, \quad \frac{c}{N} \leq s_{ij} \leq \frac{C}{N}$$

If h_{ij} are i.i.d. then it is called **Wigner ensemble**.

Solution to the Universality conjecture

Theorem [Erdos-Schlein-Y-Yin, 2009-2010] Local eigenvalue statistics are **universal for generalized Wigner ensembles**. in the bulk (**and edge**). Matrices with Bernoulli entries with varying variances are included. Assumption on the distribution of matrix elements: $4 + \varepsilon$ moments are uniformly bounded.

[Tao-Vu, 2009-10] Wigner matrices with **four moments matching those of $N(0, 1)$** . Hermitian case via combination with Johansson's result.

Extensions to sparse matrices, β -ensembles, single gap universality, the edge universality [Bourgade, Erdos, Knowles, Yin]

Theorem [Bourgade-Y, 2013] [Probabilistic local QUE]

$u_i, i = 1, \dots, N$: normalized e-vectors of generalized Wigner matrices. $\mathcal{N}_1, \mathcal{N}_2$: independent normal distribution $N(0, 1)$. Then for any j, k in the bulk and \mathbf{q} a unit vector in \mathbb{R}^N , we have

$$\sqrt{N} \left(|\langle \mathbf{q}, u_j \rangle|, |\langle \mathbf{q}, u_k \rangle| \right) \rightarrow \left(|\mathcal{N}_1|, |\mathcal{N}_2| \right)$$

For any k in the bulk, the eigenvector u_k is flat with high probability, i.e., probabilistic local QUE holds.

There is ε such that for any $\delta > 0$ and any k in the bulk and $A \subset \llbracket 1, \dots, N \rrbracket$ with $|A| \geq N^{C\varepsilon}$ we have,

$$\mathbb{P} \left(\left| \frac{1}{|A|} \sum_{a \in A} \left(N |u_k(a)|^2 - 1 \right) \right| > \delta \right) \leq C_\delta N^{-\varepsilon}.$$

Delocalization [Erdos-Schlein-Y]

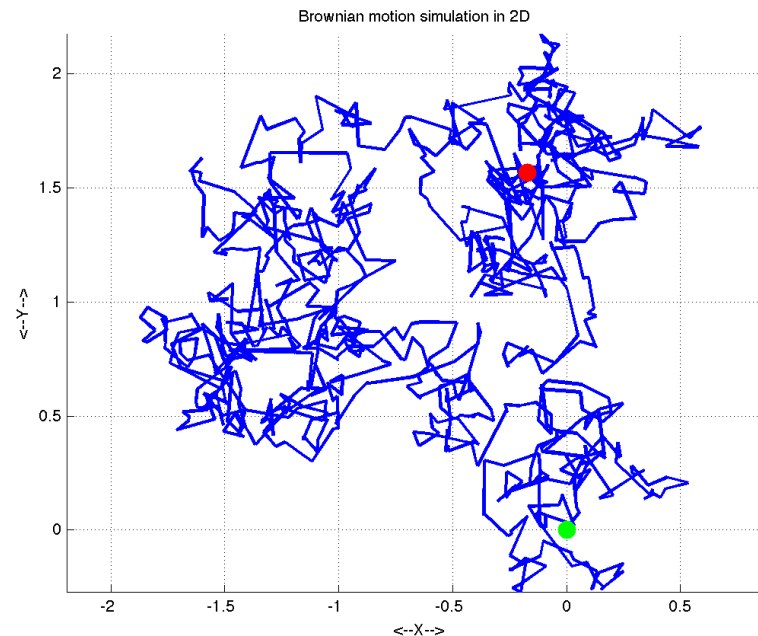
Edge QUE for Wigner ensembles [Knowles-Yin].

Matrix Dyson Brownian Motion (Matrix-DBM)

Evolve the matrix with a matrix Ornstein-Uhlenbeck process :

$$dH_t = \frac{1}{\sqrt{N}} dB_t - \frac{1}{2} H_t dt$$

The distribution of $H_t \sim e^{-t/2} H_0 + \sqrt{1 - e^{-t}} V$ where V is a GOE.



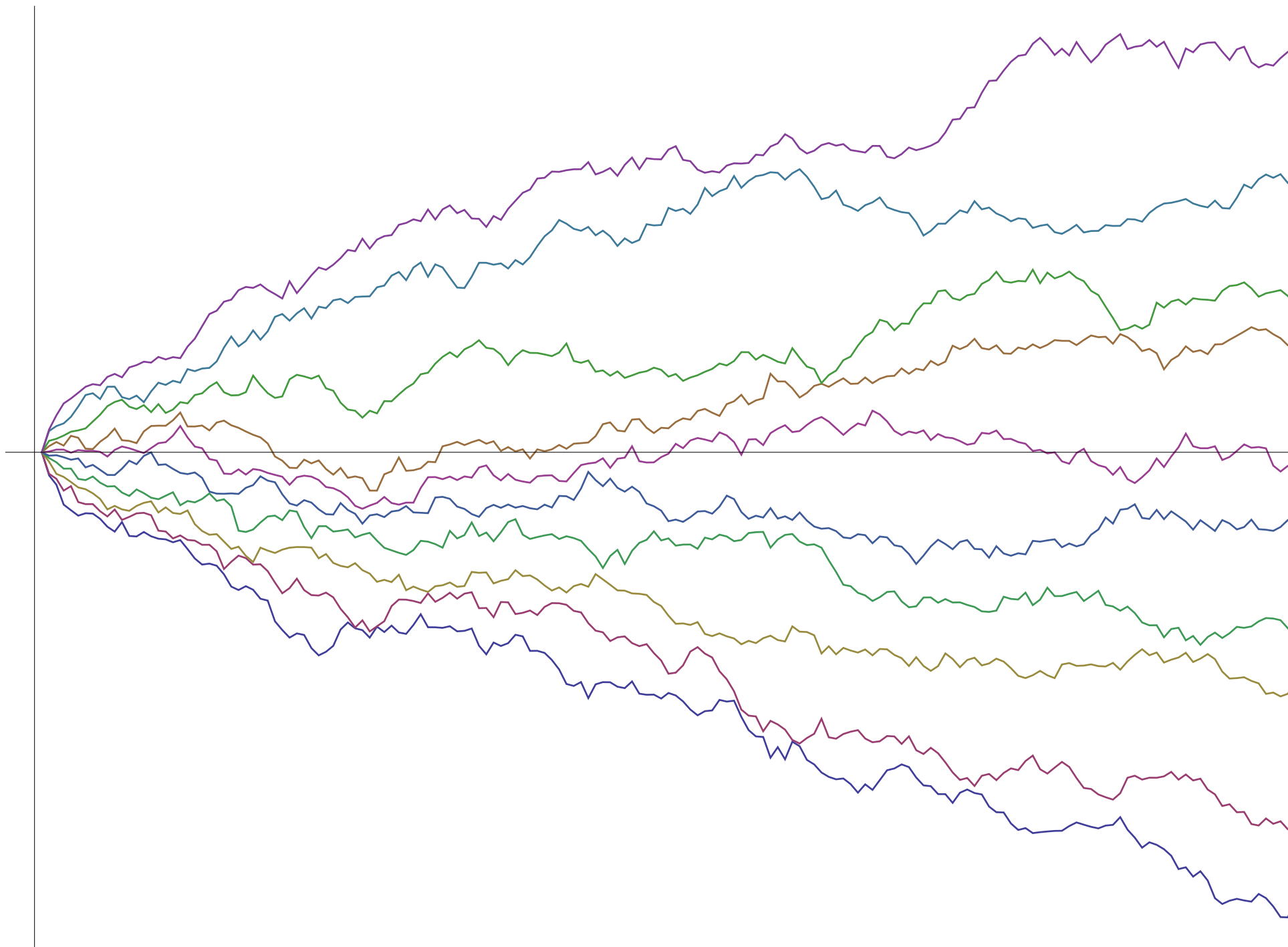
B_{ij} : symmetric independent Brownian motions

$$d\lambda_k = \frac{dB_{kk}}{\sqrt{N}} + \left(\frac{1}{N} \sum_{\ell \neq k} \frac{1}{\lambda_k - \lambda_\ell} - \frac{1}{2} \lambda_k \right) dt$$

Dyson Brownian Motion (1962) E-value equations are autonomous!

$$du_k = \frac{1}{\sqrt{N}} \sum_{\ell \neq k} \frac{u_\ell dB_{k\ell}}{\lambda_k - \lambda_\ell} - \frac{1}{2N} \sum_{\ell \neq k} \frac{u_k dt}{(\lambda_k - \lambda_\ell)^2}$$

Dyson e-vector flow E-vector equations depend on the e-values.



Dyson: The classical Coulomb gas is invariant under the DBM:

$$\mu_\beta \sim e^{-\beta N \mathcal{H}(\lambda)}, \quad \mathcal{H} = \sum_i \frac{\lambda_i^2}{4} - \frac{1}{N} \sum_{i < j} \log(\lambda_j - \lambda_i)$$

prob. density for the classical ensembles with $\beta = 1$ for the GOE

Dyson's conjecture: Time to "local equilibrium" for DBM is $\gtrsim N^{-1}$.

Erdős-Schlein-Y-Yin, 2009-10: **Dyson conjecture holds for $\beta > 0$.**

The Coulomb interactions drive the system locally very fast!

\implies Wigner-Dyson-Mehta conjecture holds for H_t with $t \gtrsim N^{-1}$.

Question: How to connect to $t = 0$?

Theorem* [Continuity of matrix-DBM]

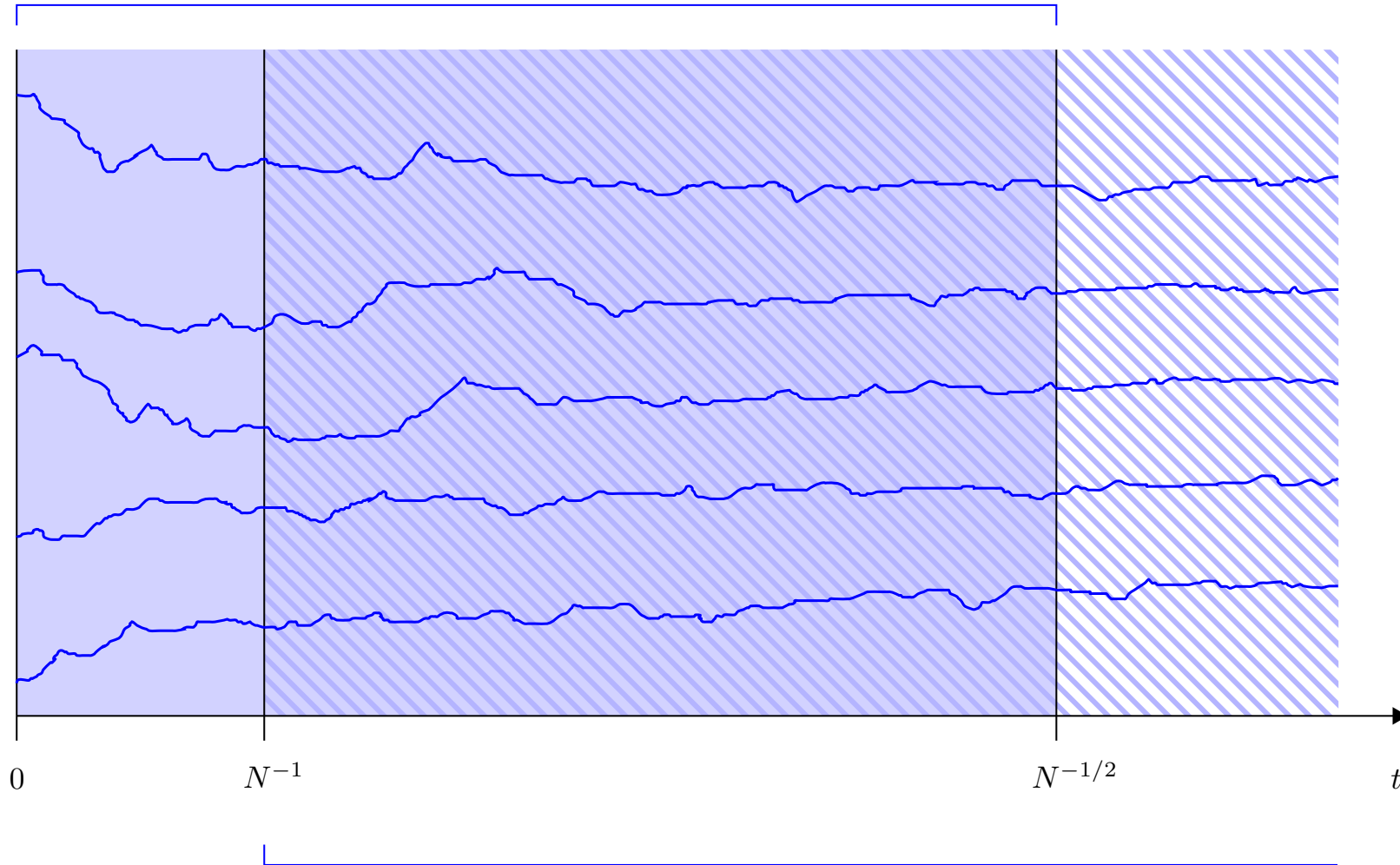
F : function of matrix elements with uniformly bounded 3rd derivatives. If $t \ll N^{-1/2}$ then

$$\mathbb{E}F(H_t) - \mathbb{E}F(H_0) \rightarrow 0$$

Proof: By Ito's formula and local semicircle laws.

* [Bourgade-Y 2013 (also a special case of Erdos-Knowles-Y-Yin)]

Continuity of matrix DBM



Equilibration of DBM

No contradiction since we use the matrix structure.

[Bourgade-Y 2013] For any unit vector \mathbf{q} fixed, define

$$f(t, j) = E \left[|\mathbf{q} \cdot u_j(t)|^2 | \boldsymbol{\lambda}(\cdot) \right].$$

The **eigenvector moment flow** (in the rescaled time $t \rightarrow Nt$)

$$\partial_t f(t, j) = \sum_{k \neq j} \frac{f(t, k) - f(t, j)}{(\lambda_j - \lambda_k)^2(t)}$$

(non-local) Random walk in random environments with singular coefficients.

A very simple equation in the indices j .

There are analogues for higher moments.

Theorem [Bourgade-Erdos-Y 2012-13] With high probability, for any $|i - j| \ll t$,

$$|f(t, k) - f(t, j)| \leq C \|f(t)\|_\infty t^{-\varepsilon}$$

ε is a Hölder regularity exponent from **De Giorgi-Nash-Moser** (Caffarelli-Chan-Vasseur) theory!

\implies **local QUE holds for H_t** ; $|\mathbf{q} \cdot u_j(t)|^2 \rightarrow N(0, 1)^2$.

How can one get Gaussian by "regularity"?

By continuity of matrix-DBM \implies local QUE holds for H_0 .

The mechanism underlying of the **universality** and **QUE** for random matrices: **Dyson Brownian motion (A dynamical idea)**

Universality: *Fast relaxation of DBM*

+ continuity of matrix-DBM

QUE: *Holder regularity of e-vector moment flow*

+ continuity of matrix-DBM

- Dynamical idea to solve time independent problems: *study the limits of the flow as $t \rightarrow \infty$.*
- Dyson Brownian Motion: *study the initial layer to understand $t = 0$.*

If you admit that the Wigner ensemble gives a completely wrong answer for the level density, why do you believe any of the other predictions of random-matrix theory?



George Uhlenbeck

- Local theory for universality was developed. A general theory ...